

DARK MATTER PRODUCTION OUT OF KINETIC EQUILIBRIUM: LATEST DEVELOPMENTS

Andrzej Hryczuk



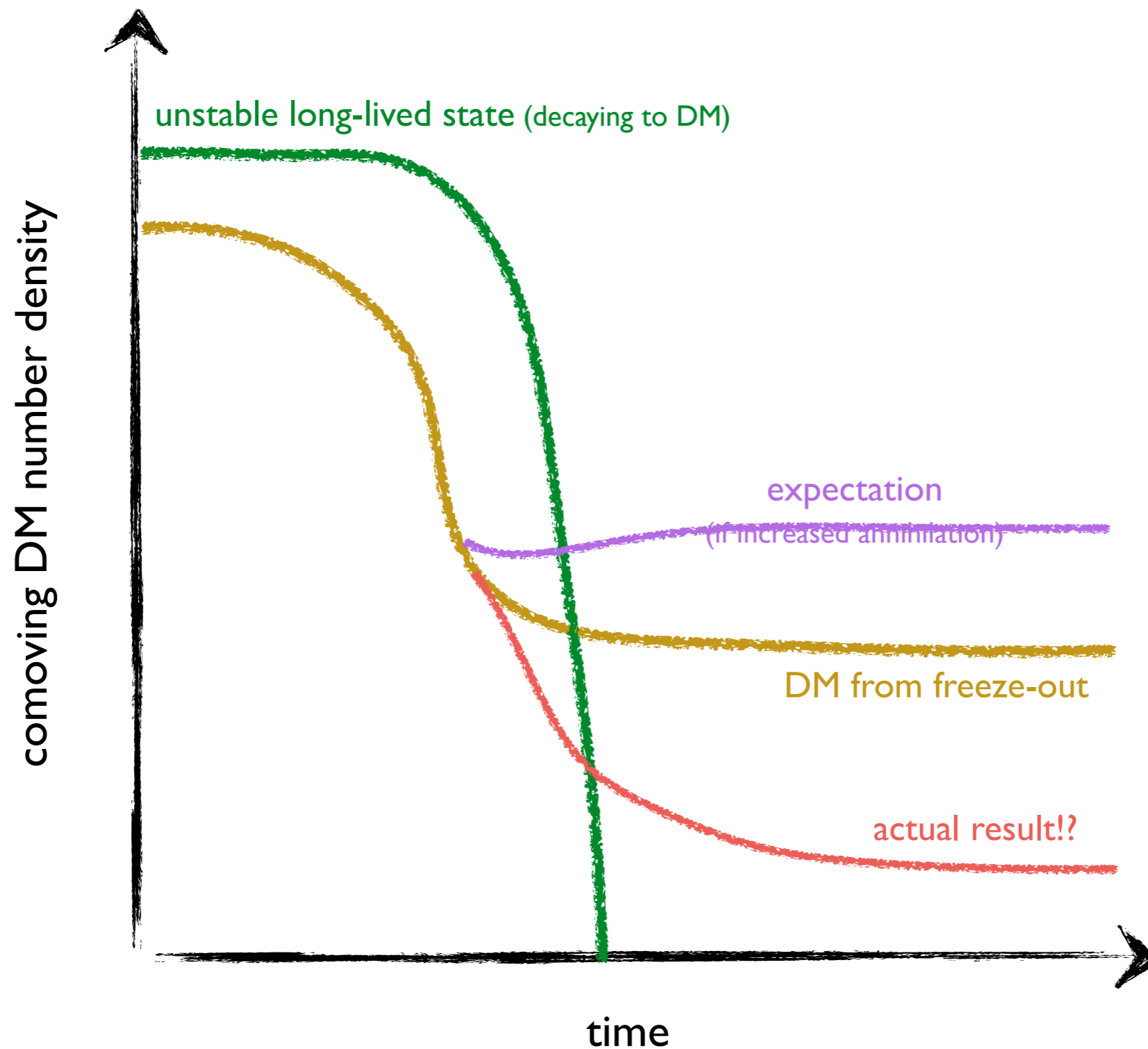
Based on:

T. Binder, T. Bringmann, M. Gustafsson & A.H. [1706.07433](#), [2103.01944](#)

A.H. & M. Laletin [2204.07078](#), [2104.05684](#)

work in progress with **S. Chatterjee**

IN CASE YOU'RE NOT INTERESTED IN WHAT FOLLOWS...



TO SEE WHY AND LEARN MORE STAY TUNED :)

THERMAL RELIC DENSITY

STANDARD SCENARIO

modified expansion rate

e.g., relentless DM, D'Eramo et al. '17, ...

numerical codes e.g.,
DarkSUSY, micrOMEGAs, MadDM, SuperISOrelic, ...

$$\frac{dn_\chi}{dt} + 3Hn_\chi = - \langle \sigma_{\chi\bar{\chi} \rightarrow ij} \sigma_{\text{rel}} \rangle^{\text{eq}} (n_\chi n_{\bar{\chi}} - n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}})$$

general multi-component dark sector

$$\frac{dn_\chi}{dt} + 3Hn_\chi$$

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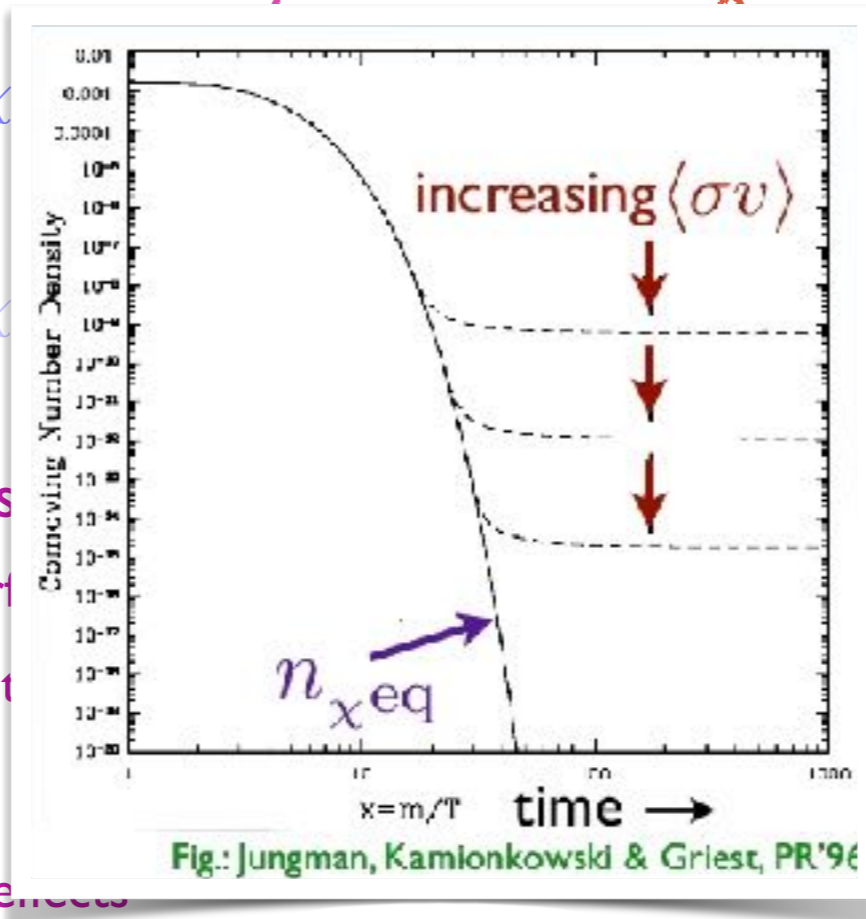
modified cross

Sommerfeld

Bound States

NLO

finite T effects



$$n_{\bar{\chi}} - n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}}$$

$$n_{\bar{\chi}} - n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}}$$

breakdown of necessary assumptions leading to different form of the equation, e.g. violation of kinetic equilibrium

where the thermally averaged cross section:

$$\langle \sigma_{\chi\bar{\chi} \rightarrow ij} v_{\text{rel}} \rangle^{\text{eq}} = \frac{h_\chi^2}{n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}}} \int \frac{d^3 \vec{p}_\chi}{(2\pi)^3} \frac{d^3 \vec{p}_{\bar{\chi}}}{(2\pi)^3} \sigma_{\chi\bar{\chi} \rightarrow ij} v_{\text{rel}} f_\chi^{\text{eq}} f_{\bar{\chi}}^{\text{eq}}$$

THERMAL RELIC DENSITY

STANDARD APPROACH

Boltzmann equation for $f_\chi(p)$:

$$E (\partial_t - H\vec{p} \cdot \nabla_{\vec{p}}) f_\chi = \mathcal{C}[f_\chi]$$

integrate over p
(i.e. take 0th moment)

$$\frac{dn_\chi}{dt} + 3Hn_\chi = - \langle \sigma_{\chi\bar{\chi} \rightarrow ij} \sigma_{\text{rel}} \rangle^{\text{eq}} (n_\chi n_{\bar{\chi}} - n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}})$$

where the thermally averaged cross section:

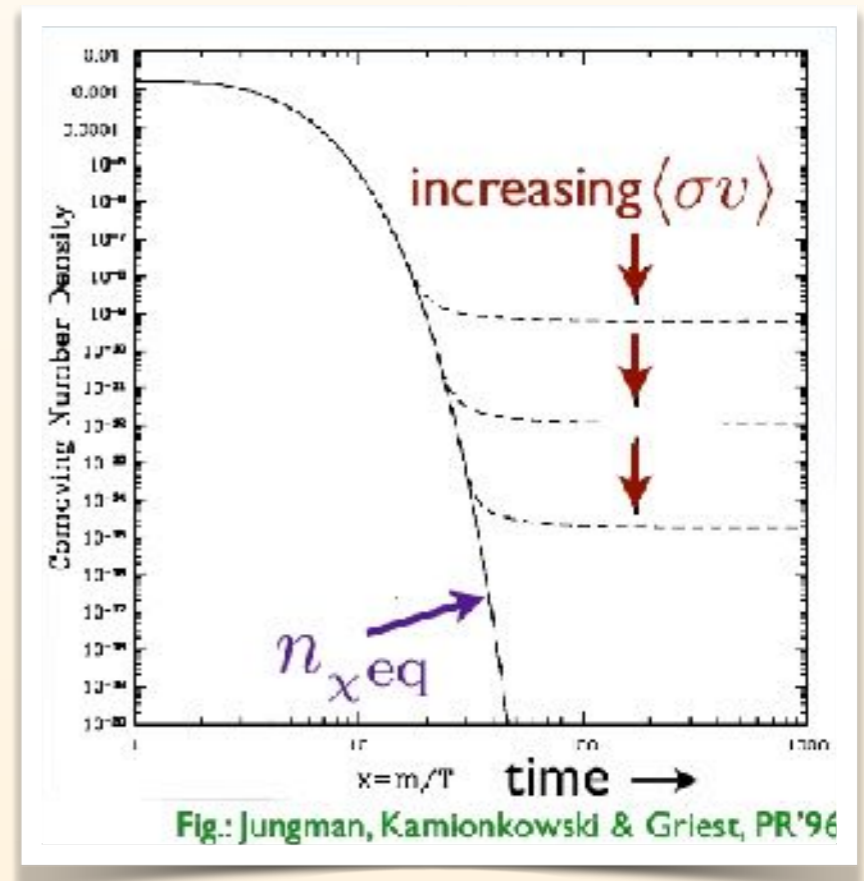
$$\langle \sigma_{\chi\bar{\chi} \rightarrow ij} v_{\text{rel}} \rangle^{\text{eq}} = - \frac{h_\chi^2}{n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}}} \int \frac{d^3\vec{p}_\chi}{(2\pi)^3} \frac{d^3\vec{p}_{\bar{\chi}}}{(2\pi)^3} \sigma_{\chi\bar{\chi} \rightarrow ij} v_{\text{rel}} f_\chi^{\text{eq}} f_{\bar{\chi}}^{\text{eq}}$$

Critical assumption:
kinetic equilibrium at chemical decoupling

$$f_\chi \sim a(T) f_\chi^{\text{eq}}$$

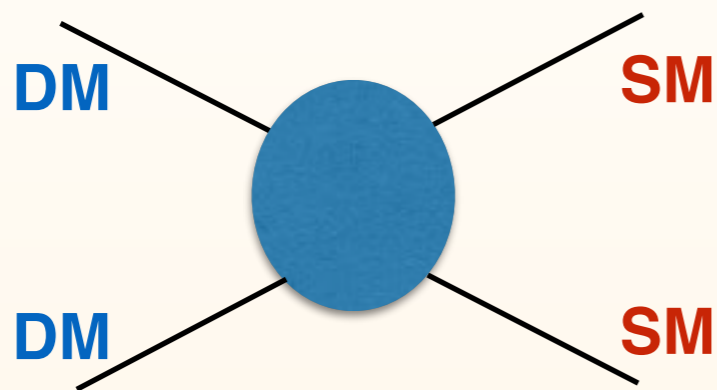
*assumptions for using Boltzmann eq:
classical limit, molecular chaos,...

...for derivation from thermal QFT
see e.g., 1409.3049



FREEZE-OUT vs. DECOUPLING

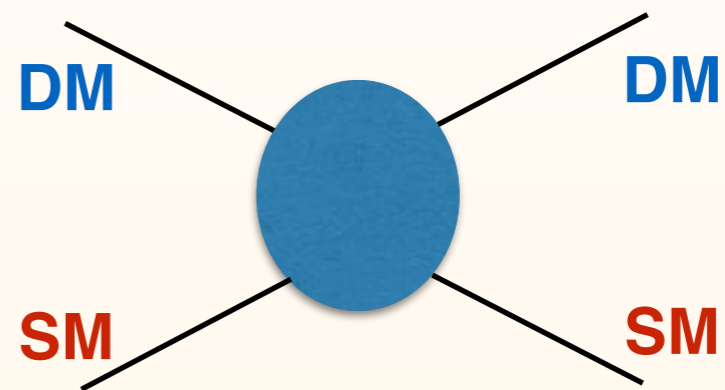
annihilation



$$\sum_{\text{spins}} |\mathcal{M}^{\text{pair}}|^2 = F(p_1, p_2, p'_1, p'_2)$$

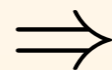
crossing sym.
 \longleftrightarrow

(elastic) scattering



$$\sum_{\text{spins}} |\mathcal{M}^{\text{scatt}}|^2 = F(k, -k', p', -p)$$

Boltzmann suppression of **DM** vs. **SM**



scatterings typically more frequent

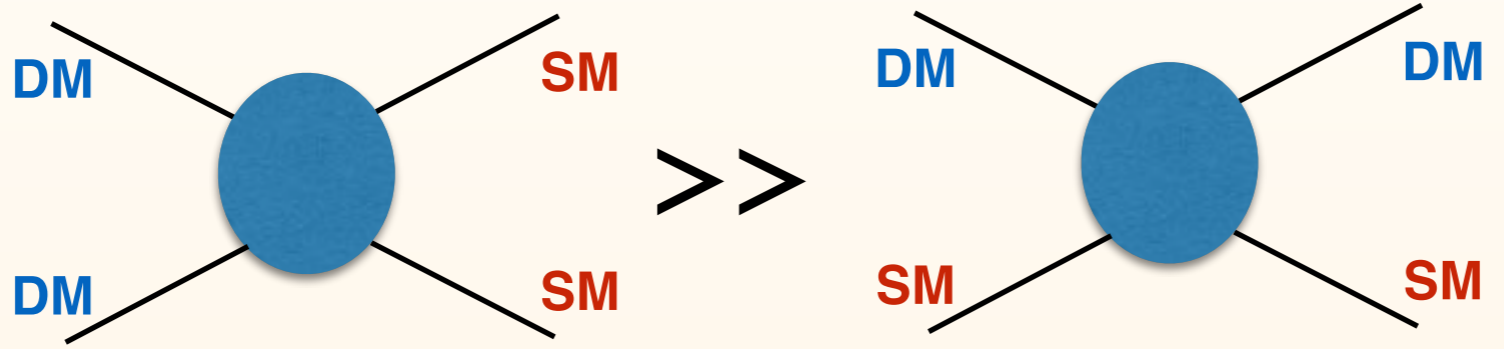
dark matter frozen-out but typically
 still kinetically coupled to the plasma

Schmid, Schwarz, Widern '99; Green, Hofmann, Schwarz '05

EARLY KINETIC DECOUPLING?

A **necessary** and **sufficient** condition: scatterings weaker than annihilation
i.e. rates around freeze-out: $H \sim \Gamma_{\text{ann}} \gtrsim \Gamma_{\text{el}}$

Possibilities:

- A)  e.g., resonant annihilation
- B) Boltzmann suppression of **SM** as strong as for **DM**
e.g., below threshold annihilation (forbidden-like DM)
- C) Scatterings and annihilation have different structure
e.g., semi-annihilation, 3 to 2 models,...
- D) Multi-component dark sectors
e.g., additional sources of DM from late decays, ...

HOW TO GO BEYOND KINETIC EQUILIBRIUM?

All information is in the full BE:

both about chemical ("normalization") and kinetic ("shape") equilibrium/decoupling

$$E (\partial_t - H\vec{p} \cdot \nabla_{\vec{p}}) f_{\chi} = \mathcal{C}[f_{\chi}]$$

contains both scatterings and annihilations

Two possible approaches:

fBE

solve numerically
for full $f_{\chi}(p)$

have insight on the distribution
no constraining assumptions

numerically challenging
often an overkill

CBE

consider system of equations
for moments of $f_{\chi}(p)$

partially analytic/much easier numerically
manifestly captures all of the relevant physics

finite range of validity
no insight on the distribution

0-th moment: n_{χ}
2-nd moment: T_{χ}

...

NEW TOOL!

GOING BEYOND THE STANDARD APPROACH

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Dark matter Relic Abundance beyond Kinetic Equilibrium

Authors: Tobias Binder, Torsten Bringmann, Michael Gustafsson and Andrzej Hryczuk

DRAKE is a numerical precision tool for predicting the dark matter relic abundance also in situations where the standard assumption of kinetic equilibrium during the freeze-out process may not be satisfied. The code comes with a set of three dedicated Boltzmann equation solvers that implement, respectively, the traditionally adopted equation for the dark matter number density, fluid-like equations that couple the evolution of number density and velocity dispersion, and a full numerical evolution of the phase-space distribution. The code is written in Wolfram Language and includes a Mathematica notebook example program, a template script for terminal usage with the free Wolfram Engine, as well as several concrete example models. DRAKE is a free software licensed under GPL3.

If you use DRAKE for your scientific publications, please cite

- **DRAKE: Dark matter Relic Abundance beyond Kinetic Equilibrium**, Tobias Binder, Torsten Bringmann, Michael Gustafsson and Andrzej Hryczuk, [[arXiv:2103.01944](#)]

Currently, an user guide can be found in the Appendix A of this reference. Please cite also quoted other works applying for specific cases.

v1.0 « [Click here to download DRAKE](#)

(March 3, 2021)

<https://drake.hepforge.org>

Applications:

DM relic density for
any (user defined) model*

Interplay between chemical and
kinetic decoupling

Prediction for the DM
phase space distribution

Late kinetic decoupling
and impact on cosmology

see e.g., [1202.5456](#)

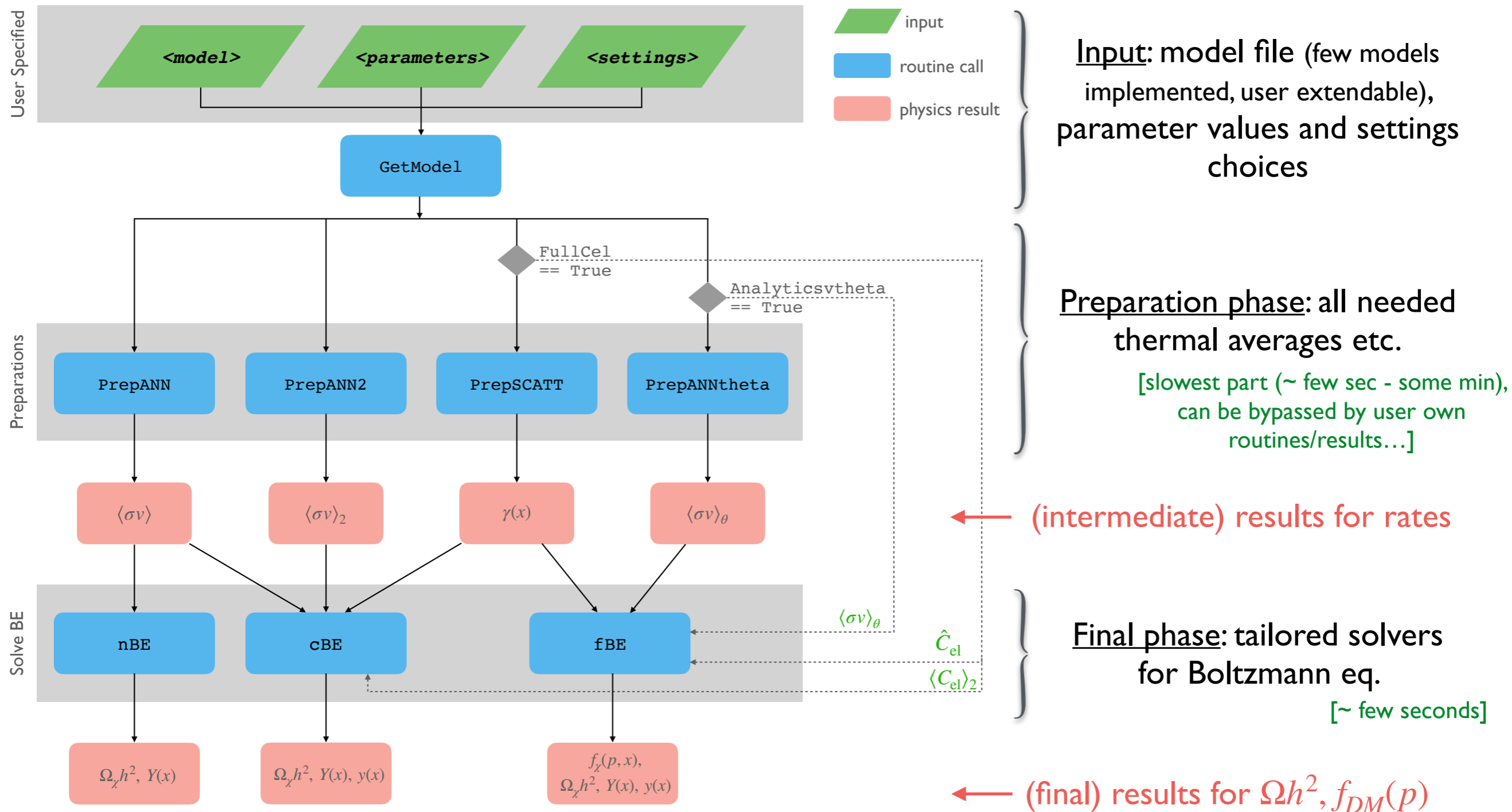
...

(only) prerequisite:
Wolfram Language (or Mathematica)

*at the moment for a single DM species and w/o
co-annihilations... but stay tuned for extensions!

FEW WORDS ABOUT THE CODE

written in *Wolfram Language*, lightweight, modular and simple to use both via script and front end usage

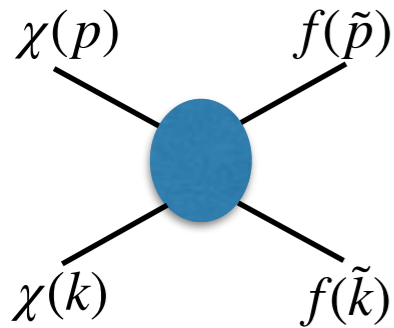


COLLISION TERM

$$E (\partial_t - H\vec{p} \cdot \nabla_{\vec{p}}) f_\chi = \mathcal{C}[f_\chi]$$

contains both scatterings and annihilations

Annihilation:



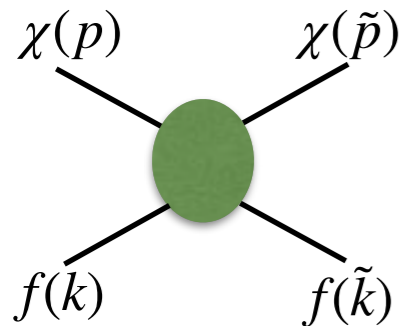
$$C_{\text{ann}} \sim \int d\tilde{\Pi} |\mathcal{M}|_{\chi\chi \leftrightarrow f\tilde{f}}^2 \left(\underbrace{f_f^{\text{eq}}(\tilde{p})}_{\text{easy}} \underbrace{f_f^{\text{eq}}(\tilde{k})}_{\text{easy}} - f_\chi(p) f_\chi(k) \right)$$

easy: no unknown f_χ under integral
 \Rightarrow 1D integration

medium: no unknown f_χ under integral
 \Rightarrow 2-3D integration

hard: unknown f_χ under integral
 \Rightarrow 2-4D integration

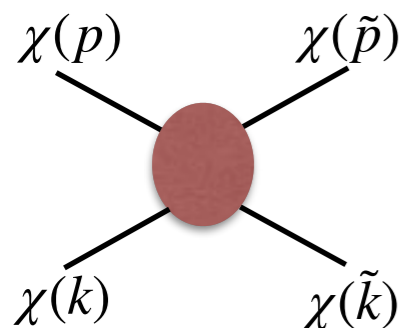
El. scattering (on SM particles):



$$C_{\text{el}} \sim \int d\tilde{\Pi} |\mathcal{M}|_{\chi f \leftrightarrow \chi f}^2 \left(\underbrace{f_\chi(\tilde{p}) f_f^{\text{eq}}(\tilde{k}) (1 \pm f_f^{\text{eq}}(k))}_{\text{hard}} - \underbrace{f_\chi(p) f_f^{\text{eq}}(k) (1 \pm f_f^{\text{eq}}(\tilde{k}))}_{\text{medium}} \right)$$

An approximate method needed!

El. self-scattering (DM on DM):



$$C_{\text{self}} \sim \int d\tilde{\Pi} |\mathcal{M}|_{\chi\chi \leftrightarrow \chi\chi}^2 \left(\underbrace{f_\chi(\tilde{p}) f_\chi(\tilde{k})}_{\text{hard}} - f_\chi(p) f_\chi(k) \right)$$

$$d\tilde{\Pi} = d\Pi_{\tilde{p}} d\Pi_k d\Pi_{\tilde{k}} \delta^{(4)}(\tilde{p} + p - \tilde{k} - k)$$

APPROACHES

I) Expand in „small momentum transfer”

$$M_{\text{DM}} \gg |\vec{q}| \sim T \gg m_{\text{SM}}$$

Bringmann, Hofmann '06

$$\delta^{(3)}(\vec{p} + \vec{k} - \vec{p} - \vec{k}) \approx \sum_n \frac{1}{n!} (\vec{q} \cdot \nabla_{\vec{p}})^n \delta^{(3)}(\vec{p} - \vec{p})$$

typical momentum transfer

Kasahara '09; Binder, Covi, Kamada, Murayama, Takahashi, Yoshida '16

$$f_3 \simeq f_1 + \tilde{\mathbf{q}}_i \frac{\partial f_1}{\partial \mathbf{p}_{1i}} + \frac{1}{2} \tilde{\mathbf{q}}_i \tilde{\mathbf{q}}_j \frac{\partial^2 f_1}{\partial \mathbf{p}_{1i} \partial \mathbf{p}_{1j}}$$

A.H. & S. Chatterjee, work in progress...

(on different expansion schemes)

⇒ all lead to Fokker-Planck type eq.

II) Replace the backward term with a simpler one (i.e. a **relaxation-like approximation**)

⇒ simpler, but generally incorrect

Ala-Mattinen, Kainulainen '19

Ala-Mattinen, Heikinheimo, Kainulainen, Tuominen '22

$$\begin{aligned} \hat{C}_{\text{E},m}(p_1, t) &\rightarrow -\delta f(p_1, t) \Gamma_{\text{E}}^m(p_1, t) \\ &= (g_m(t) f_{\text{eq}}(p_1, t) - f(p_1, t)) \Gamma_{\text{E}}^m(p_1, t) \end{aligned}$$

III) Langevin simulations

Kim, Laine '23

$$(\hat{p}^i)' = -\hat{\eta} \hat{p}^i + \hat{f}^i, \quad \langle \hat{f}^i(x_1) \hat{f}^j(x_2) \rangle = \hat{\zeta} \delta^{ij} \delta(x_1 - x_2)$$

⇒ perhaps promising...

stochastic term, taking care of detailed balance

IV) Fully numerical implementation

A.H. & M. Laletin [2204.07078](#) (focus on DM self-scatterings)

Ala-Mattinen, Heikinheimo, Kainulainen, Tuominen '22

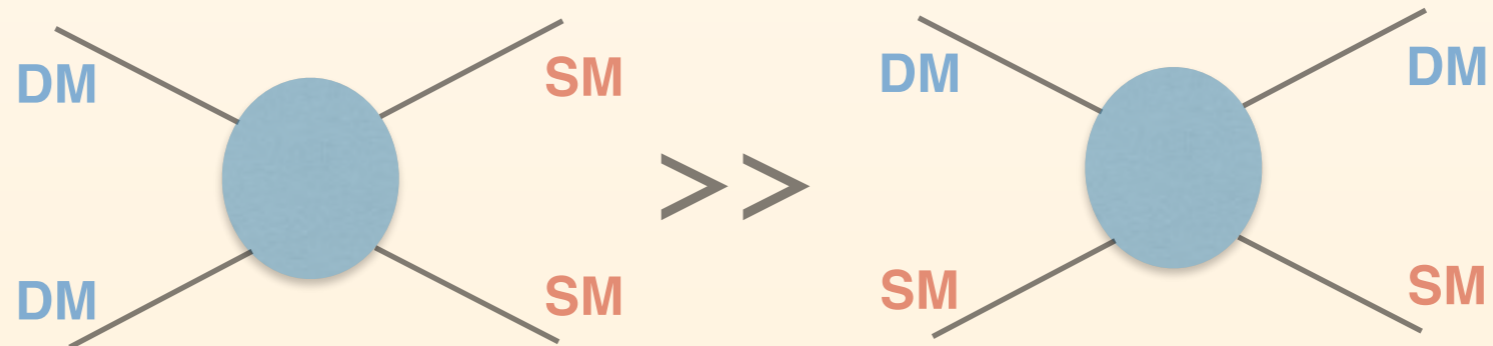
Du, Huang, Li, Li, Yu '21

Aboubrahim, Klasen, Wiggering '23

⇒ doable, but (very) CPU expensive

EXAMPLE A: SCALAR SINGLET DM

A)



EXAMPLE A

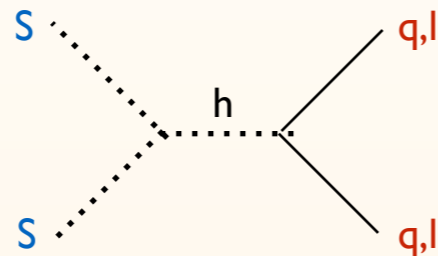
SCALAR SINGLET DM

To the SM Lagrangian add one singlet scalar field S with interactions with the Higgs:

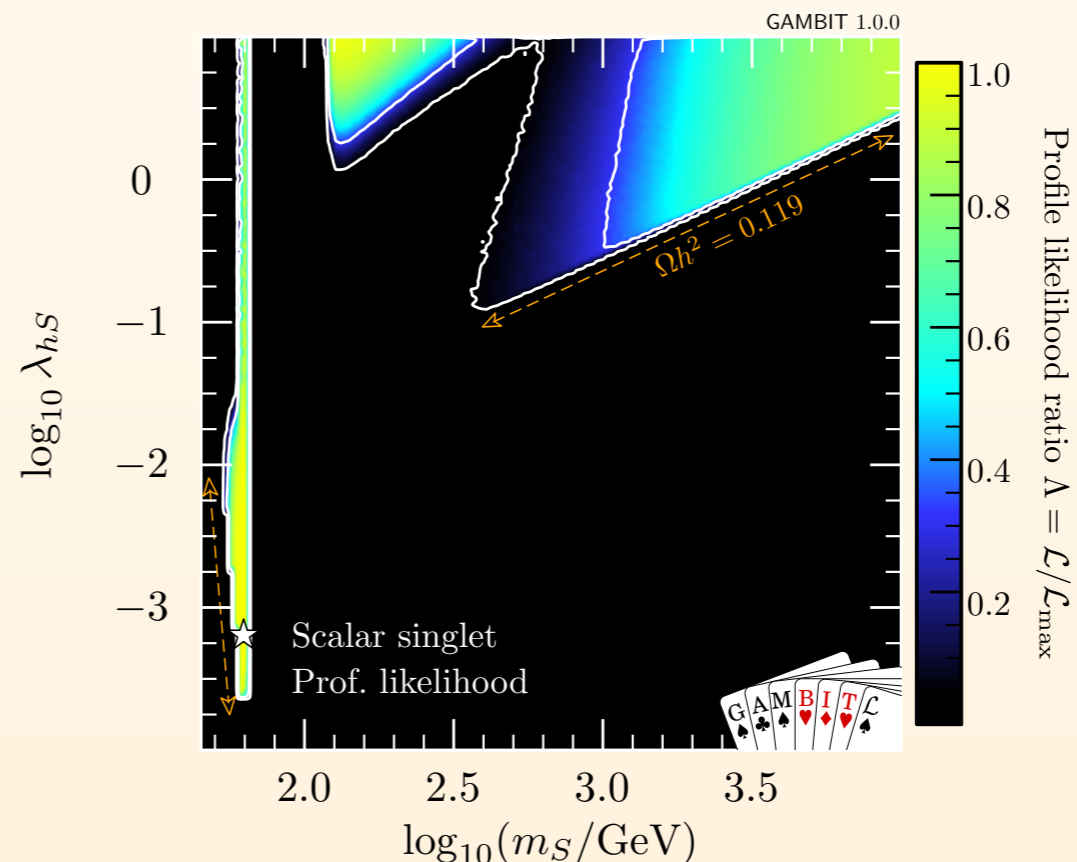
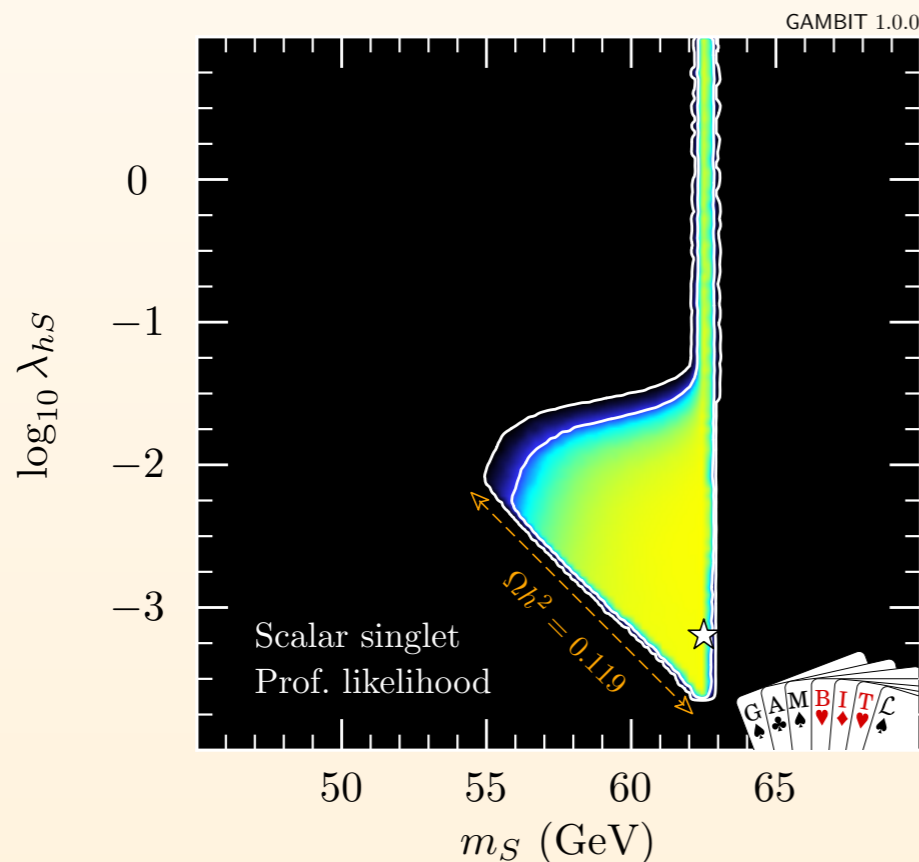
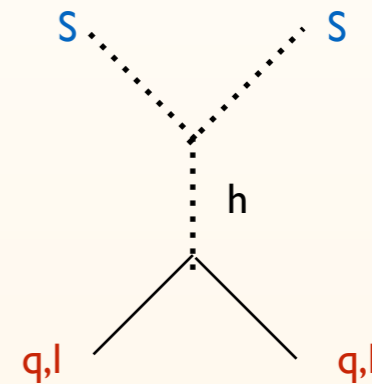
$$\mathcal{L}_S = \frac{1}{2} \partial_\mu S \partial^\mu S - \frac{1}{2} \mu_S^2 S^2 - \frac{1}{2} \lambda_s S^2 |H|^2$$

$$m_s = \sqrt{\mu_S^2 + \frac{1}{2} \lambda_s v_0^2}$$

Annihilation
processes:
resonant



El. scattering
processes:
non-resonant

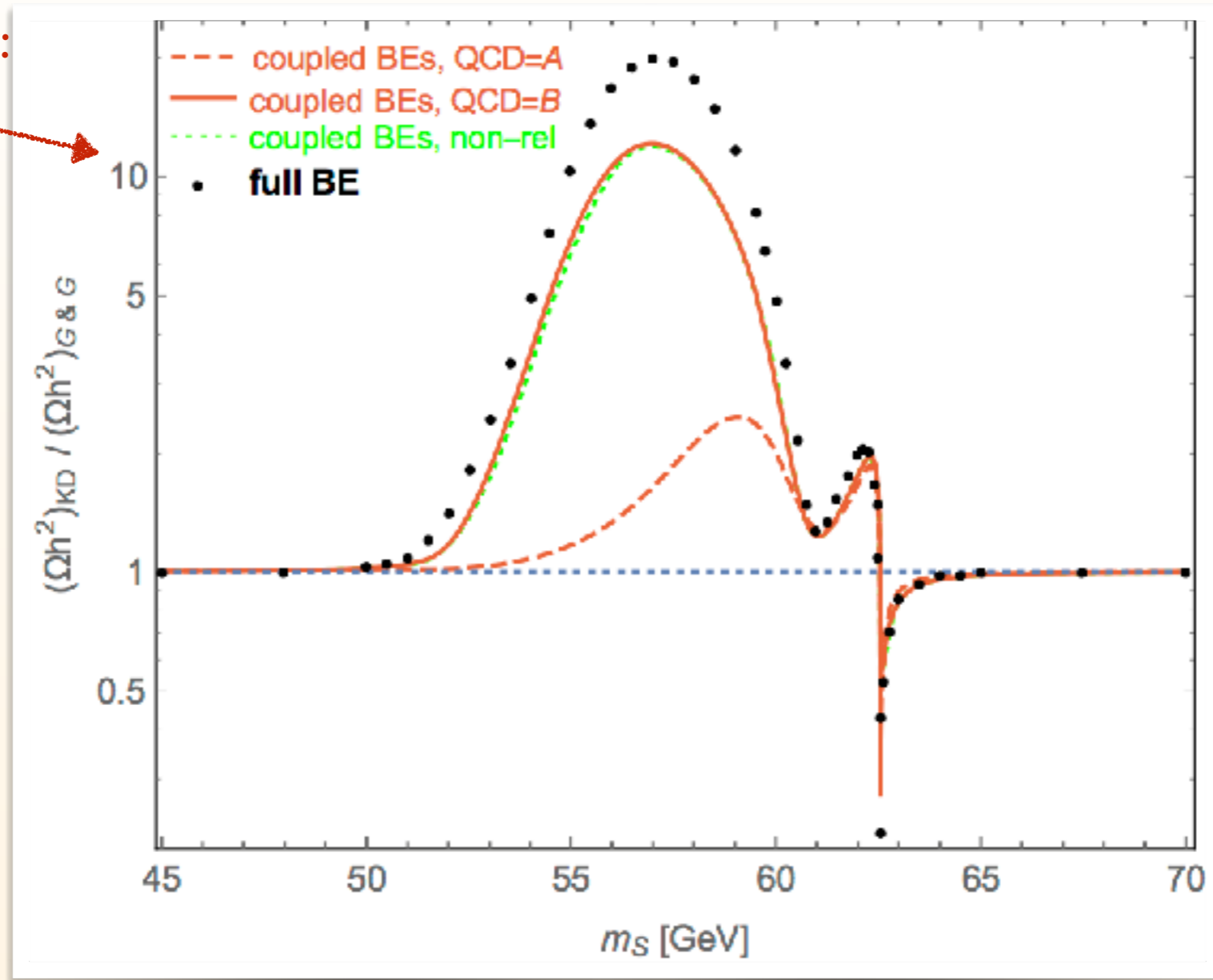


GAMBIT collaboration
1705.07931

RESULTS

EFFECT ON THE Ωh^2

effect on relic density:
up to $O(\sim 10)$



[... Freeze-out at few GeV \rightarrow what is the abundance of heavy quarks in QCD plasma?

two scenarios: QCD = A - all quarks are free and present in the plasma down to $T_c = 154$ MeV
 QCD = B - only light quarks contribute to scattering and only down to $4T_c$...] 14

EXAMPLE D: WHEN ADDITIONAL INFLUX OF DM ARRIVES

D) Multi-component dark sectors

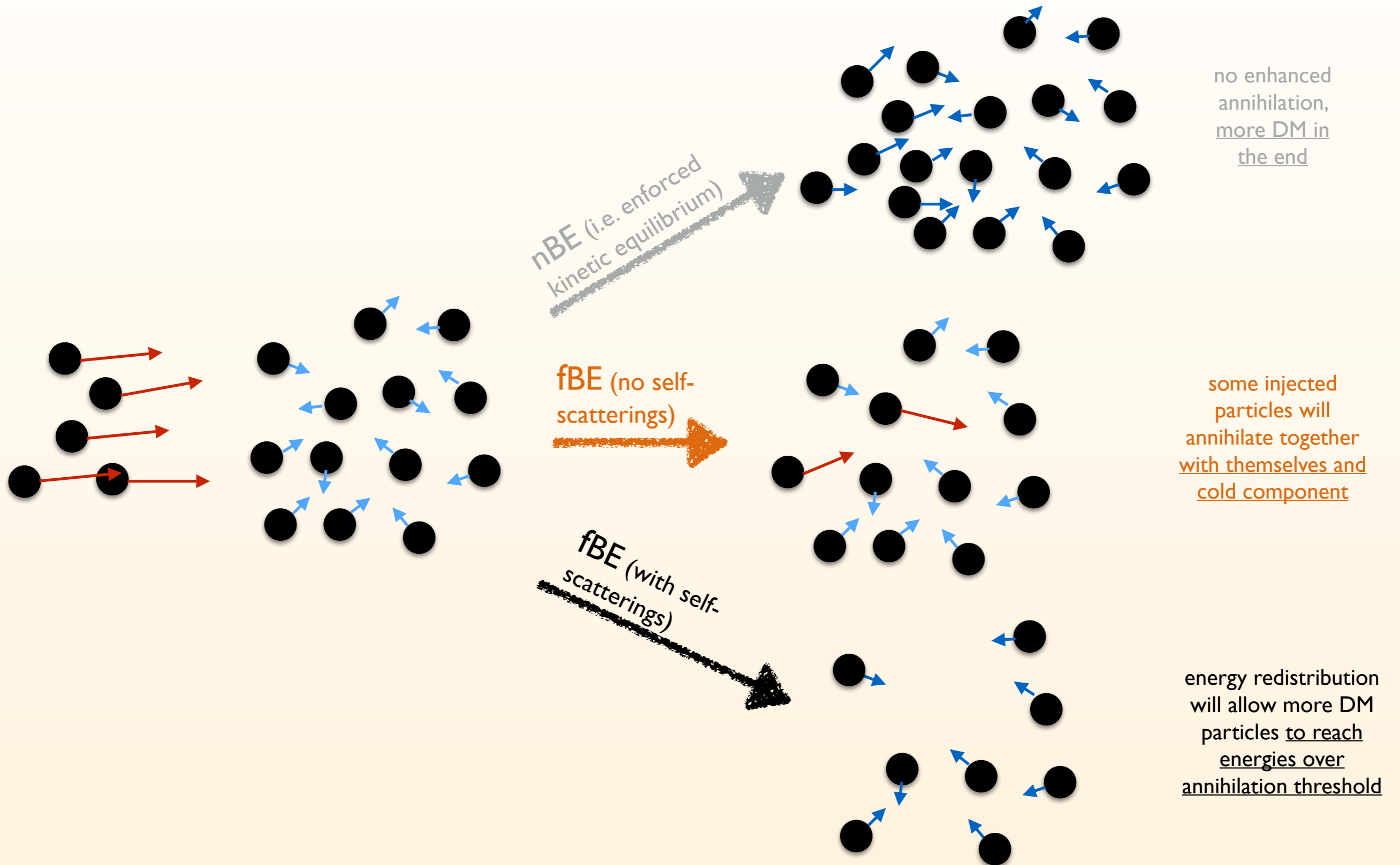
Sudden injection of more DM particles **distorts** $f_\chi(p)$
(e.g. from a decay or annihilation of other states)

- this can **modify the annihilation rate** (if still active)
- how does the **thermalization** due to elastic scatterings happen?

1) DM produced via:

- 1st component from **thermal freeze-out**
- 2nd component from **a decay $\phi \rightarrow \bar{\chi}\chi$**

2) DM annihilation has a **threshold**
 e.g. $\chi\bar{\chi} \rightarrow f\bar{f}$ with $m_\chi \lesssim m_f$



EXAMPLE EVOLUTION

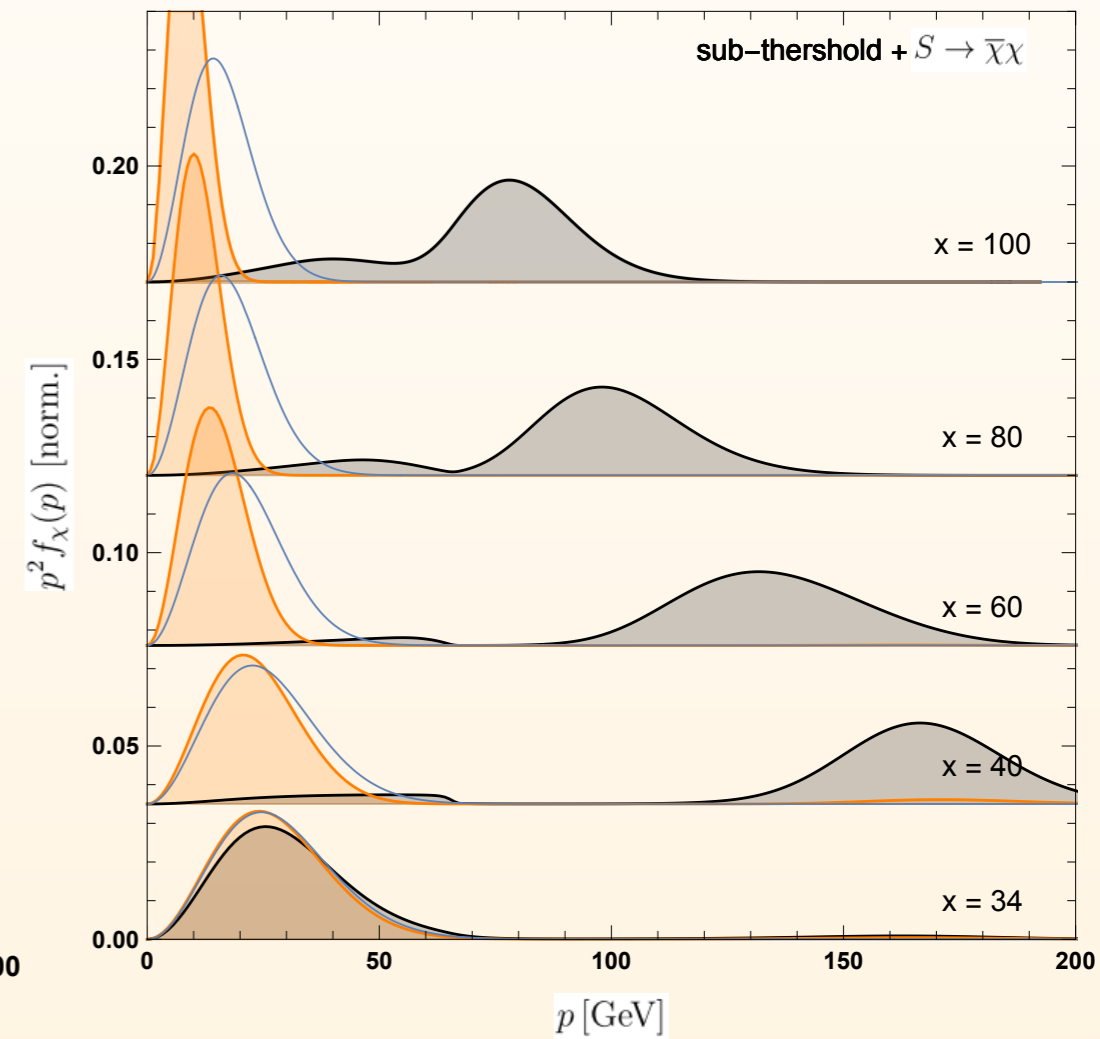
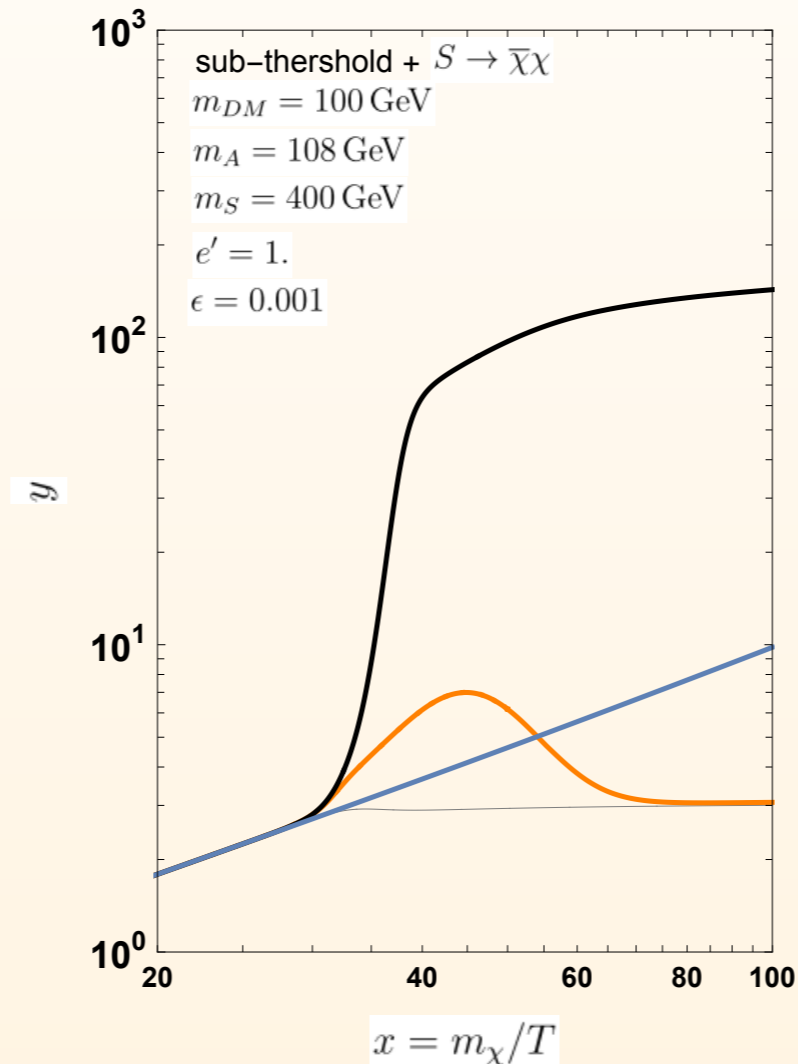
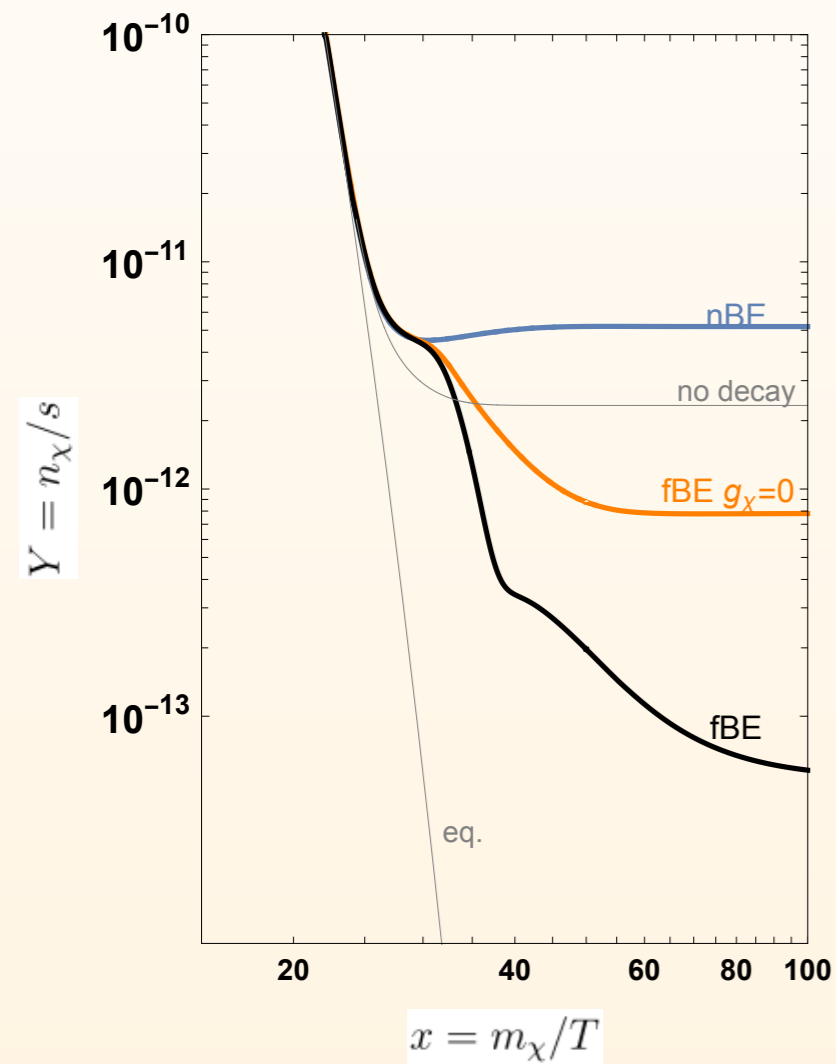
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- 2) DM annihilation has a **threshold**
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$Y \sim$ number density

$y \sim$ temperature

$p^2 f(p) \sim$ momentum distribution



SUMMARY

1. In recent years a **significant progress** in refining the relic density calculations (not yet fully implemented in public codes!)

2. **Kinetic equilibrium** is a necessary (often implicit) assumption for standard relic density calculations in all the numerical tools...
...while it is not always warranted!

3. Introduced coupled **system of Boltzmann eqs. for 0th and 2nd moments (cBE)** allows for much more accurate treatment while the **full phase space Boltzmann equation (fBE)** can be also successfully solved for higher precision and/or to obtain result for $f_{DM}(p)$

(we also introduced  a new tool to extend the current capabilities to the regimes **beyond kinetic equilibrium**)

4. **Multi-component sectors**, when studied **at the fBE level**, can reveal quite unexpected behavior