Kinetic Decoupling of Dark Matter

and its impact on the Relic Density

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based on: **T. Binder, T. Bringmann, M. Gustafsson and AH,**
Phys.Rev. D96 (2017) 115010, astro-ph.co/1706.07433

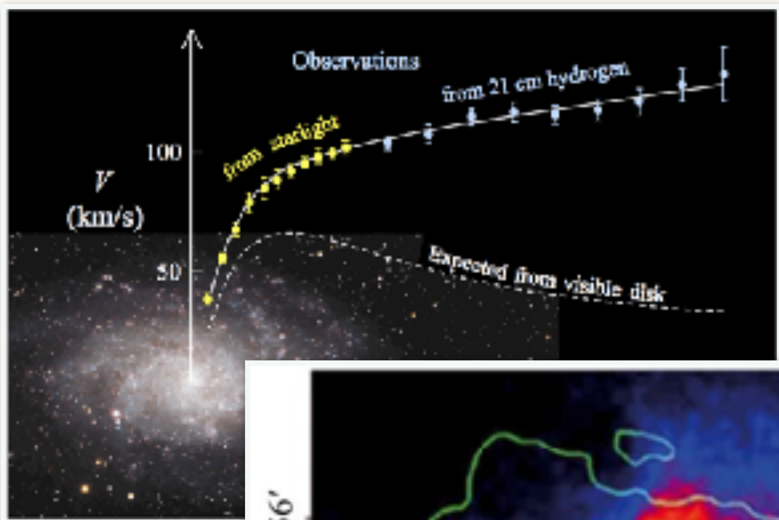
University of Warsaw, 21st December 2017

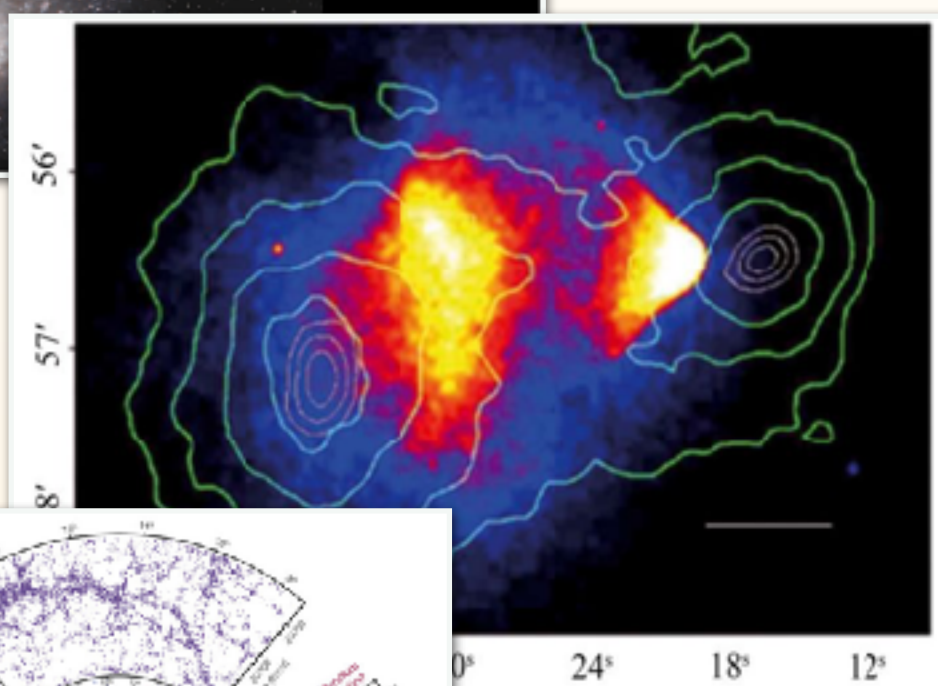
OUTLINE

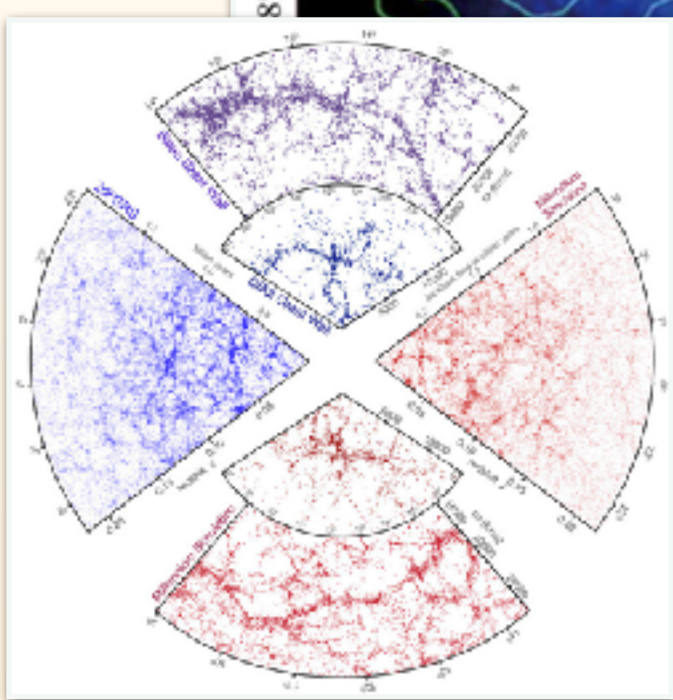
1. Introduction
   - standard approach to thermal relic density
   - recent novel models/ideas
2. Kinetic decoupling
   - freeze-out vs. decoupling
   - significance for cosmology
3. ***n**-th* Exception
   - early kinetic decoupling with
   - velocity dependent annihilation
4. Summary

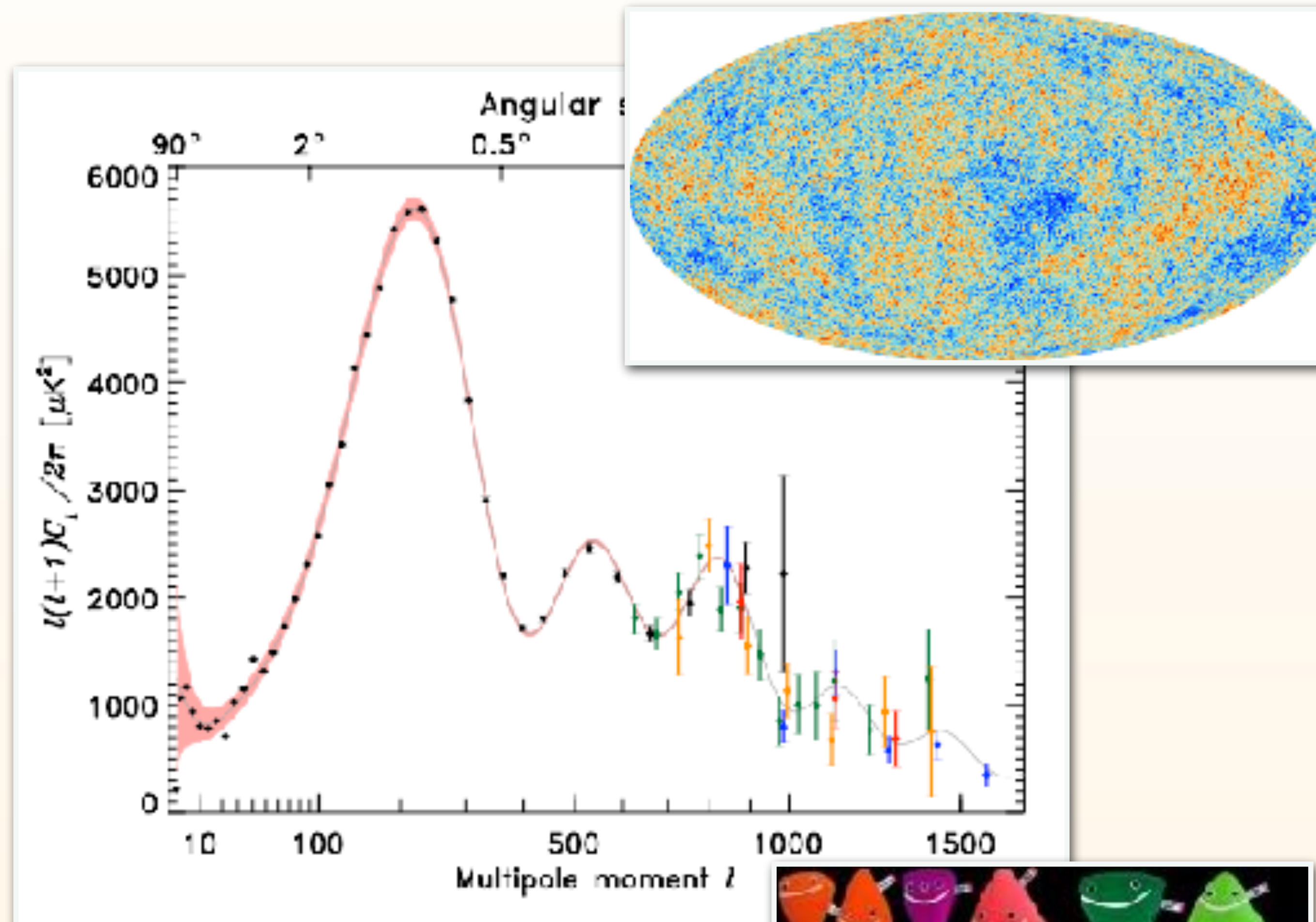
DARK MATTER

IS EVERYWHERE!









$\Longrightarrow$ Evidence on all scales!

The Origin of Dark Matter

Dark matter could be created in many different ways…

…but <u>every</u> massive particle with not-too-weak interactions with the SM will be produced thermally, with relic abundance:

Lee, Weinberg '77; + others

$$\Omega_\chi h^2 \approx 0.1 \, \frac{3 \times 10^{-26} \mathrm{cm}^3 \mathrm{s}^{-1}}{\langle \sigma v \rangle}$$

It is very natural to expect that this mechanism is responsible for the origin of all of dark matter

…but even if not, it still is present nevertheless and it's important to be able to correctly determine thermal population abundance

Historical Prelude

Three Exceptions

Griest & Seckel '91

1. Co-annihilations

if more than one state share a conserved quantum number making DM stable

$$\langle \sigma_{\rm eff} v \rangle = \sum_{ij} \langle \sigma_{ij} v_{ij} \rangle \frac{n_i^{\rm eq} n_j^{\rm eq}}{n_{\rm eq}^2}$$

with: $\sigma_{ij} = \sum_X \sigma(\chi_i \chi_j \to X)$

e.g., SUSY

2. Annihilation to forbidden channels

if DM is slightly below mass threshold for annihilation $\Longrightarrow$ „forbidden" channel can still be accessible in thermal bath

recent e.g., 1505.07107

3. Annihilation near poles

expansion in velocity (s-wave, p-wave, etc.) not safe

(more historical issue: these days most people use numerical codes)

Thermal Relic Density

modern "Exceptions"

1. Non-standard cosmology

   many works… very recent e.g., D'Eramo, Fernandez, Profumo '17

2. Bound State Formation

   recent e.g., Petraki at al. '15, '16; An et al. '15, '16; Cirelli et al. '16; …

3. $3 \to 2$ and $4 \to 2$ annihilation

   e.g., D'Agnolo, Ruderman '15; Cline at al. '17; Choi at al. '17; …

4. Second era of annihilation

   Feng et al. '10; Bringmann et al. '12; …

5. Semi-annihilation

   D'Eramo, Thaler '10; …

6. Cannibalization

   e.g., Kuflik et al. '15; Pappadopulo et al. '16; …

7. …

In other words: whenever studying non-minimal scenarios "exceptions" appear — but most of them come from interplay of new added effects, while do not affect the foundations of modern calculations

THERMAL RELIC DENSITY

STANDARD APPROACH

T ↑		time ↓
$\Gamma_{\rm ann} > H$	DM in equilibrium
$\Gamma_{\rm ann} \sim H$	**chemical decoupling**
$\Gamma_{\rm ann} < H$	freeze-out

non-equilibrium

time evolution of $f_\chi(p)$ in kinetic theory:

$$E\left(\partial_t - H\vec{p}\cdot\nabla_{\vec{p}}\right) f_\chi = \mathcal{C}[f_\chi] \quad \Rightarrow \quad \frac{dn_\chi}{dt} + 3Hn_\chi = C$$

Liouville operator in FRW background

the collision term (integrated)

assumptions for using Boltzmann eq: classical limit, molecular chaos,...

…for derivation from thermal QFT see e.g., M. Beneke, F. Dighera, AH; JHEP 1410 (2014) 45

THERMAL RELIC DENSITY

THE COLLISION TERM

for $2 \leftrightarrow 2$ CP invariant process:

$$C_{\rm LO} = -h_\chi^2 \int \frac{d^3\vec{p}_\chi}{(2\pi)^3} \frac{d^3\vec{p}_{\bar\chi}}{(2\pi)^3}\, \sigma_{\chi\bar\chi\to ij} v_{\rm rel} \; [f_\chi f_{\bar\chi}(1\pm f_i)(1\pm f_j) - f_i f_j (1\pm f_\chi)(1\pm f_{\bar\chi})]$$

assuming kinetic equilibrium at chemical decoupling: $f_\chi \sim a(\mu) f_\chi^{\rm eq}$

$$C_{\rm LO} = -\langle \sigma_{\chi\bar\chi\to ij} v_{\rm rel}\rangle^{\rm eq} \left(n_\chi n_{\bar\chi} - n_\chi^{\rm eq} n_{\bar\chi}^{\rm eq}\right)$$

where the thermally averaged cross section:

$$\langle \sigma_{\chi\bar\chi\to ij} v_{\rm rel}\rangle^{\rm eq} = -\frac{h_\chi^2}{n_\chi^{\rm eq} n_{\bar\chi}^{\rm eq}} \int \frac{d^3\vec{p}_\chi}{(2\pi)^3} \frac{d^3\vec{p}_{\bar\chi}}{(2\pi)^3}\, \sigma_{\chi\bar\chi\to ij} v_{\rm rel}\; f_\chi^{\rm eq} f_{\bar\chi}^{\rm eq}$$

THERMAL RELIC DENSITY

BOLTZMANN EQ.

$$\frac{dn_\chi}{dt} + 3Hn_\chi = -\langle \sigma_{\chi\bar{\chi}\to ij}\sigma_{\rm rel}\rangle^{\rm eq}\left(n_\chi n_{\bar{\chi}} - n_\chi^{\rm eq} n_{\bar{\chi}}^{\rm eq}\right)$$

Re-written for the comoving number density:

$$\frac{dY}{dx} = \sqrt{\frac{g_*\pi m_\chi^2}{45G}}\,\frac{\langle \sigma_{\chi\bar{\chi}\to ij}\sigma_{\rm rel}\rangle^{\rm eq}}{x^2}\left(Y^2 - Y_{\rm eq}^2\right)$$

$$\lim_{x\to 0} Y = Y_{\rm eq} \qquad \lim_{x\to\infty} Y = \text{const}$$

Recipe:
compute annihilation cross-section,
take a thermal bath average,
throw it into BE… and voilà

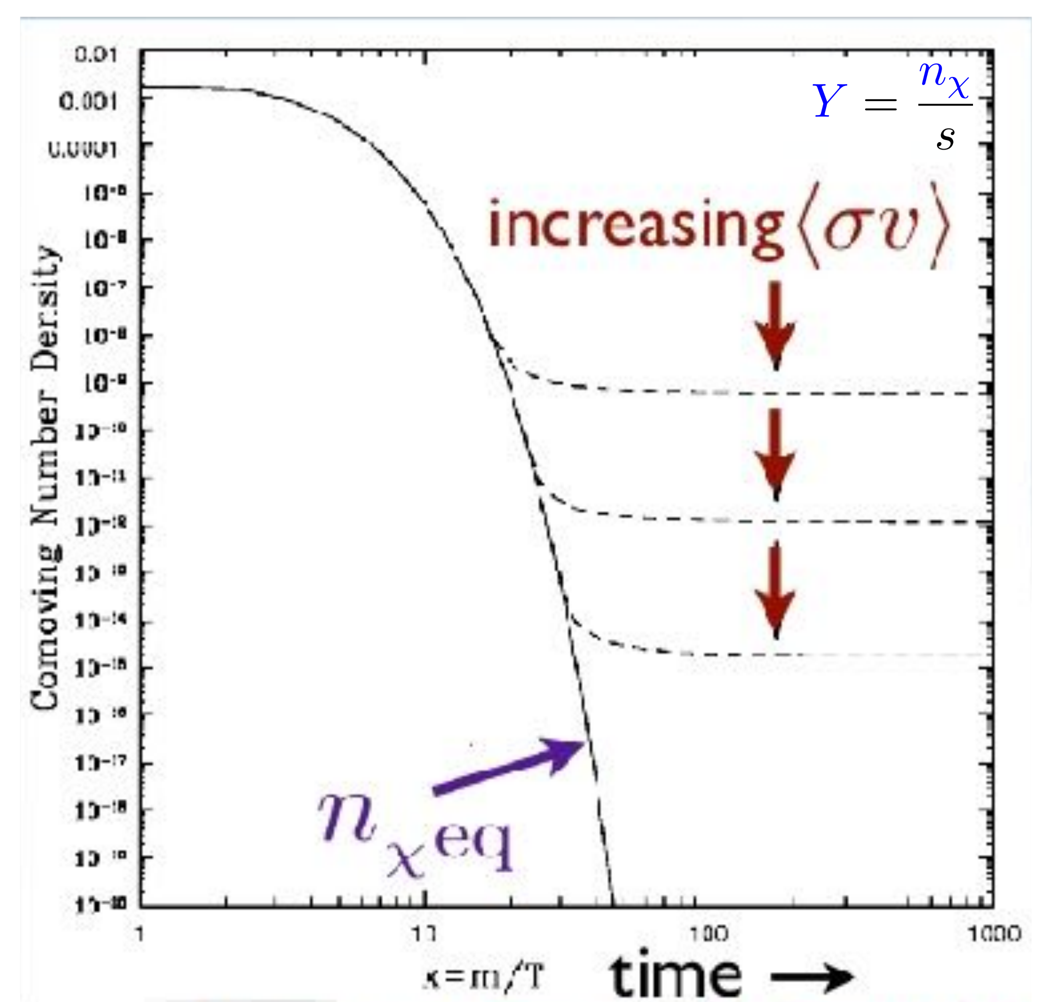
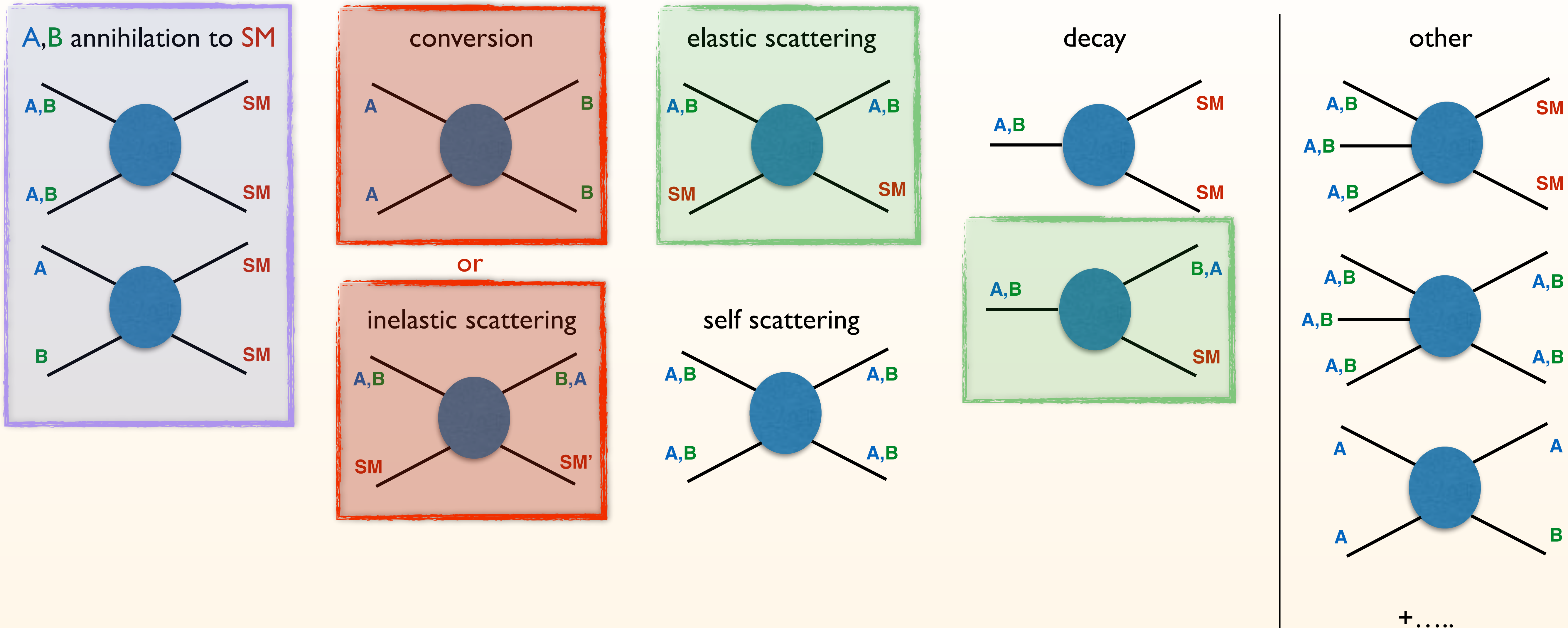


Fig.: Jungman, Kamionkowski & Griest, PR'96

WHAT IF NON-MINIMAL SCENARIO?



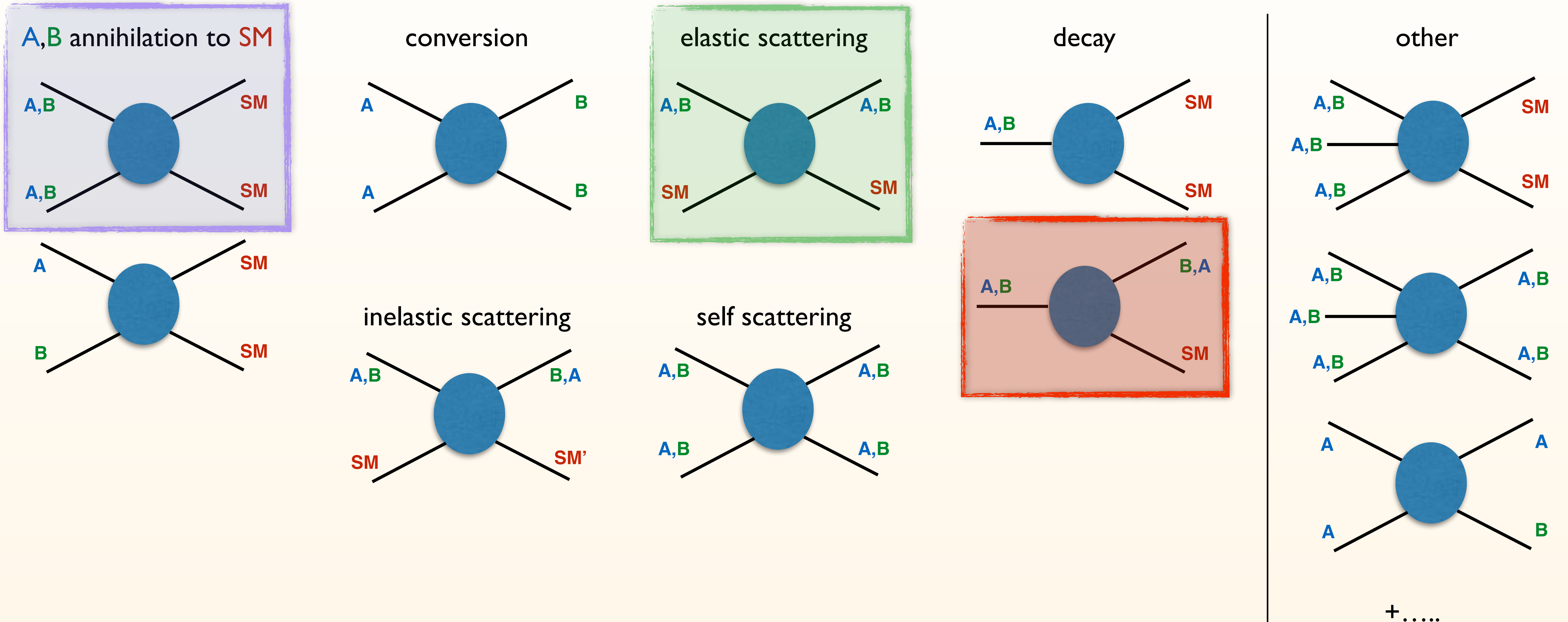
↳ typically forbidden by symmetry

Co-annihilation ⟶ due to efficient conversion processes one can trace only number density of sum of the states with shared conserved quantum number using weighted annihilation cross section

Griest, Seckel '91

- needed to be efficient for mechanism to work
- setting the relic density
- assumed in computation

WHAT IF NON-MINIMAL SCENARIO?



+…..

↳ typically forbidden by symmetry

SuperWIMP → DM abundance is an effect of complete decay of heavier state which freezes-out as standard WIMP

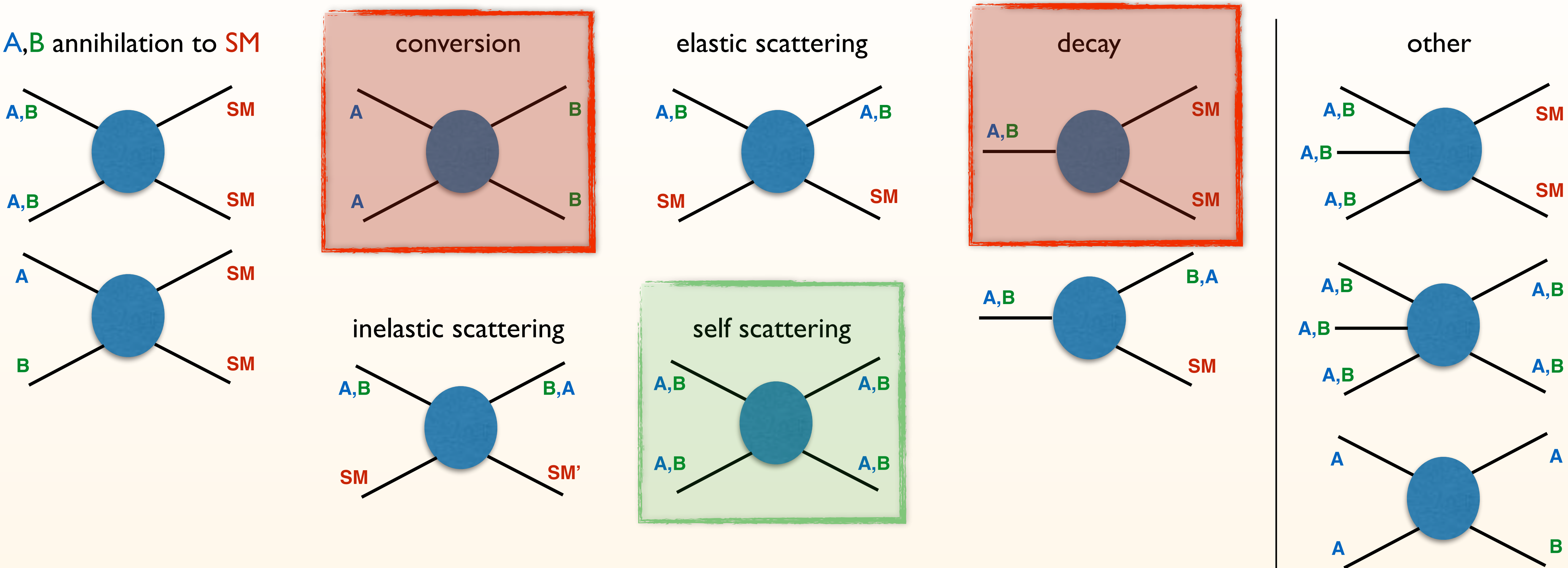
Feng, Rajaraman, Takayama '03

needed to be efficient for mechanism to work

setting the relic density

assumed in computation

What if non-minimal scenario?



+…..

↳ typically forbidden by symmetry

Co-decaying → DM decouples when relativistic but then one of the dark sector states decays and this effect important as long as conversions are

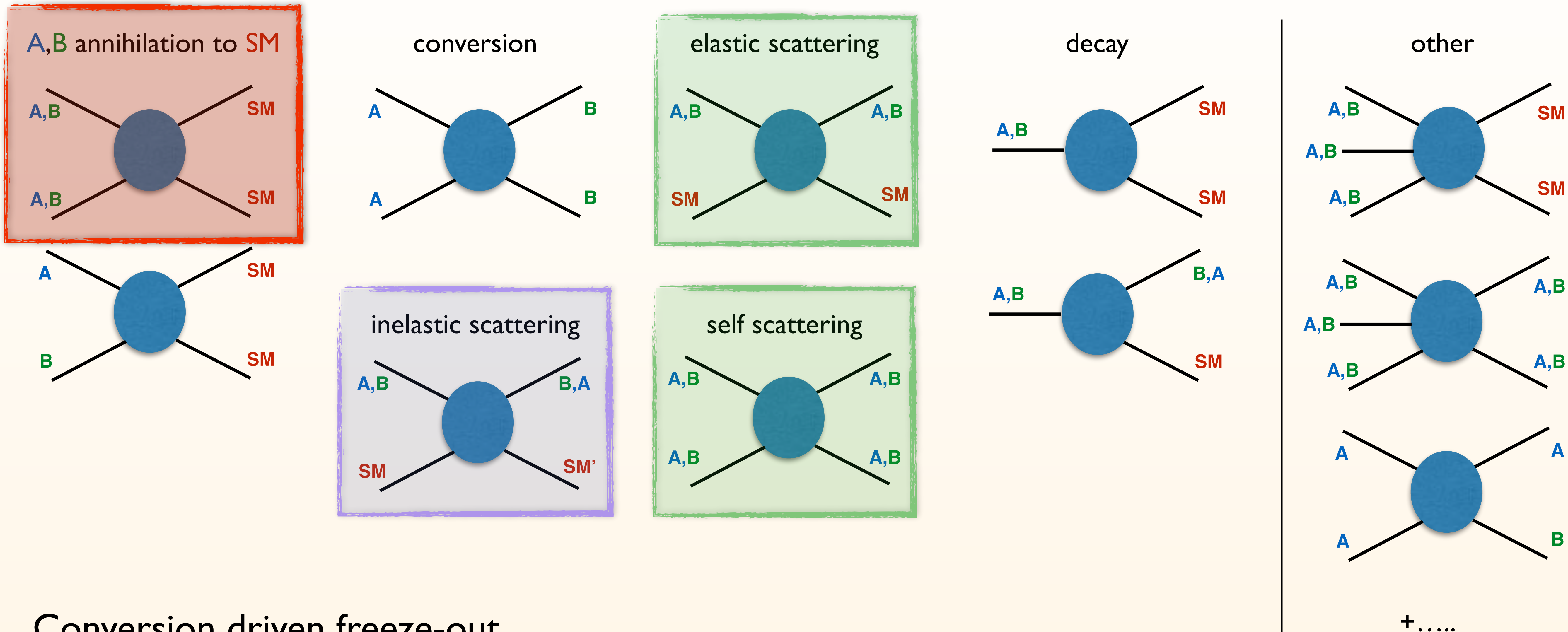
Dror, Kuflik, Ng '16

needed to be efficient for mechanism to work

setting the relic density

assumed in computation

WHAT IF NON-MINIMAL SCENARIO?



Conversion driven freeze-out

Garny, Heisig, Lulf, Vogl '17

Co-scattering

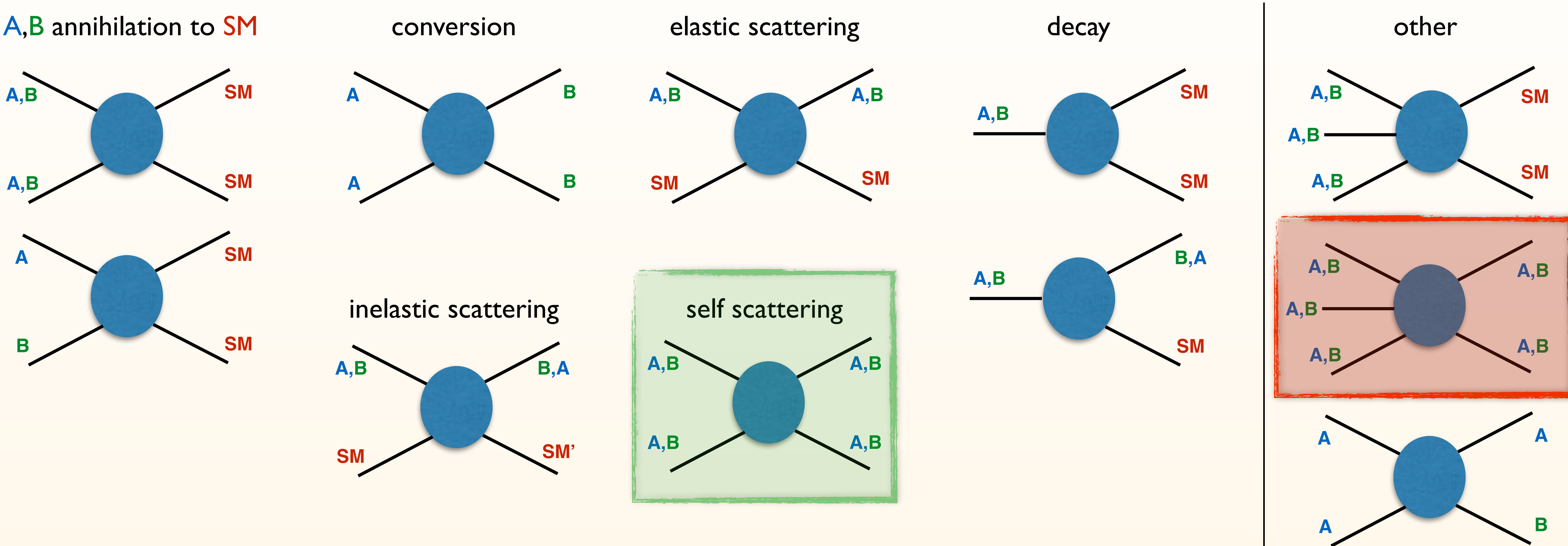
D'Agnolo, Pappadopulo, Ruderman '17

→ only one of the dark sector states annihilates efficiently, but also conversions stop being efficient which blocks co-annihilation

↳ typically forbidden by symmetry

- needed to be efficient for mechanism to work (red box)
- setting the relic density (purple box)
- assumed in computation (green box)

WHAT IF NON-MINIMAL SCENARIO?



+…..

typically forbidden by symmetry

Cannibalization → secluded dark sector, with efficient 3-2 annihilation leading to DM zero chemical potential and keeping DM at much higher temperature than SM plasma

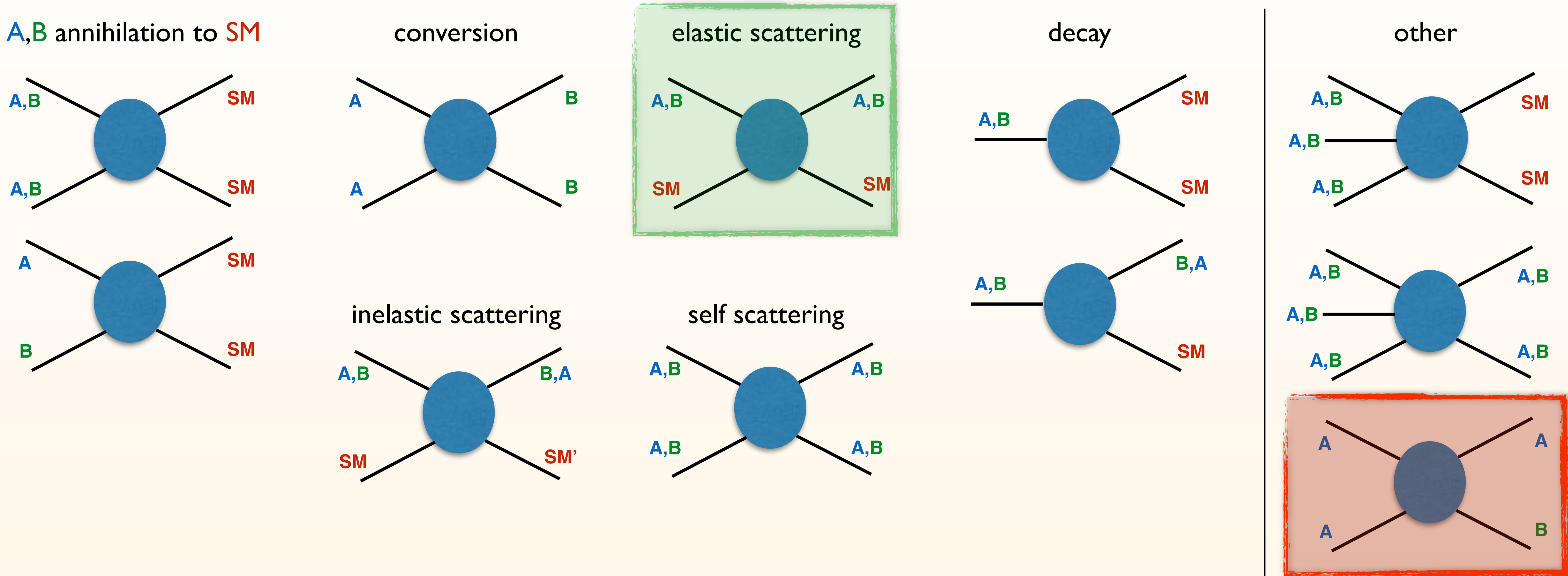
Kuflik et al. '15

needed to be efficient for mechanism to work

setting the relic density

assumed in computation

WHAT IF NON-MINIMAL SCENARIO?



+.....

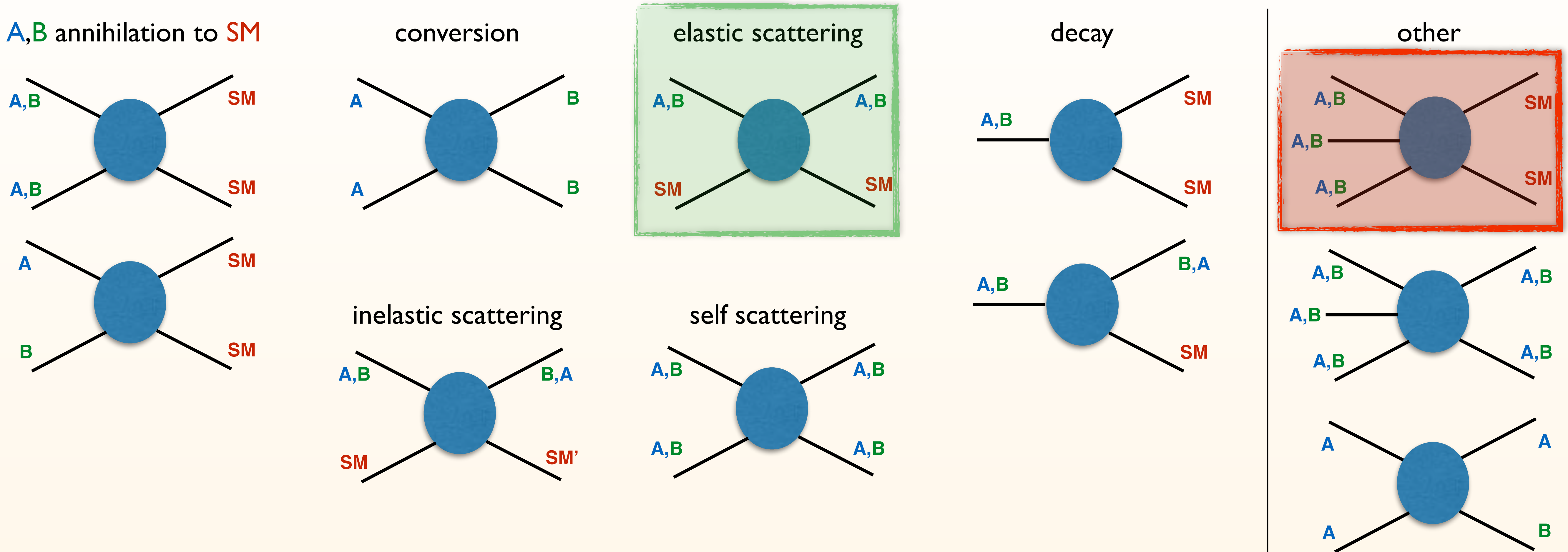
↳ typically forbidden by symmetry

Semi-annihilation ⟶ new type of annihilation precess that can dominate the freeze-out dynamics; occurs when new „flavour" or „baryon" structure in dark sector

D'Eramo, Thaler '10

- needed to be efficient for mechanism to work
- setting the relic density
- assumed in computation

What if non-minimal scenario?



Forbidden-like DM ⟶ annihilation to a state heavier than DM, possible from tail of the Boltzmann distribution or when 3-2 processes are non-negligible

e.g., D'Agnolo, Ruderman '15; Cline at al. '17; Choi at al. '17;

+…..

↳ typically forbidden by symmetry

needed to be efficient for mechanism to work

setting the relic density

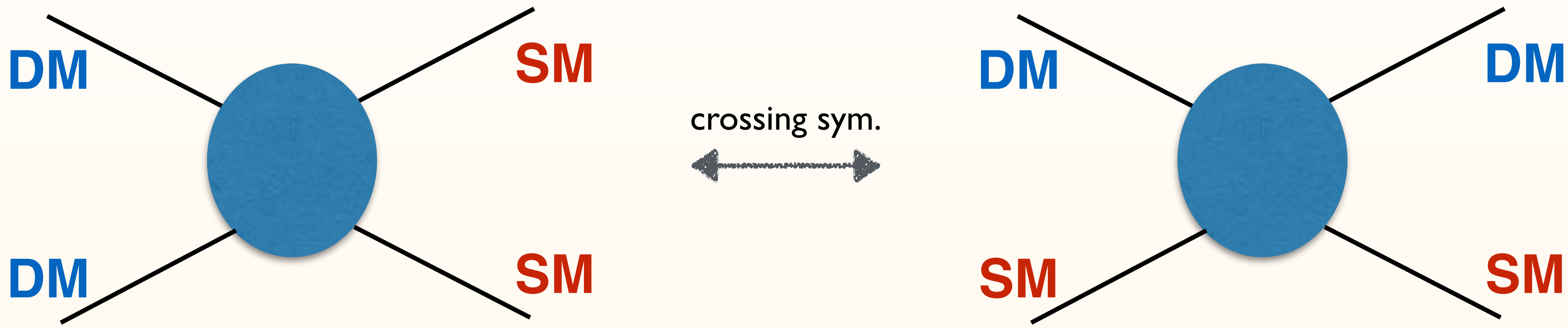
assumed in computation

Exception *N*:

Early kinetic decoupling

FREEZE-OUT *VS*. DECOUPLING

annihilation (elastic) scattering



$$\sum_{\text{spins}} \left|\mathcal{M}^{\text{pair}}\right|^2 = F(p_1, p_2, p_1', p_2') \quad \sim \quad \sum_{\text{spins}} \left|\mathcal{M}^{\text{scatt}}\right|^2 = F(k, -k', p', -p)$$

Boltzmann suppression of DM vs. SM $\Rightarrow$ scatterings typically more frequent

dark matter frozen-out but typically still kinetically coupled to the plasma

$$\tau_r(T_{\text{kd}}) \equiv N_{\text{coll}}/\Gamma_{\text{el}} \sim H^{-1}(T_{\text{kd}})$$

Schmid, Schwarz, Widern '99; Green, Hofmann, Schwarz '05

Two consequences:

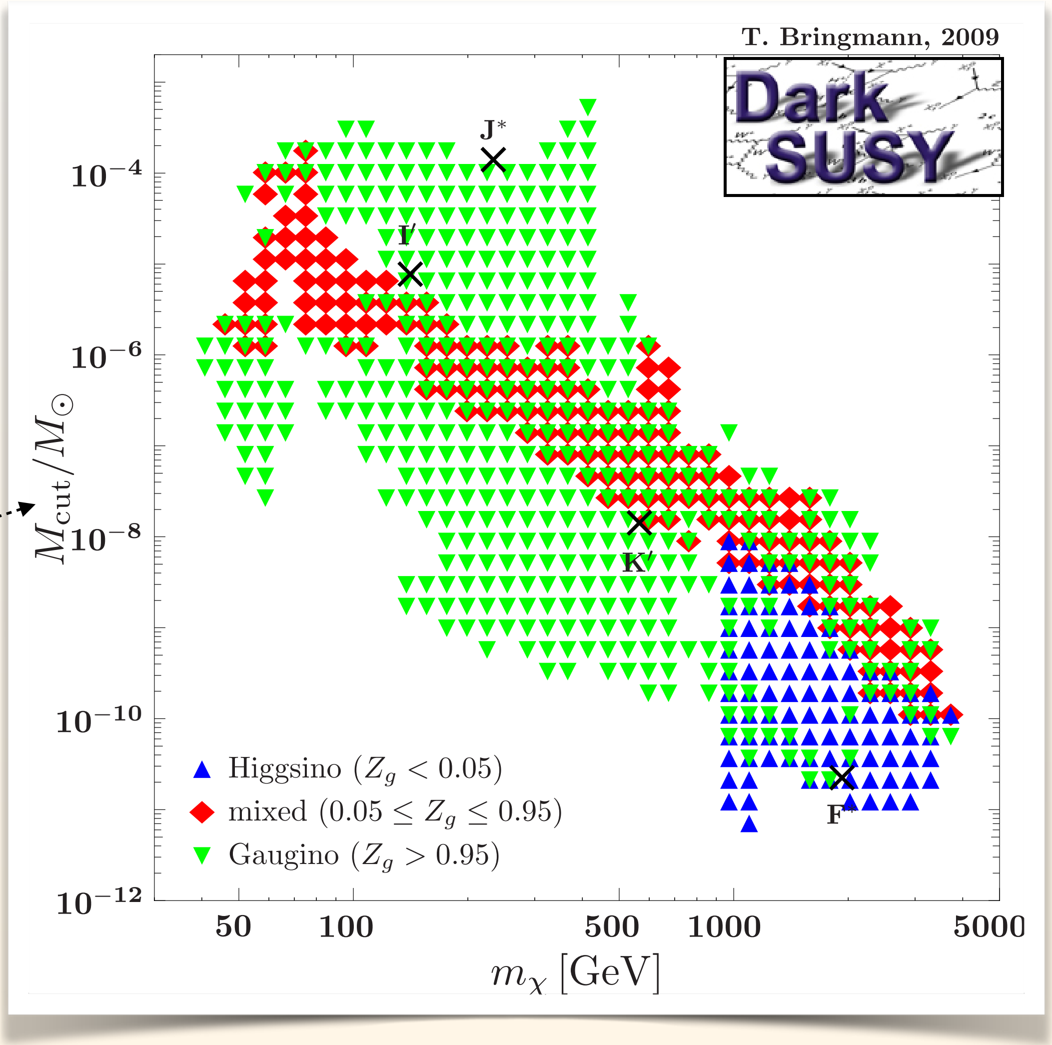
1. During freeze-out (chemical decoupling) typically: $f_\chi \sim a(\mu) f_\chi^{\text{eq}}$
2. If kinetic decoupling much, much later: possible impact on the matter power spectrum

i.e. kinetic decoupling can have observable consequences and affect e.g. missing satellites problem

see e.g., Bringmann, Ihle, Karsten, Walia '16

IMPLICATIONS OF KINETIC DECOUPLING

E.g. for SUSY neutralino: Bringmann '09



Free-streaming of DM after KD washes out density contrasts at small scales (similarly to baryonic oscillations)

Green, Hofmann, Schwarz '05

Cut-off in the power spectrum corresponding to smallest gravitationally bound objects

„Typical" values for WIMPs are relatively small → small substructures expected

→ but bad for missing satellites problem

$\Longrightarrow$ moment of KD leaves important imprint on the Universe

A Pitfall in a Nutshell

T ↑ | $\Gamma_{\mathrm{ann}} > H$ | DM in equilibrium | time ↓

$\Gamma_{\mathrm{ann}} \sim H$ | chemical decoupling

$\Gamma_{\mathrm{ann}} < H$ | freeze-out

can it happen? / typically

kinetic decoupling

If KD happens around CD → what would be the relic density?

↓

how to even compute that? ⇒ need for refined treatment of solving the Boltzmann eq.

assuming kinetic equilibrium at chemical decoupling: $f_\chi \sim a(\mu) f_\chi^{\mathrm{eq}}$

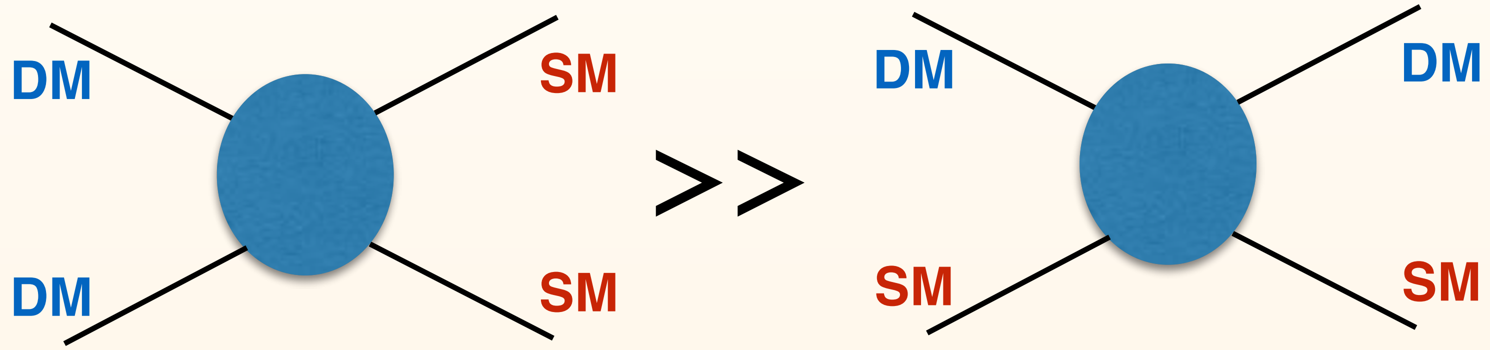
$$C_{\mathrm{LO}} = -\langle \sigma_{\chi\bar{\chi}\to ij} v_{\mathrm{rel}} \rangle^{\mathrm{eq}} \left(n_\chi n_{\bar{\chi}} - n_\chi^{\mathrm{eq}} n_{\bar{\chi}}^{\mathrm{eq}} \right)$$

EARLY KINETIC DECOUPLING?

A necessary and sufficient condition: scatterings weaker than annihilation

i.e. rates around freeze-out: $H \sim \Gamma_{\rm ann} \gtrsim \Gamma_{\rm el}$

Possibilities:

A) 

e.g., resonant annihilation

B) Boltzmann suppression of SM as strong as for DM

e.g., below threshold annihilation (forbidden-like DM)

C) Scatterings and annihilation have different structure

e.g., semi-annihilation, 3 to 2 models,…

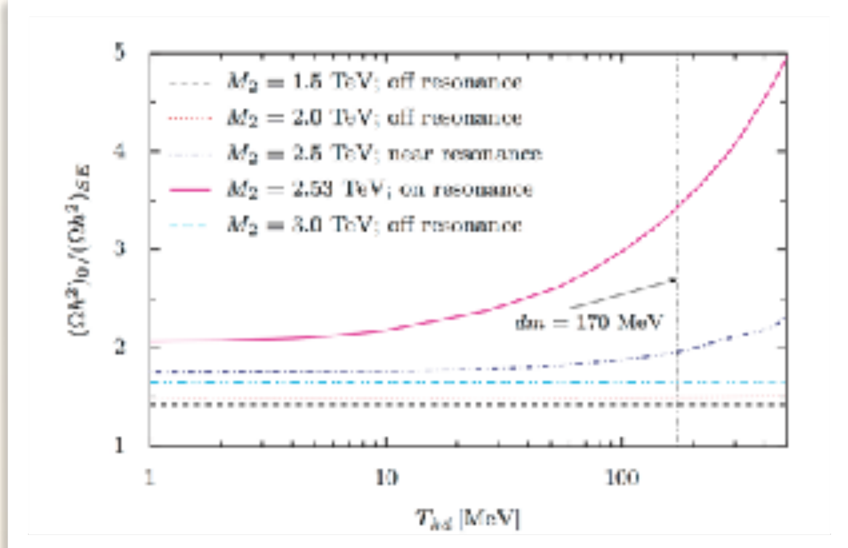
Early KD and resonance

our work wasn't the first to realize that resonant annihilation can lead to early kinetic decoupling…

Feng, Kaplinghat, Yu '10 — noted that for Sommerfeld-type resonances KD can happen early

Dent, Dutta, Scherrer '10 — looked at potential effect of KD on thermal relic density

Since then people were aware of this effect and sometimes tried to estimate it assuming instantaneous KD, e.g., in the case of Sommerfeld effect in the MSSM:



AH, Iengo, Ullio '11

but no systematic studies of decoupling process were performed, until…

…models with very late KD become popular, in part to solve „missing satellites" problem

van den Aarssen et al '12; Bringmann et al '16, x2; Binder et al '16

this progress allowed for better approach to early KD scenarios as well and was applied to the resonant annihilation case in

Duch, Grządkowski '17

… but we developed a dedicated accurate method/code to deal with this and other similar situations

COMMENT:

COMPLEMENTARITY OF DG VS. BBGH

The effect of early KD was approached from two very different perspectives:

Duch, Grządkowski '17	Binder, Bringmann, Gustafsson, AH '17
motivated by phenomenology of self-interacting DM	motivated by the question: can KD happen before CD?
evolved into more in-depth study of effects related to general resonant annihilation: KD and energy dependent width	evolved into more in-depth study of the time evolution of the phase space distribution function for a general case of velocity dependent annihilation
as an example case studied a vector DM model and looked more onto self-interactions	as an example case studied a very specific resonance model and concentrated on the relic density

How to describe KD?

All information is in full BE:

both about chemical ("normalization") and kinetic ("shape") equilibrium/decoupling

$$E\,(\partial_t - H\vec{p}\cdot\nabla_{\vec{p}})\, f_\chi = \mathcal{C}[f_\chi]$$

contains both scatterings and annihilation

Two possible approaches:

solve numerically for full $f_\chi(p)$

have insight on the distribution

no constraining assumptions

numerically challenging

typically overkill

consider system of equations for moments of $f_\chi(p)$

0-th moment: n_χ

2-nd moment: T_χ

...

partially analytic/much easier numerically

manifestly captures all of the relevant physics

finite range of validity

no insight on the distribution

SCATTERING

The elastic scattering collision term:

$$C_{\rm el} = \frac{1}{2g_\chi}\int\frac{d^3k}{(2\pi)^3 2\omega}\int\frac{d^3\tilde{k}}{(2\pi)^3 2\tilde{\omega}}\int\frac{d^3\tilde{p}}{(2\pi)^3 2\tilde{E}}$$
$$\times(2\pi)^4\delta^{(4)}(\tilde{p}+\tilde{k}-p-k)|\mathcal{M}|^2_{\chi f\leftrightarrow\chi f}$$
$$\times\left[(1\mp g^\pm)(\omega)\,g^\pm(\tilde{\omega})f_\chi(\tilde{\mathbf{p}})-(\omega\leftrightarrow\tilde{\omega},\mathbf{p}\leftrightarrow\tilde{\mathbf{p}})\right]$$

→ equilibrium functions for SM particles

Expanding in **NR** and small **momentum transfer**:

Bringmann, Hofmann '06

$$C_{\rm el}\simeq\frac{m_\chi}{2}\gamma(T)\left[Tm_\chi\partial_p^2+\left(p+2T\frac{m_\chi}{p}\right)\partial_p+3\right]f_\chi$$

physical interpretation: scattering rate

More generally, Fokker-Planck scattering operator (relativistic, but still small **momentum transfer**):

Binder et al. '16

$$C_{\rm el}\simeq\frac{E}{2}\nabla_{\mathbf{p}}\cdot\left[\gamma(T,\mathbf{p})\left(ET\nabla_{\mathbf{p}}+\mathbf{p}\right)f_\chi\right]$$

Semi-relativistic: assume that scattering $\gamma(T,\mathbf{p})$ is momentum independent

KINETIC DECOUPLING 101

DM temperature Definition:

$$T_\chi \equiv \frac{g_\chi}{3 m_\chi n_\chi} \int \frac{\mathrm{d}^3 p}{(2\pi)^3} p^2 f_\chi(p) \qquad y \equiv \frac{m_\chi T_\chi}{s^{2/3}}$$

→ actually: normalized average NR energy - equals temperature at equilibrium

First take late KD scenario and consider only temperature evolution - i.e. leave out feedback **on/from** changing number density:

then 2nd moment of full BE (up to terms p^2/m_χ^2) gives:

$$\frac{y'}{y} = -\frac{Y'}{Y}\left(1 - \frac{\langle \sigma v_{\rm rel}\rangle_2}{\langle \sigma v_{\rm rel}\rangle}\right) - \left(1 - \frac{x}{3}\frac{g_{*S}'}{g_{*S}}\right)\frac{2 m_\chi c(T)}{Hx}\left(1 - \frac{y_{\rm eq}}{y}\right)$$

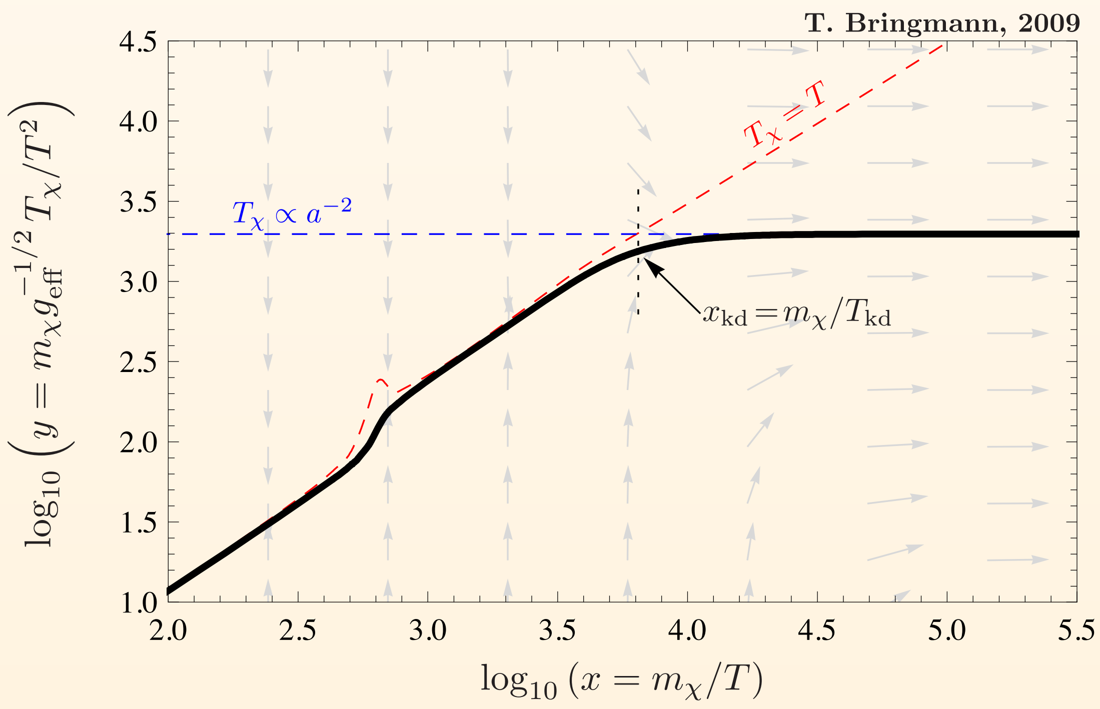
where:

$$\langle \sigma v_{\rm rel}\rangle_2 \equiv \frac{g_\chi^2}{3 T m_\chi n_\chi^2} \int \frac{d^3 p}{(2\pi)^3} \int \frac{d^3 \tilde{p}}{(2\pi)^3}\, p^2 v_{\rm rel} \sigma_{\bar{\chi}\chi \to \bar{X} X} f(E) f(\tilde{E})$$

impact of annihilation

$$c(T) = \frac{1}{12(2\pi)^3 m_\chi^4 T} \sum_X \int dk\, k^5 \omega^{-1} g^{\pm}\left(1 \mp g^{\pm}\right) \int_{-4k^2}^{0} (-t) \frac{1}{8k^4} |\mathcal{M}_{\rm el}|^2$$

impact of elastic scatterings



ONE STEP FURTHER…

Now consider general KD scenario, i.e. coupled temperature and number density evolution:

annihilation and production thermal averages done at different *T* — feedback of modified *y* evolution

$$\frac{Y'}{Y} = -\frac{1-\frac{x}{3}\frac{g'_{*S}}{g_{*S}}}{Hx} sY \left(\langle\sigma v_{\rm rel}\rangle|_{x=m_\chi^2/(s^{2/3}y)} - \frac{Y_{\rm eq}^2}{Y^2}\langle\sigma v_{\rm rel}\rangle|_x \right)$$

$$\frac{y'}{y} = -\frac{1-\frac{x}{3}\frac{g'_{*S}}{g_{*S}}}{Hx}\left[2m_\chi c(T)\left(1-\frac{y_{\rm eq}}{y}\right) - sY\left(\left(\langle\sigma v_{\rm rel}\rangle - \langle\sigma v_{\rm rel}\rangle_2\right)_{x=m_\chi^2/(s^{2/3}y)} - \frac{Y_{\rm eq}^2}{Y^2}\left(\langle\sigma v_{\rm rel}\rangle - \frac{y_{\rm eq}}{y}\langle\sigma v_{\rm rel}\rangle_2\right)_x \right)\right]$$

$$+\frac{1-\frac{x}{3}\frac{g'_{*S}}{g_{*S}}}{3m_\chi}\langle p^4/E^3\rangle_{x=m_\chi^2/(s^{2/3}y)}$$

"relativistic" term

$$T_\chi \equiv \frac{g_\chi}{3n_\chi}\int\frac{d^3p}{(2\pi)^3}\frac{p^2}{E}f_\chi(p)$$

elastic scatterings term

impact of annihilation

These equations still assume the equilibrium shape of $f_\chi(p)$ — but with variant temperature

or more accurately: that the thermal averages computed with true non-equilibrium distributions don't differ much from the above ones

NUMERICAL APPROACH

… or one can just solve full phase space Boltzmann eq.

$$\begin{aligned}\partial_x f_\chi(x,q) &= \frac{m_\chi^3}{\tilde{H}x^4}\frac{g_{\bar\chi}}{2\pi^2}\int d\tilde{q}\,\tilde{q}^2\,\frac{1}{2}\int d\cos\theta\; v_{\text{Møl}}\sigma_{\bar\chi\chi\to\bar f f}\\ &\quad\times\left[f_{\chi,\text{eq}}(q)f_{\chi,\text{eq}}(\tilde q)-f_\chi(q)f_\chi(\tilde q)\right]\\ &+\frac{2m_\chi c(T)}{2\tilde{H}x}\left[x_q\partial_q^2+\left(q+\frac{2x_q}{q}+\frac{q}{x_q}\right)\partial_q+3\right]f_\chi\\ &+\tilde g\frac{q}{x}\partial_q f_\chi,\end{aligned}$$

fully general

expanded in NR and small momentum transfer (semi-relativistic!)

discretization, ~1000 steps

$$\begin{aligned}\partial_x f_i &= \\ &\frac{m_\chi^3}{\tilde{H}x^4}\frac{g_{\bar\chi}}{2\pi^2}\sum_{j=1}^{N-1}\frac{\Delta\tilde q_j}{2}\Big[\tilde q_j^2\,\langle v_{\text{Møl}}\sigma_{\bar\chi\chi\to\bar f f}\rangle_{i,j}^{\theta}\left(f_i^{\text{eq}}f_j^{\text{eq}}-f_if_j\right)\\ &+\tilde q_{j+1}^2\,\langle v_{\text{Møl}}\sigma_{\bar\chi\chi\to\bar f f}\rangle_{i,j+1}^{\theta}\left(f_i^{\text{eq}}f_{j+1}^{\text{eq}}-f_if_{j+1}\right)\Big]\\ &+\frac{2m_\chi c(T)}{2\tilde{H}x}\left[x_{q,i}\partial_q^2+\left(q_i+\frac{2x_{q,i}}{q_i}+\frac{q_i}{x_{q,i}}\right)\partial_q+3\right]f_i\\ &+\tilde g\frac{q_i}{x}\partial_q f_i,\end{aligned}$$

Solved numerically with MatLab

Note:

can be extended to e.g. self-scatterings

very stiff, care needed with numerics

EXAMPLE:

SCALAR SINGLET DM

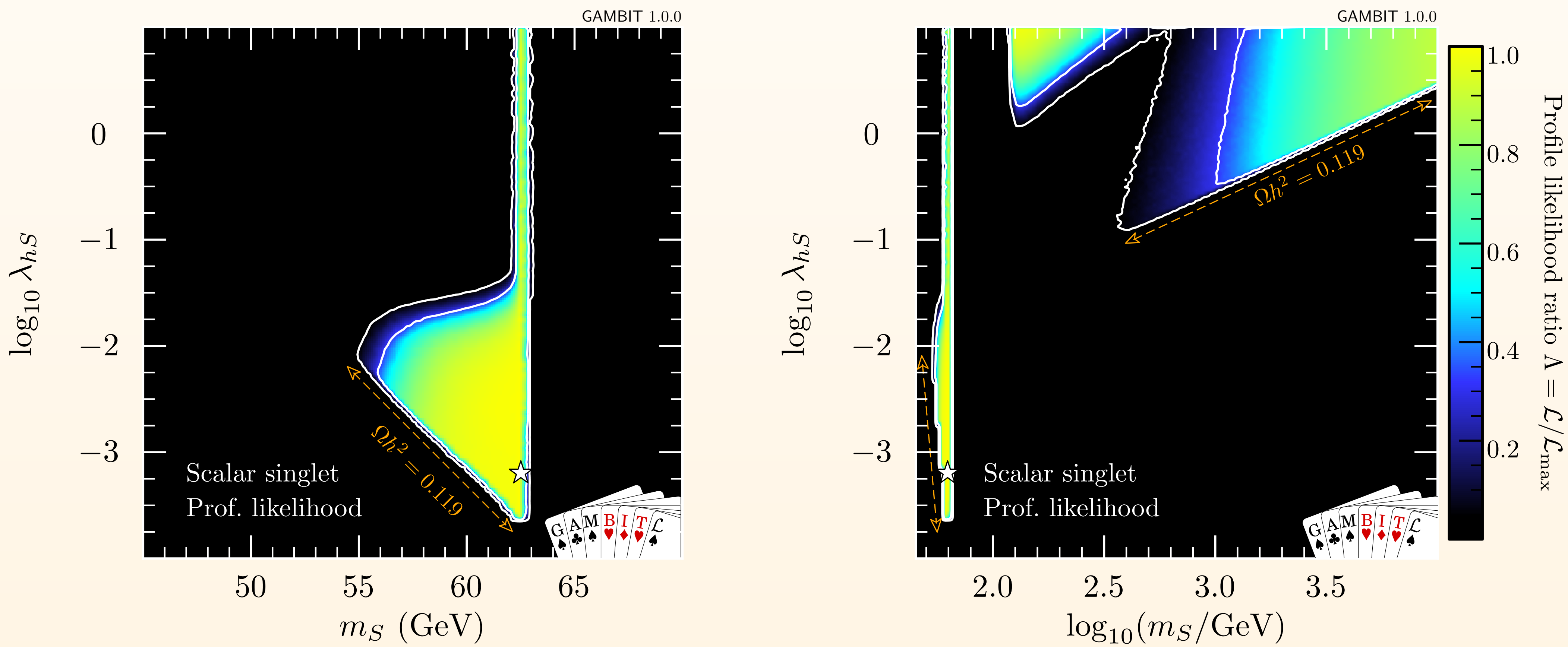
SCALAR SINGLET DM

VERY SHORT INTRODUCTION

To the SM Lagrangian add one singlet scalar field *S* with interactions with the Higgs:

$$\mathcal{L}_S = \frac{1}{2}\partial_\mu S \partial^\mu S - \frac{1}{2}\mu_S^2 S^2 - \frac{1}{2}\lambda_s S^2 |H|^2$$

$$m_s = \sqrt{\mu_S^2 + \frac{1}{2}\lambda_s v_0^2}$$



GAMBIT collaboration
1705.07931

Most of the parameter space excluded, but… even such a simple model is hard to kill

SCALAR SINGLET DM

ANNIHILATION VS. SCATTERINGS

$$\sigma v_{\rm rel} = \frac{2\lambda_s^2 v_0^2}{\sqrt{s}} |D_h(s)|^2 \Gamma_h(\sqrt{s})$$

tabulated Higgs width

with:

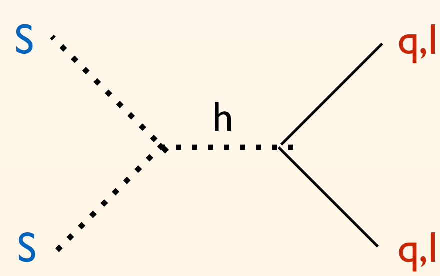
$$|D_h(s)|^2 \equiv \frac{1}{(s-m_h^2)^2 + m_h^2 \Gamma_h^2(m_h)}\,.$$

$$\left\langle |\mathcal{M}|^2 \right\rangle_t = \sum_f \frac{N_f \lambda_S^2 m_f^2}{8k^4} \Big[\frac{2k_{\rm cm}^2 - 2m_f^2 + m_h^2}{1 + m_h^2/(4k_{\rm cm}^2)} - \left(m_h^2 - 2m_f^2\right) \log\left(1 + 4k_{\rm cm}^2/m_h^2\right) \Big]\,.$$

Hierarchical Yukawa couplings: strongest coupling to more Boltzmann suppressed quarks/leptons

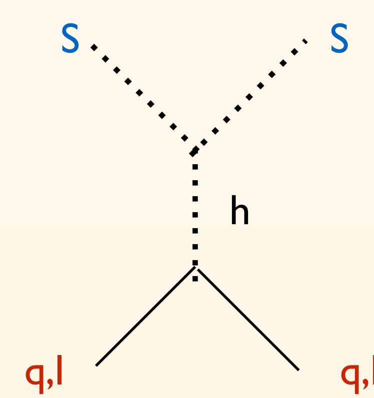
Annihilation processes:

resonant



El. scattering processes:

non-resonant



Freeze-out at few GeV ⟶ what is the abundance of heavy quarks in QCD plasma?

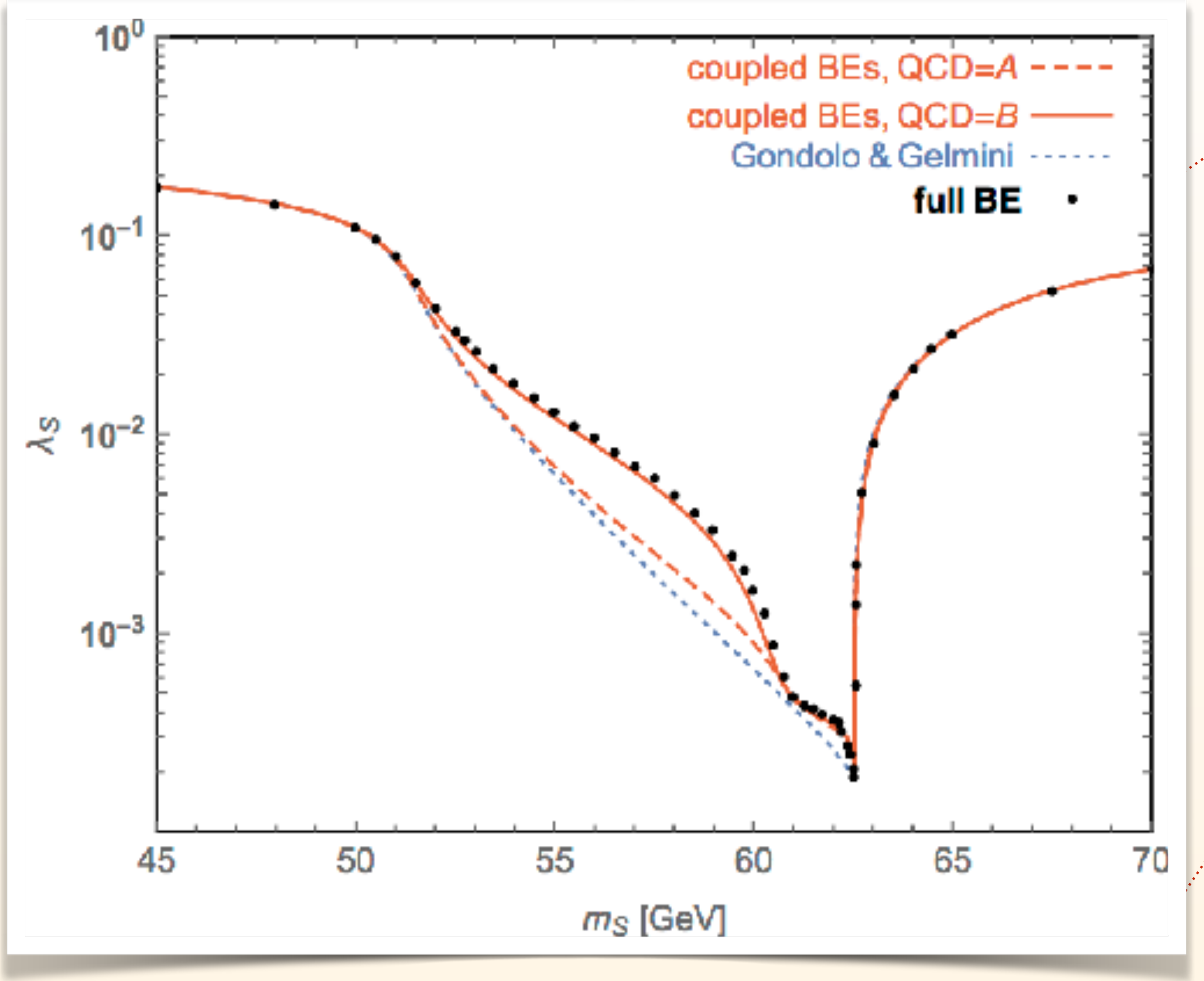
two scenarios:

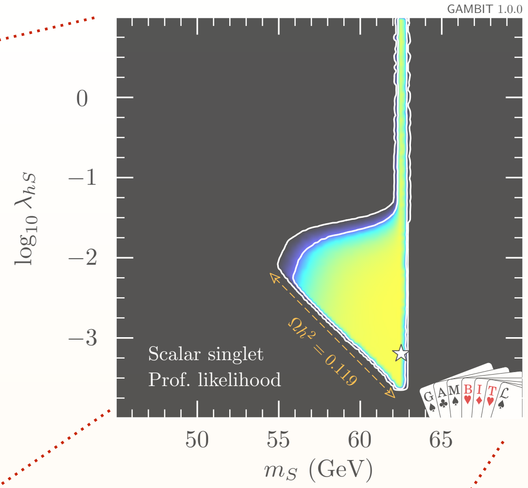
QCD = A - all quarks are free and present in the plasma down to T_c =154 MeV

QCD = B - only light quarks contribute to scattering and only down to $4T_c$

RESULTS

RD CONTOURS





essentially the only region left for this model

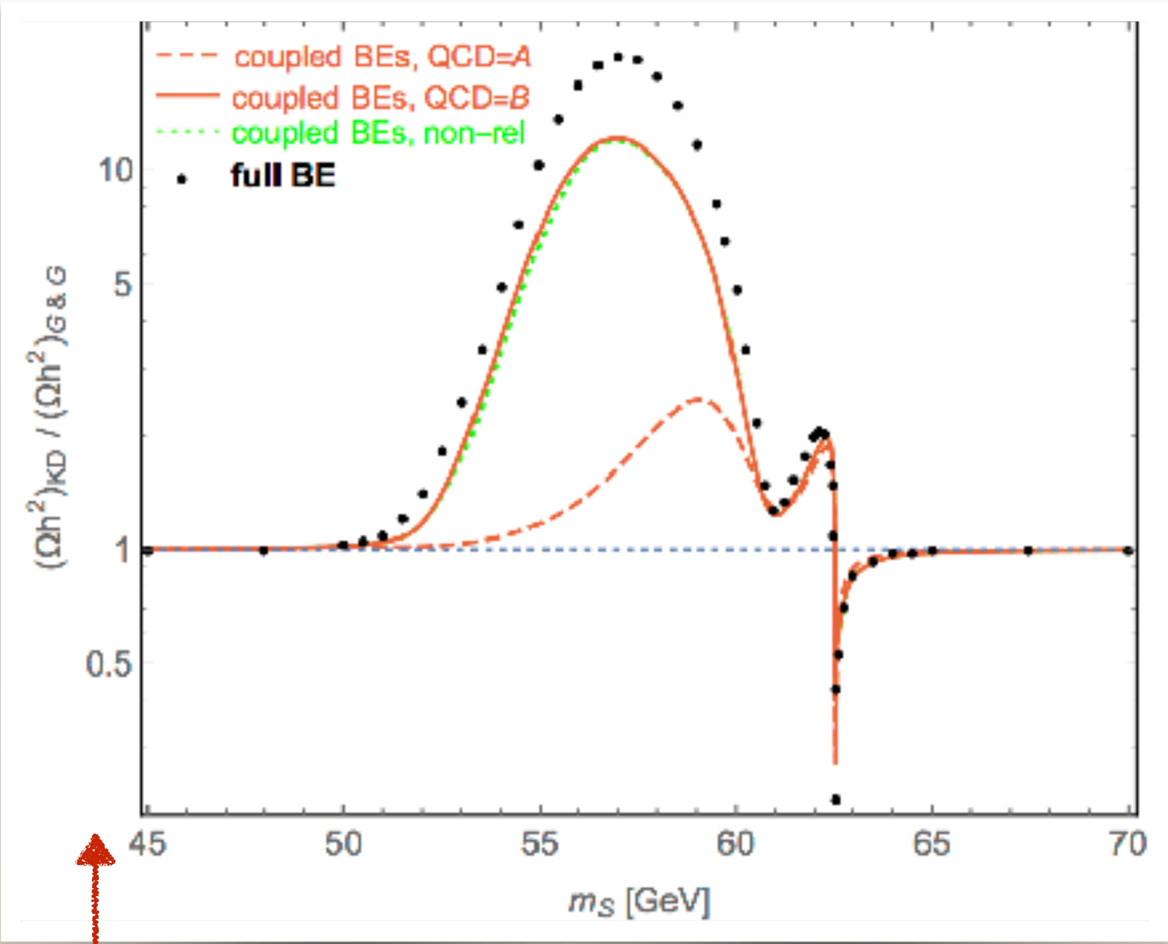
Significant <u>modification</u> of the observed relic density contour in the Scalar Singlet DM model

→ larger coupling needed → better chance for closing the last window

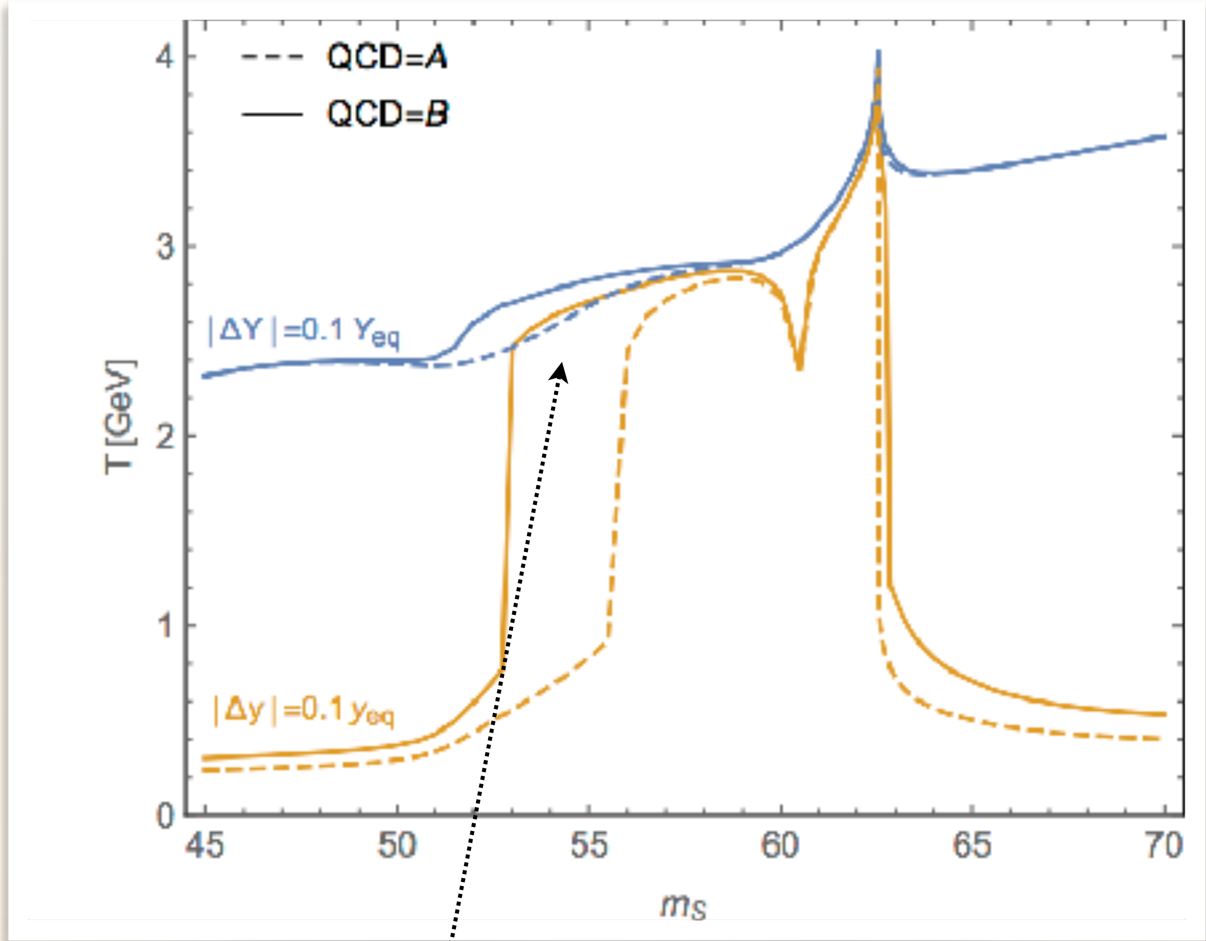
RESULTS

EFFECT

effect on relic density:



kinetic and chemical decoupling:



effect on relic density:
up to O(~10)

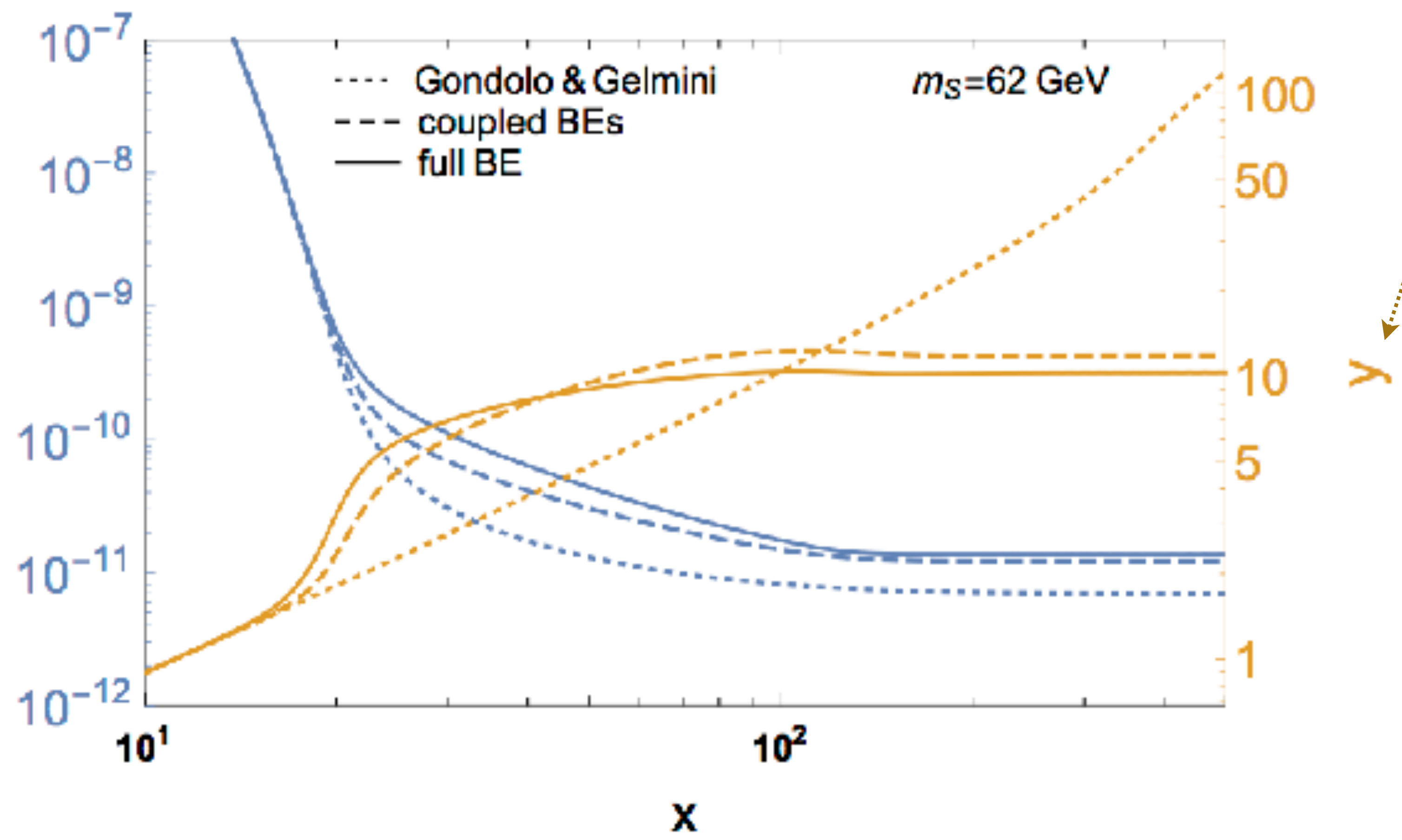
ratio approaches 1,
but does not reach it!

Why such non-trivial shape of the effect of early kinetic decoupling?

→ we'll inspect the y and Y evolution…

DENSITY AND T_{DM} EVOLUTION

for m_{DM} = 62 GeV, i.e. just below the resonance:



DM temperature

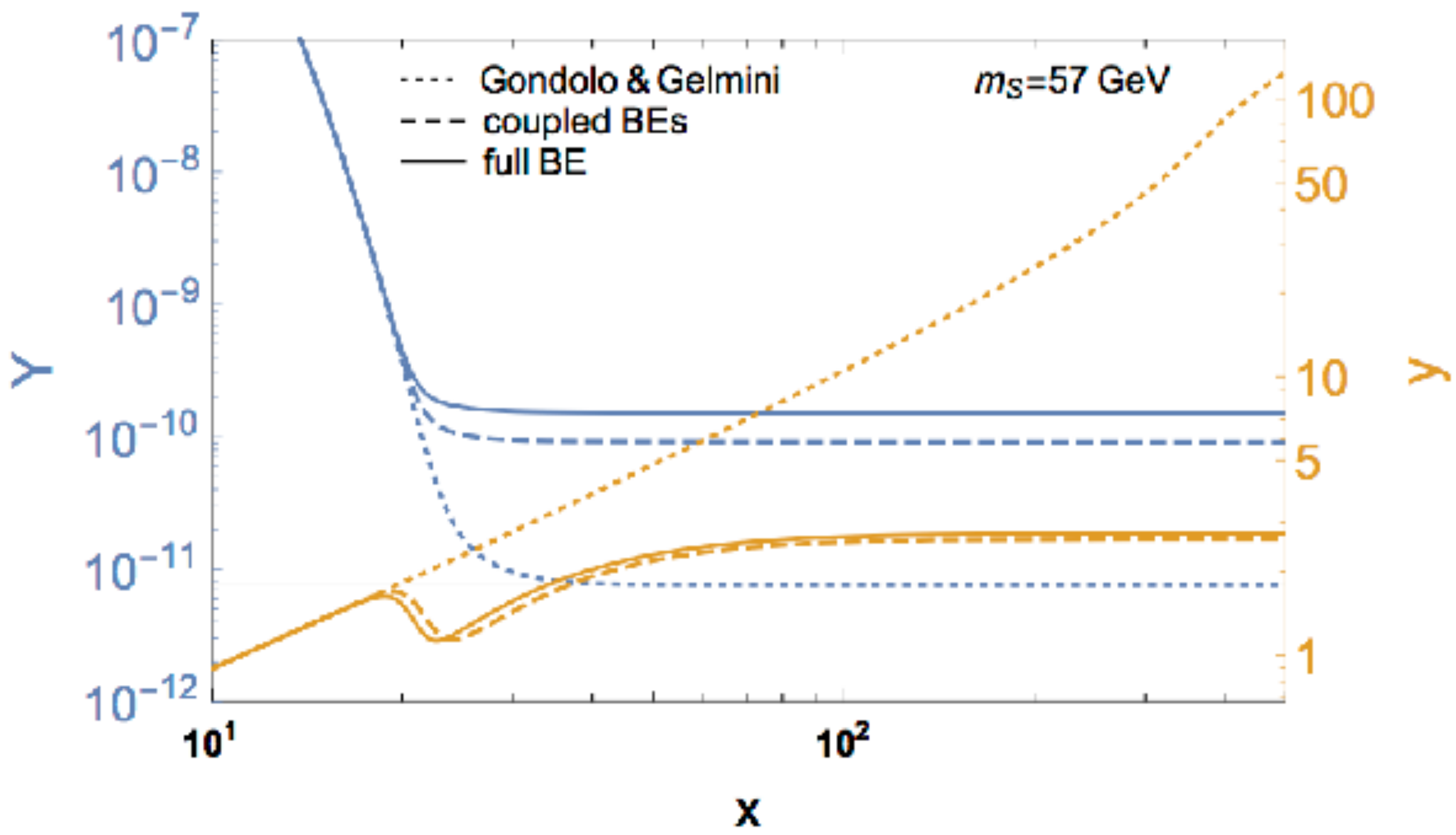
co-moving number density

Resonant annihilation **most effective** for **low momenta**

⟶ DM fluid goes through "heating" phase before leaves kinetic equilibrium

DENSITY AND T_{DM} EVOLUTION

for m_{DM} = 57 GeV, i.e. further away from the resonance:

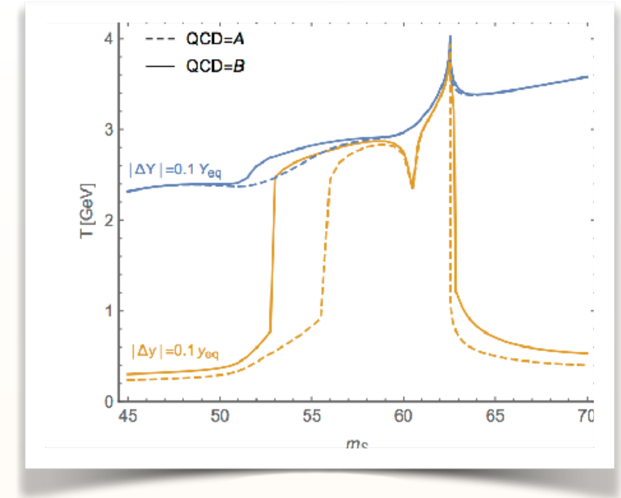


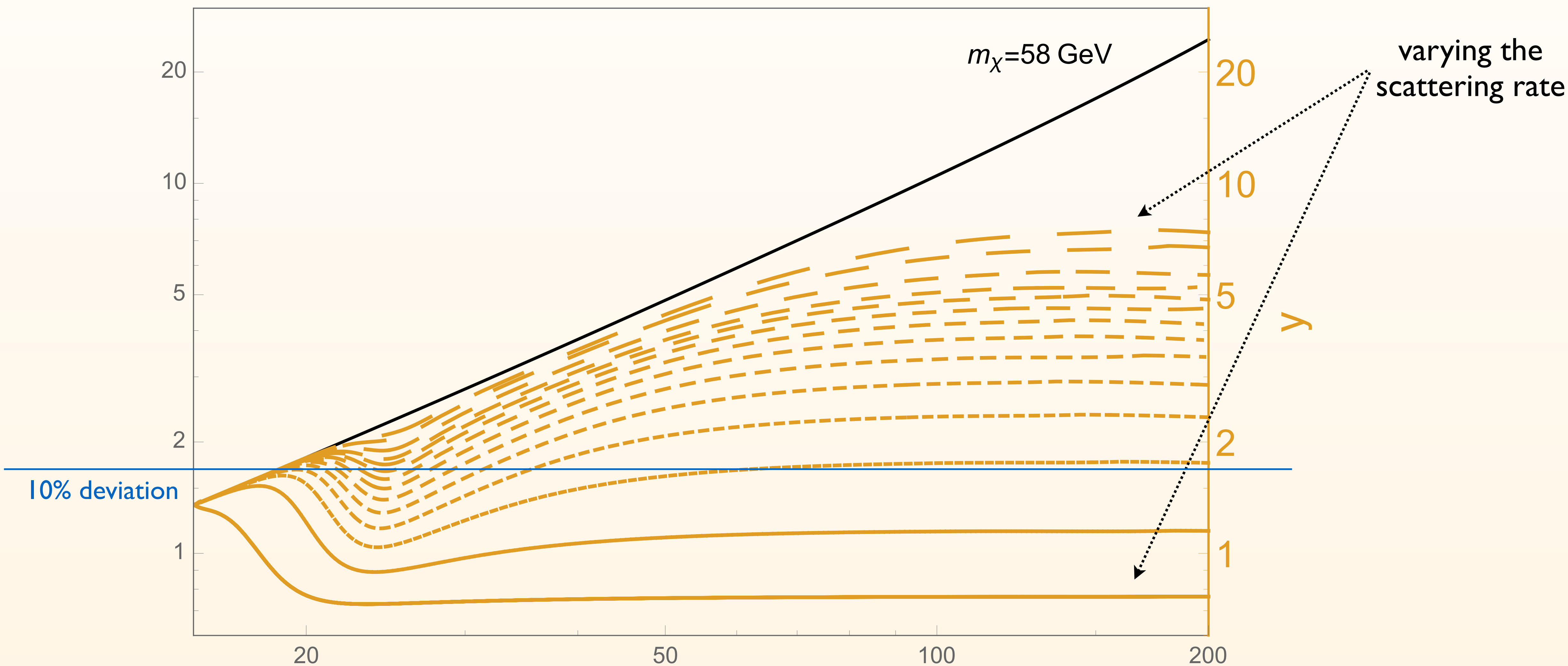
Resonant annihilation **most effective** for **high momenta**

⟶ DM fluid goes through fast "cooling" phase

after that when T_{DM} drops to much annihilation not effective anymore

WHY SPIKES IN T_{KD}?

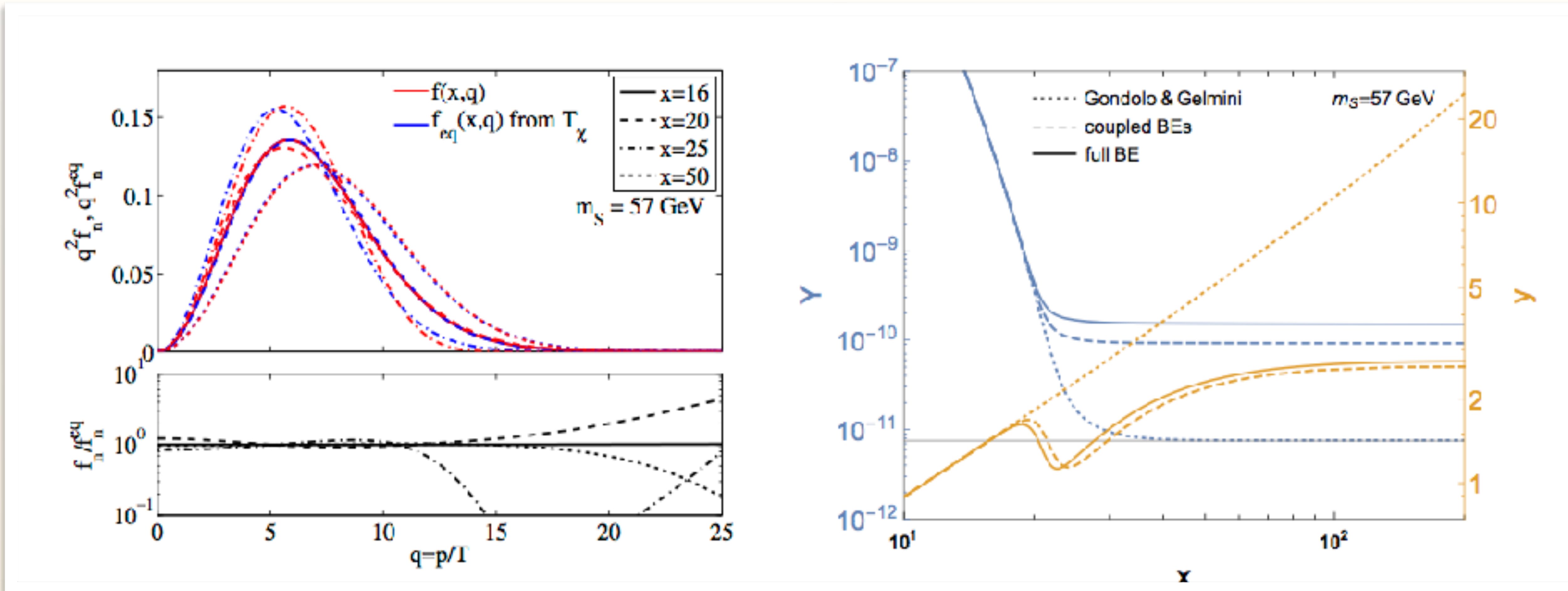




Effect resembling first order „phase transition” —
artificial as dependent on a particular choice of T_{KD} definition

Full phase-space BE solver

Solutions for full phase-space distribution function:



Results of both approaches compatible:
some deviation from equilibrium shape mildly affects the Y and y evolution

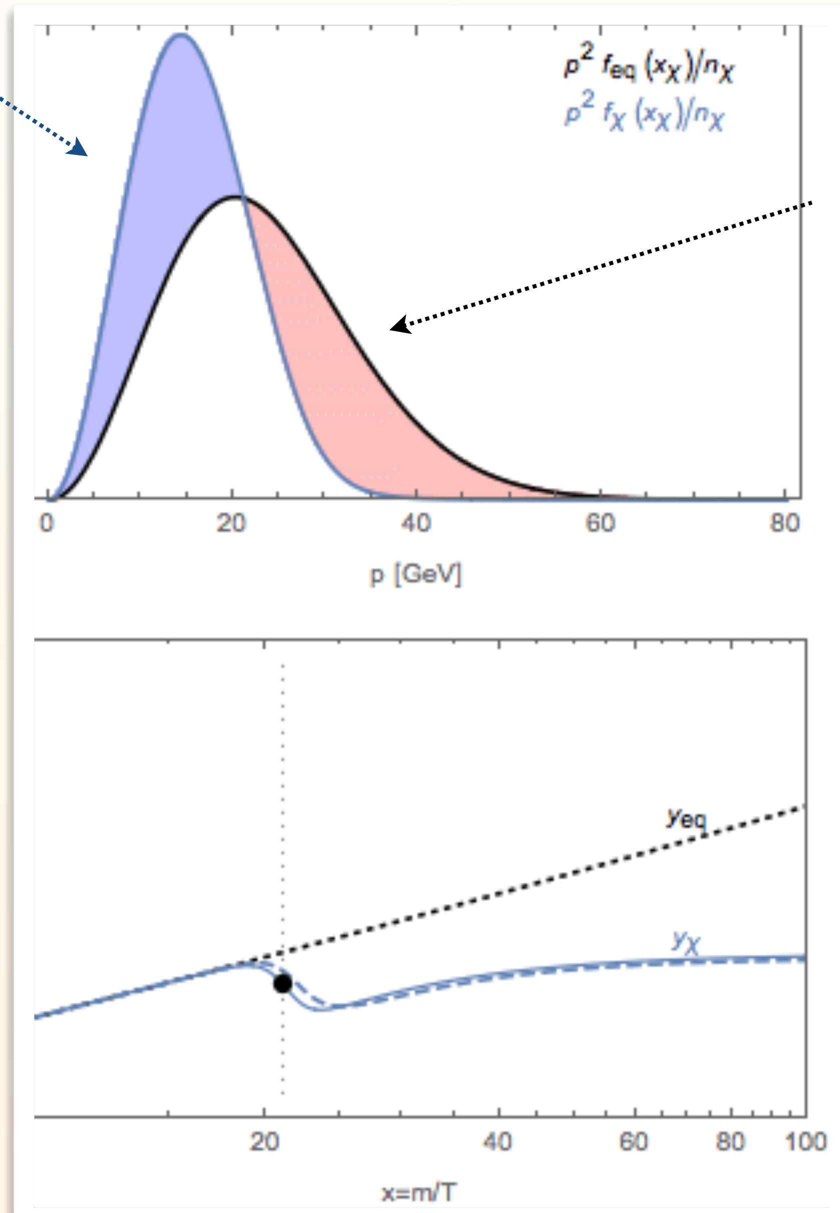
Allows to study the evolution of $f_\chi(p)$ and
the interplay between scatterings and annihilation!

FULL PHASE-SPACE EVOLUTION

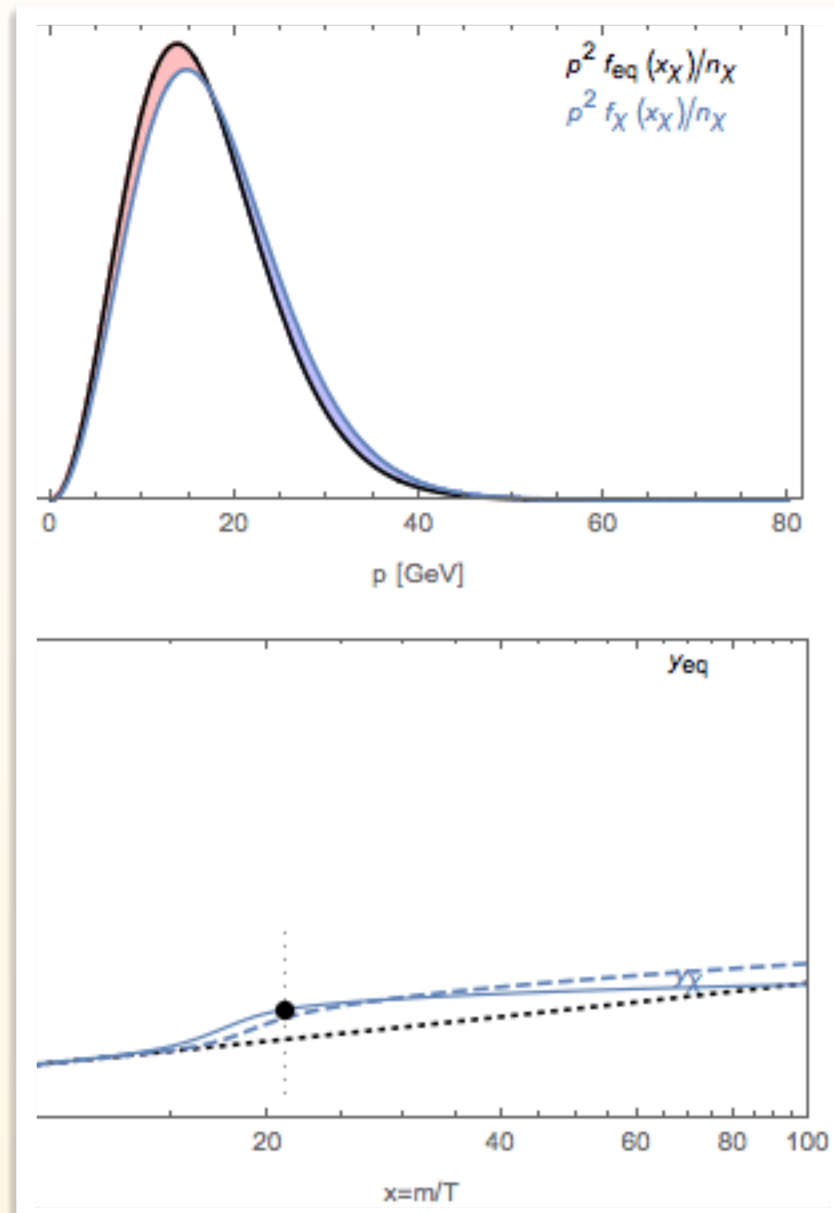
m_{DM} = 58 GeV

m_{DM} = 62.5 GeV

blue - full solution for f_{DM} at T_{DM}



black - equilibrium at T_{DM}



significant deviation from equilibrium shape already around freeze-out

⟶ effect on relic density largest, both from different T and f_{DM}

large deviations at later times, around freeze-out not far from eq. shape

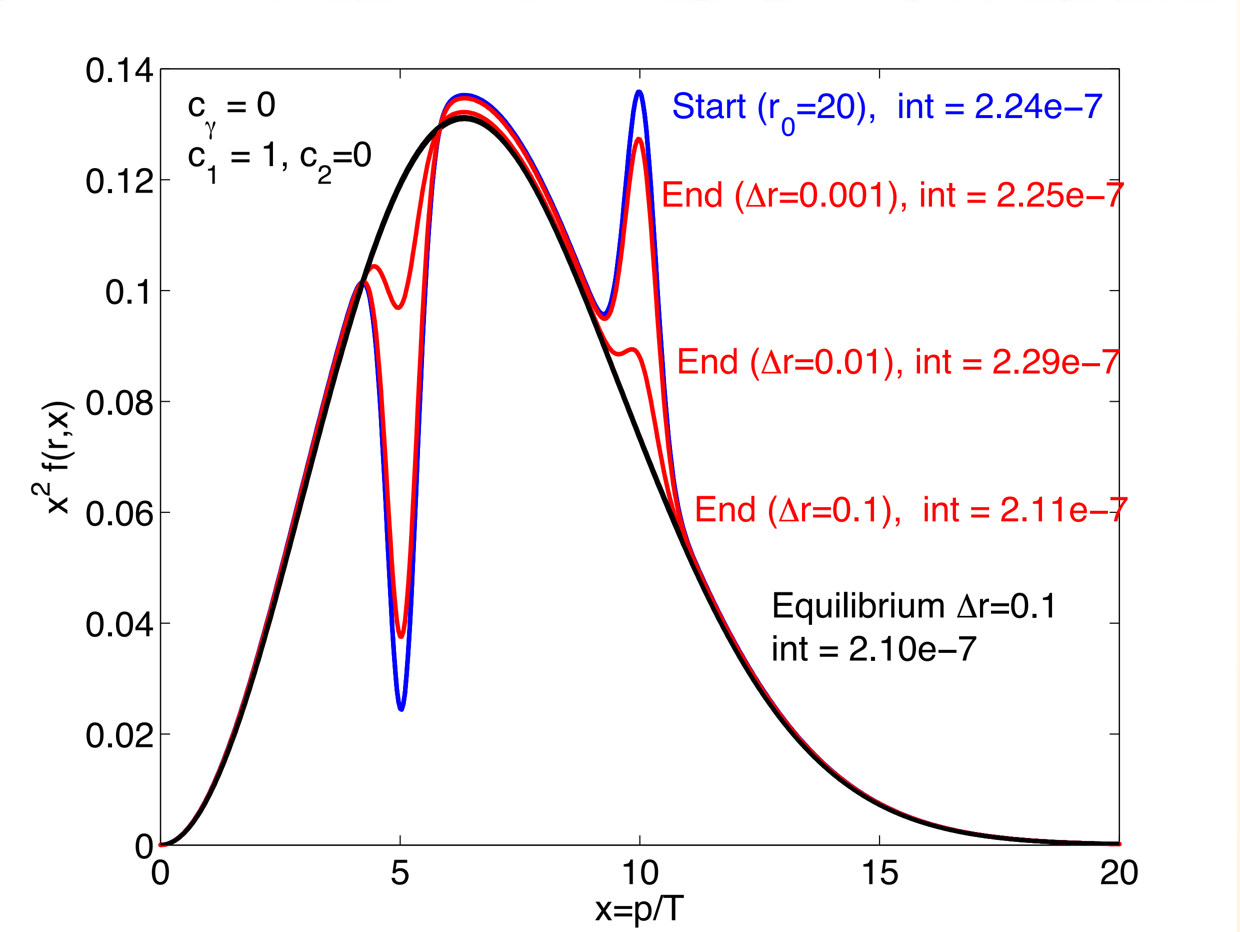
⟶ effect on relic density ~only from different T

KD BEFORE CD?

Obvious issue:
How to define exactly the kinetic and chemical decouplings and what is the significance of such definitions?

⇓

Improved question:
Can kinetic decoupling happen much earlier than chemical?



we have already seen that even if scatterings were very inefficient compared to annihilation, departure from equilibrium for both Y and y happened around the same time…

← turn off scatterings and take s-wave annihilation; look at local disturbance

annihilation/production precesses drive to restore kinetic equilibrium!

CONCLUSIONS

1. One needs to remember that kinetic equilibrium is a necessary assumption for standard relic density calculations

2. Coupled system of Boltzmann equations for 0th and 2nd moments allow for a very accurate treatment of the kinetic decoupling and its effect on relic density

3. In special cases the full phase space Boltzmann equation can be necessary — especially if one wants to trace DM temperature as well

Exception *N*:

kinetic decoupling can happen together with freeze-out…

TAKEAWAY MESSAGE

When computing relic density of dark matter one needs carefully to check if the standard treatment is sufficient for the case at hand

”Everything should be made as simple as possible, but no simpler.”

attributed to* **Albert Einstein**

*The published quote reads:
”It can scarcely be denied that the supreme goal of all theory is to make the irreducible basic elements as simple and as few as possible without having to surrender the adequate representation of a single datum of experience.”
„On the Method of Theoretical Physics" ,The Herbert Spencer Lecture, delivered at Oxford (10 June 1933); also published in *Philosophy of Science*, Vol. 1, No. 2 (April 1934), pp. 163-169., p. 165

God Jul