

DARK MATTER FREEZE-OUT AND FREEZE-IN BEYOND KINETIC EQUILIBRIUM

Andrzej Hryczuk



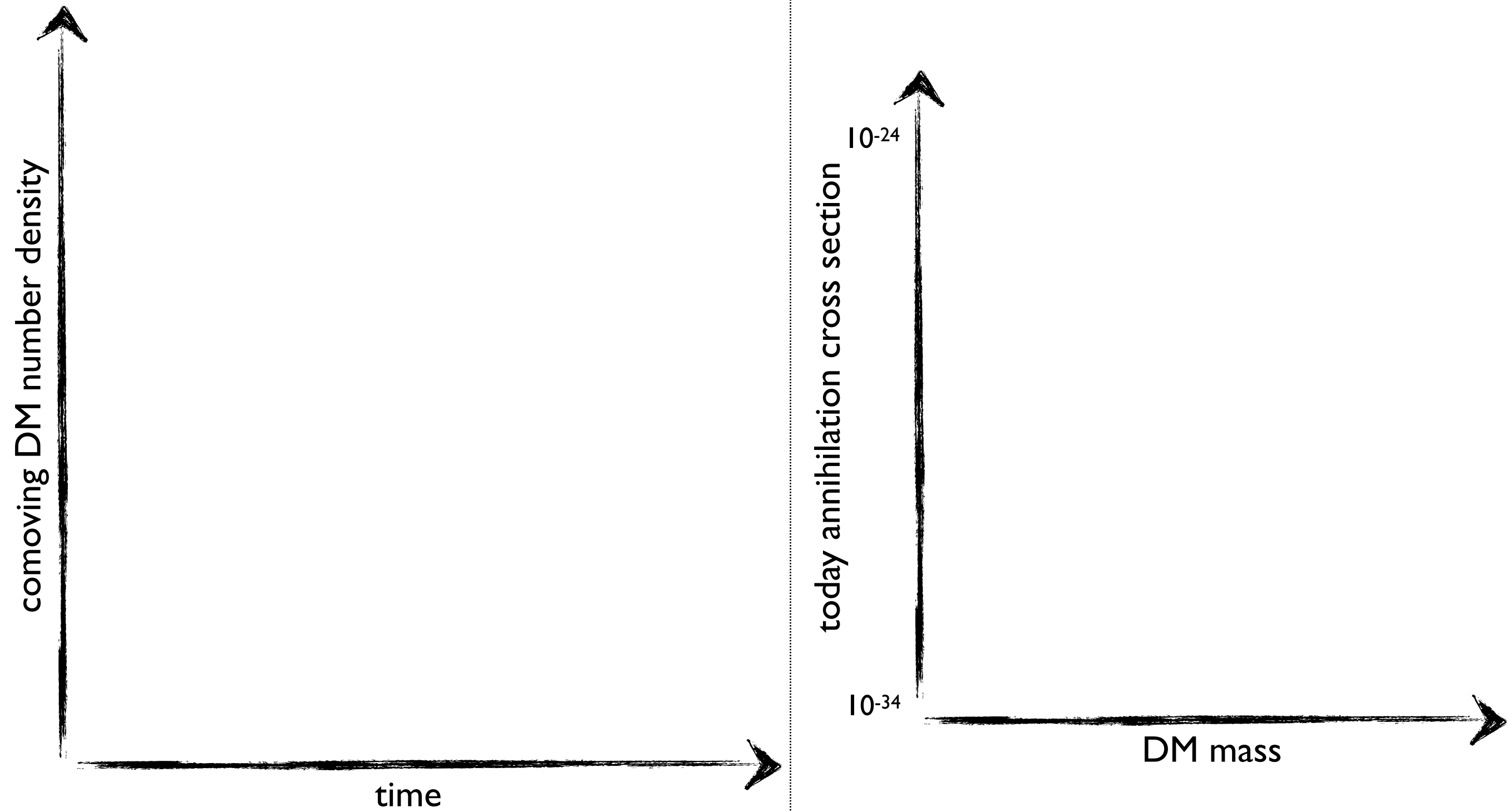
based on:

A.H. & M. Laletin [2204.07078](#)

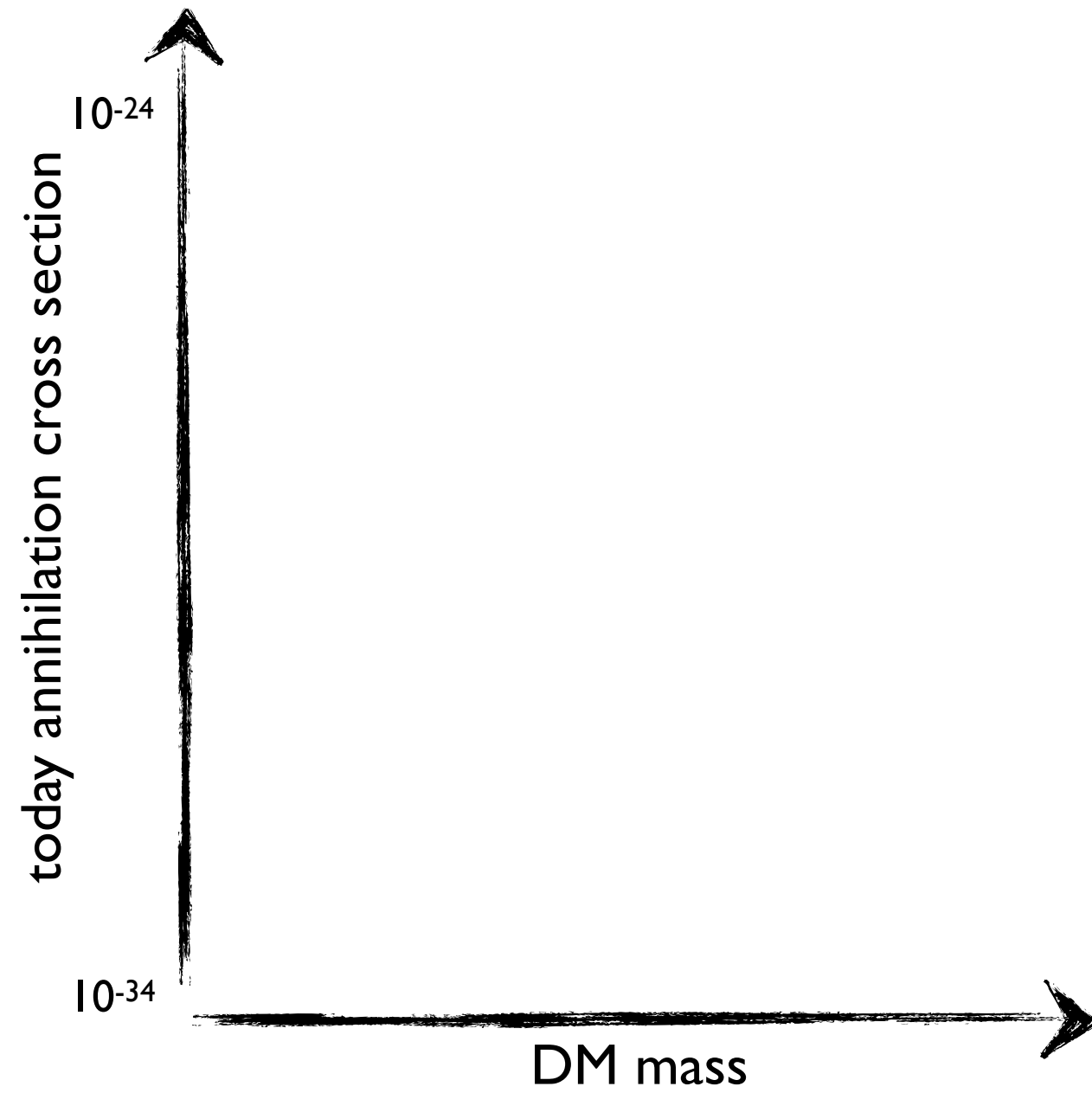
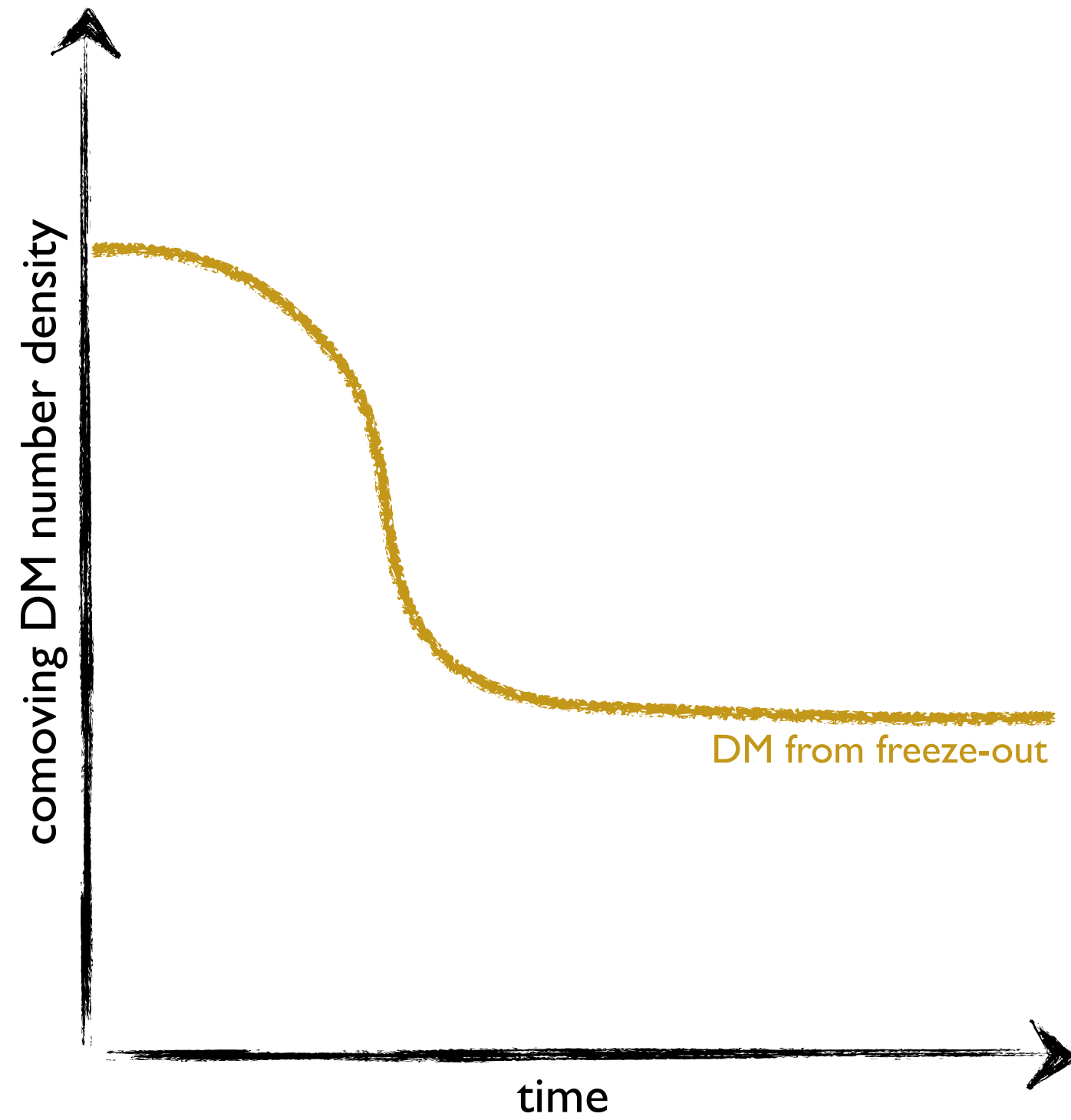
A.H. & M. Laletin [2104.05684](#)

and **T. Binder, T. Bringmann, M. Gustafsson & A.H.** [1706.07433](#), [2103.01944](#)

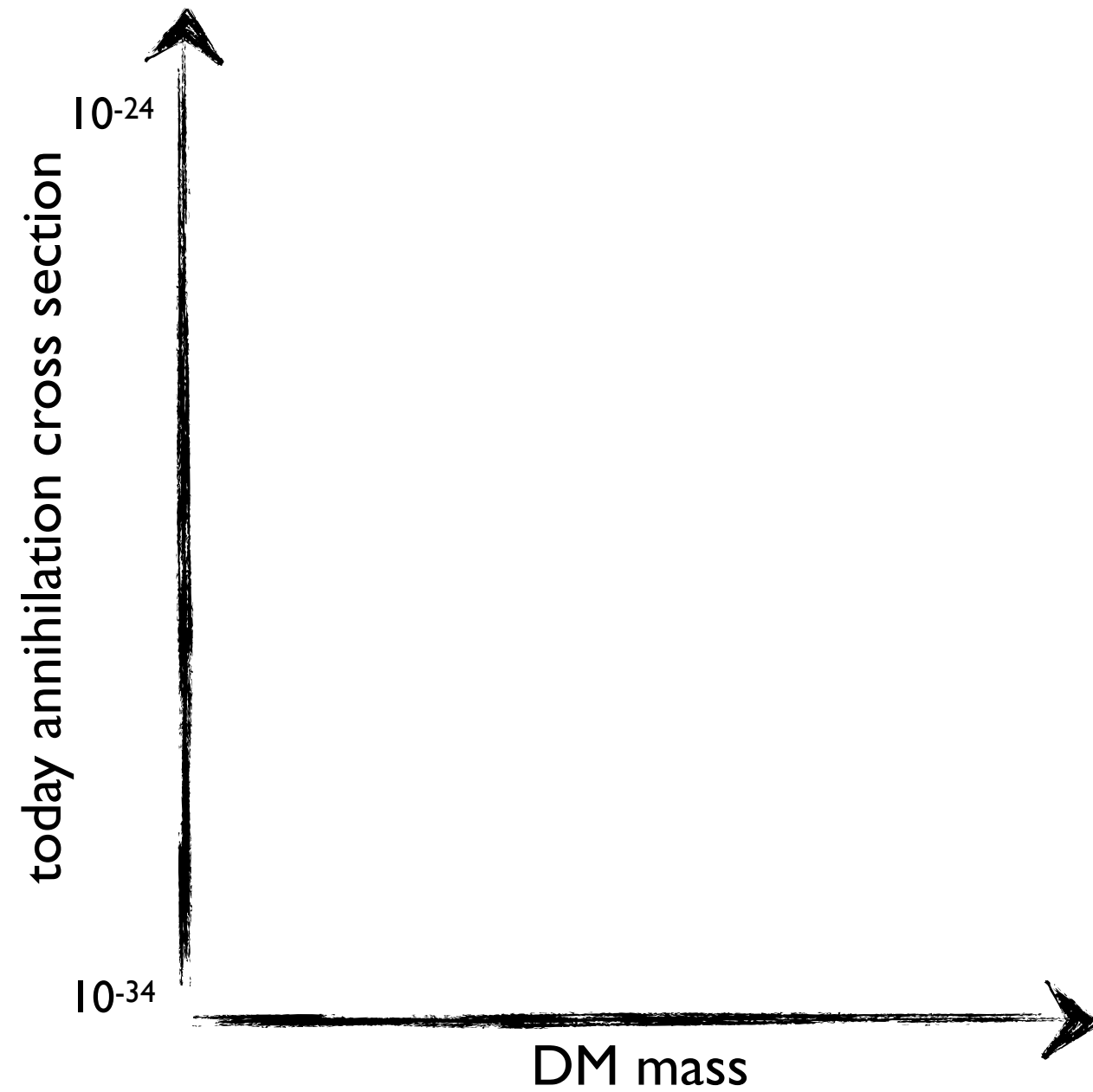
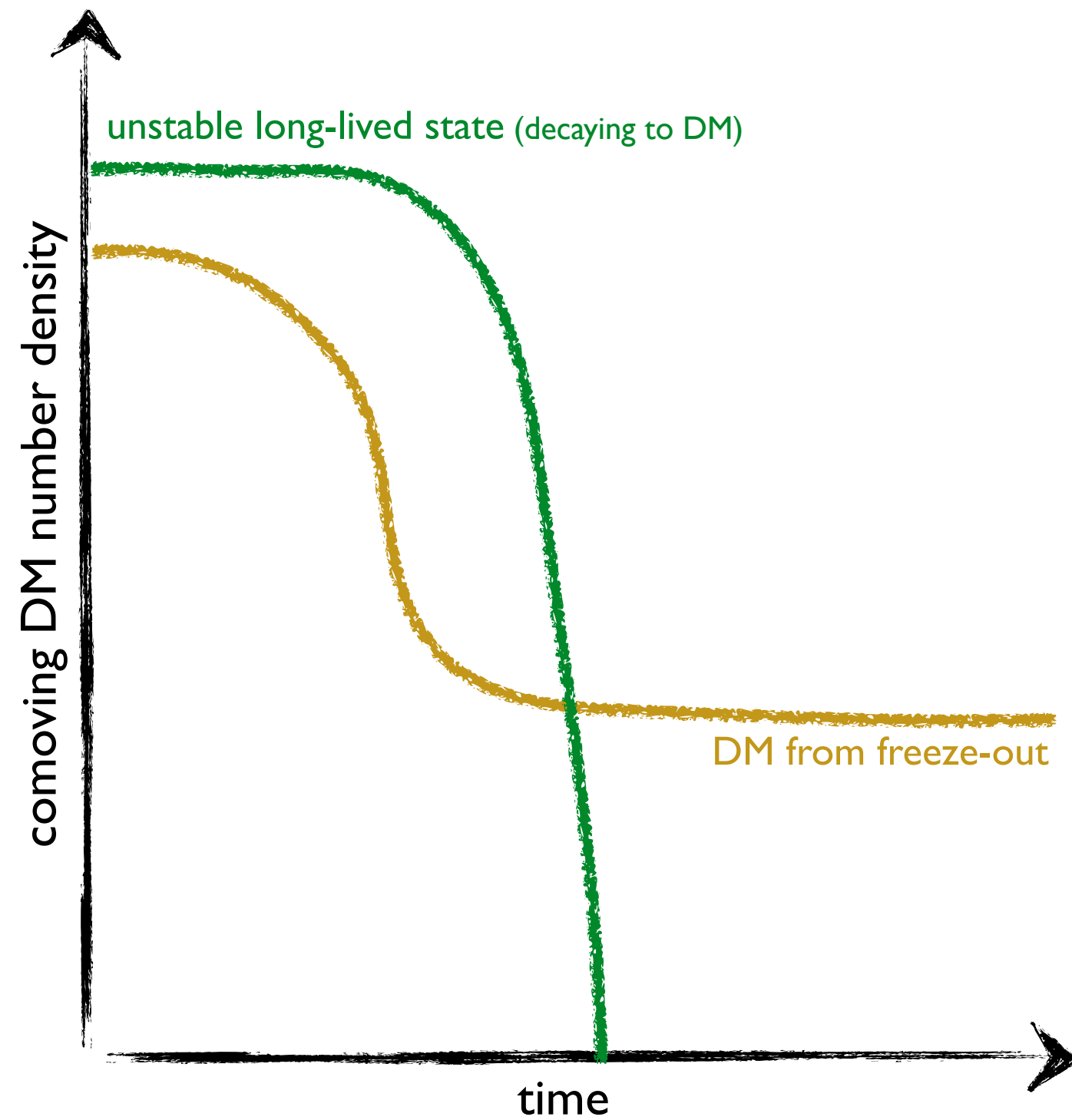
IN CASE YOU'RE NOT INTERESTED IN WHAT FOLLOWS...



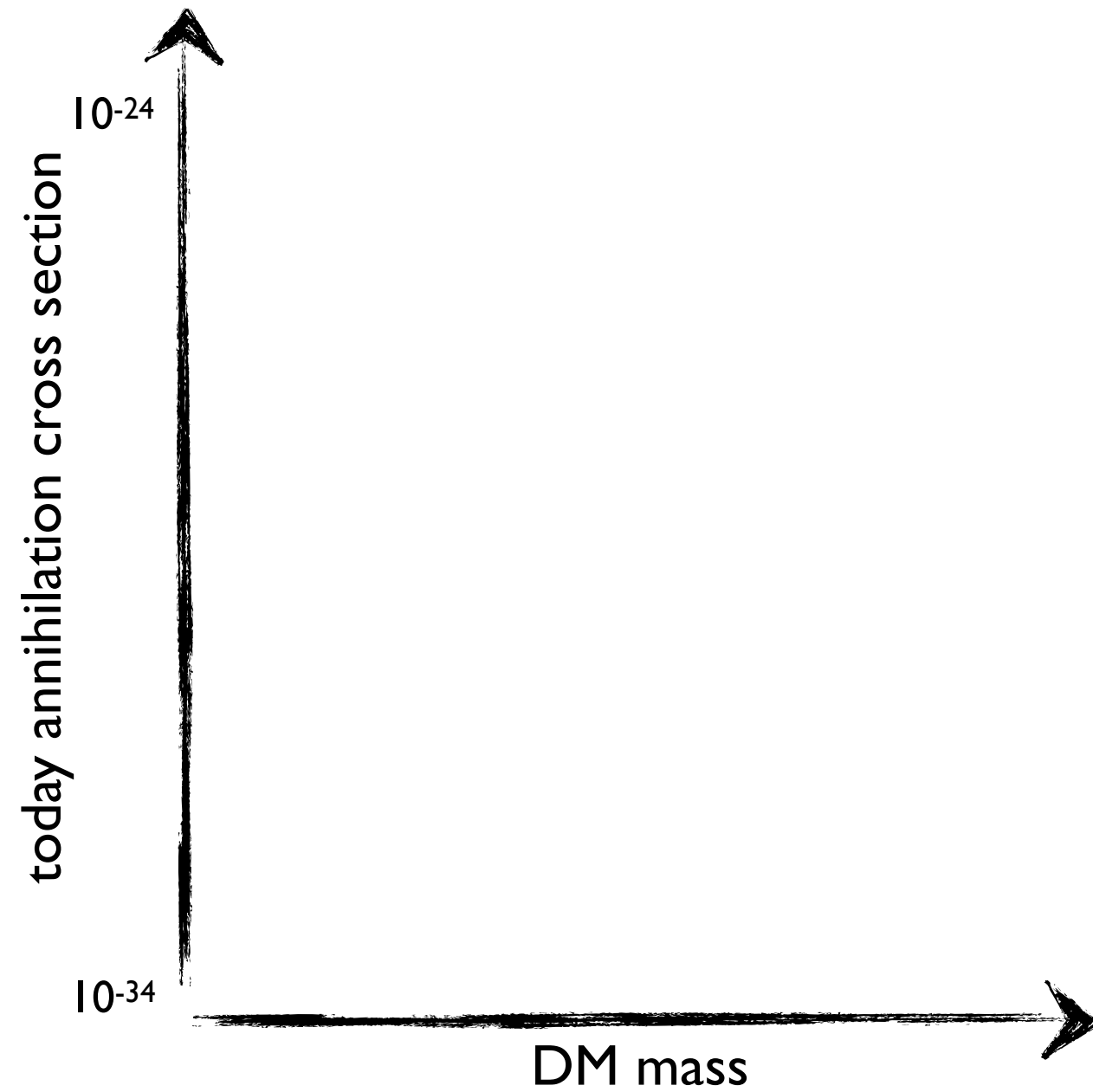
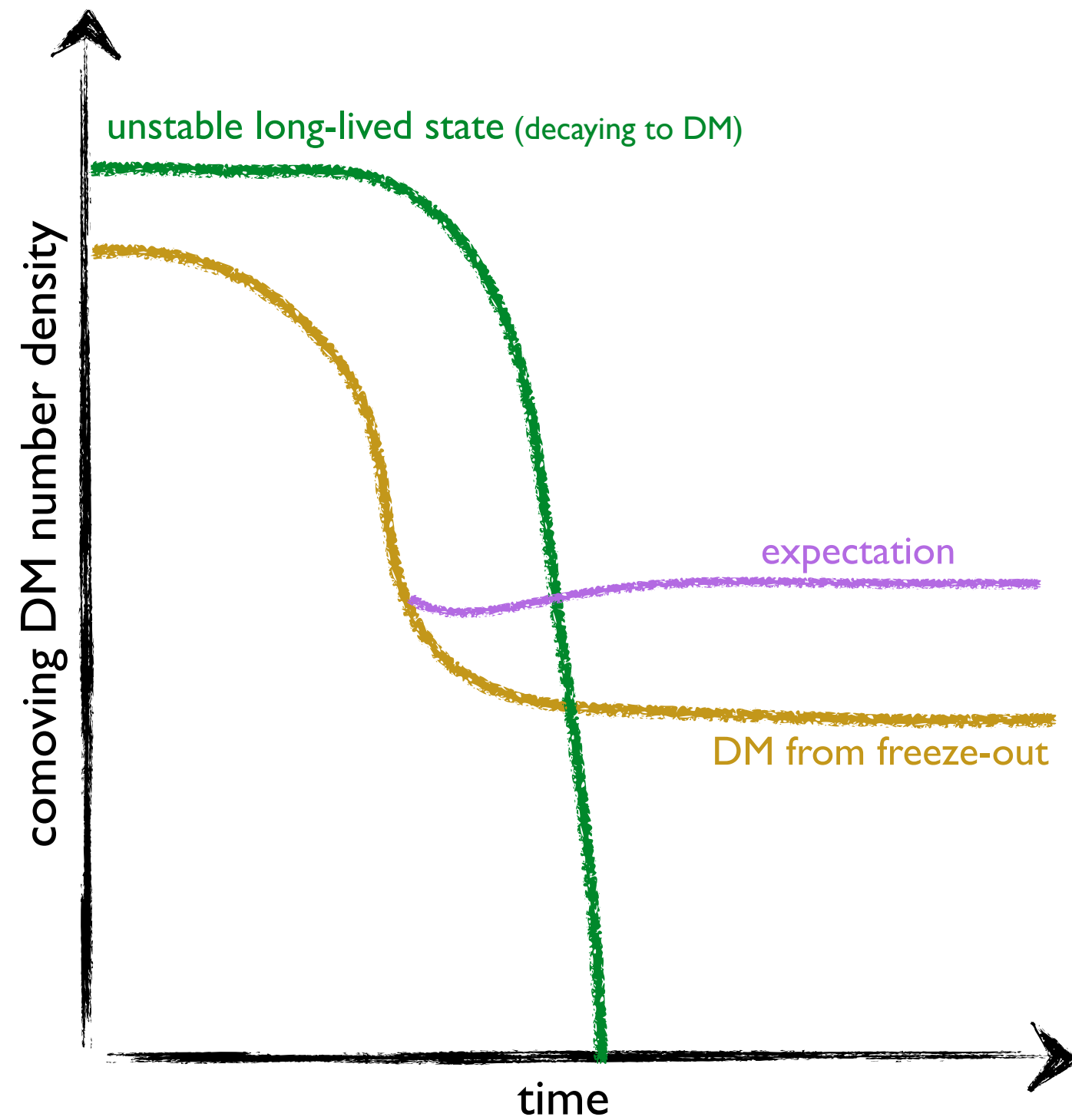
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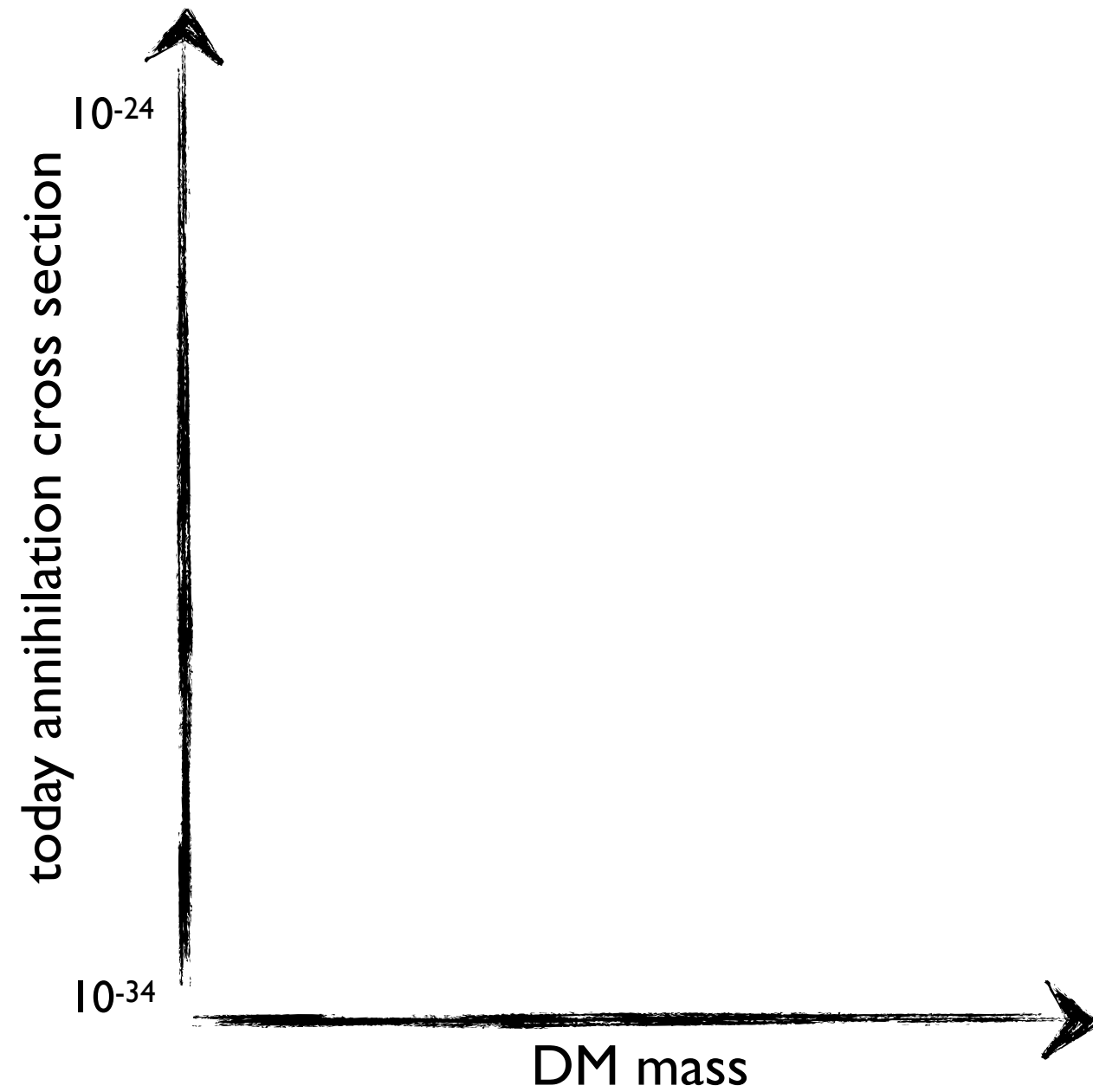
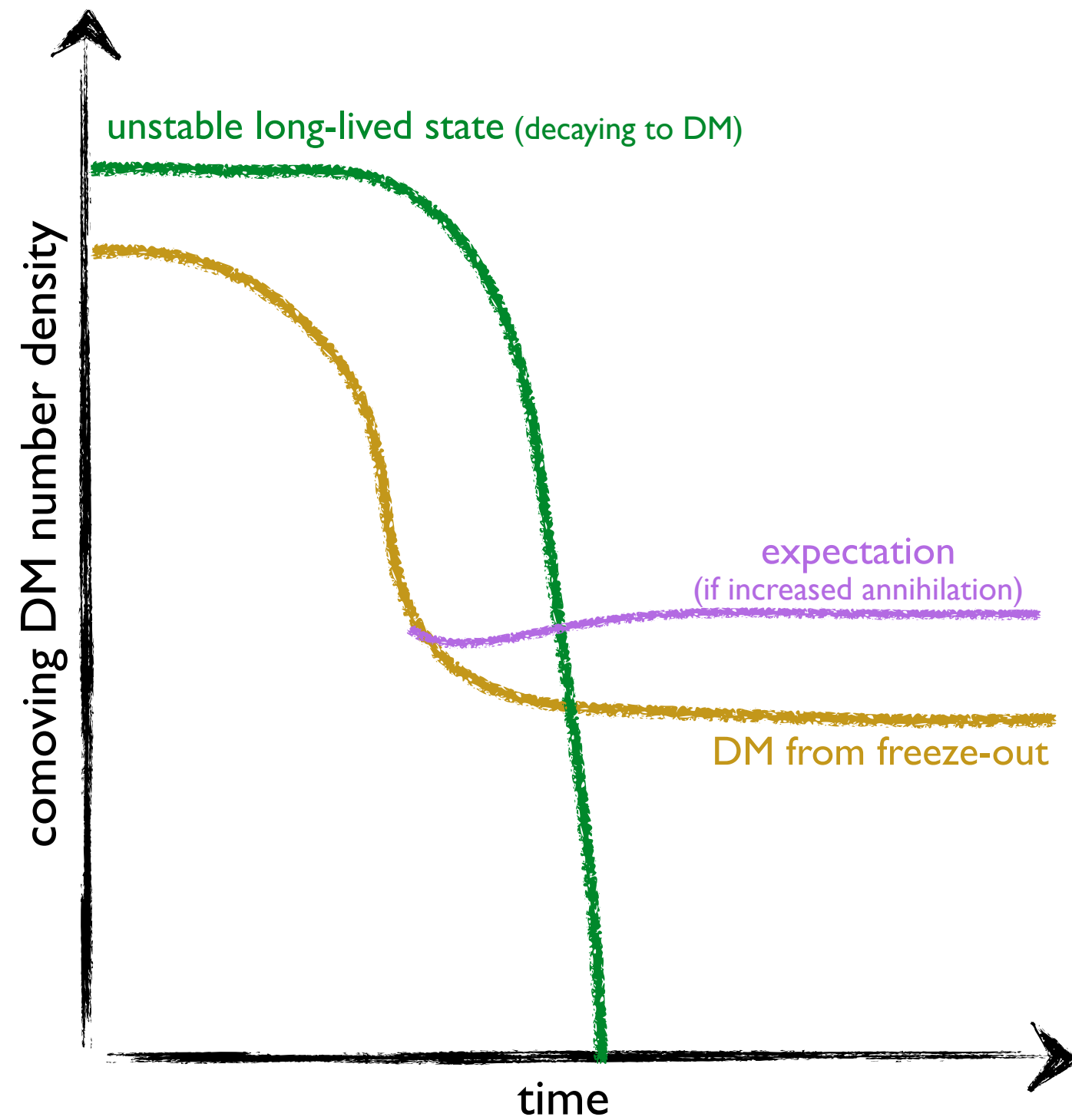
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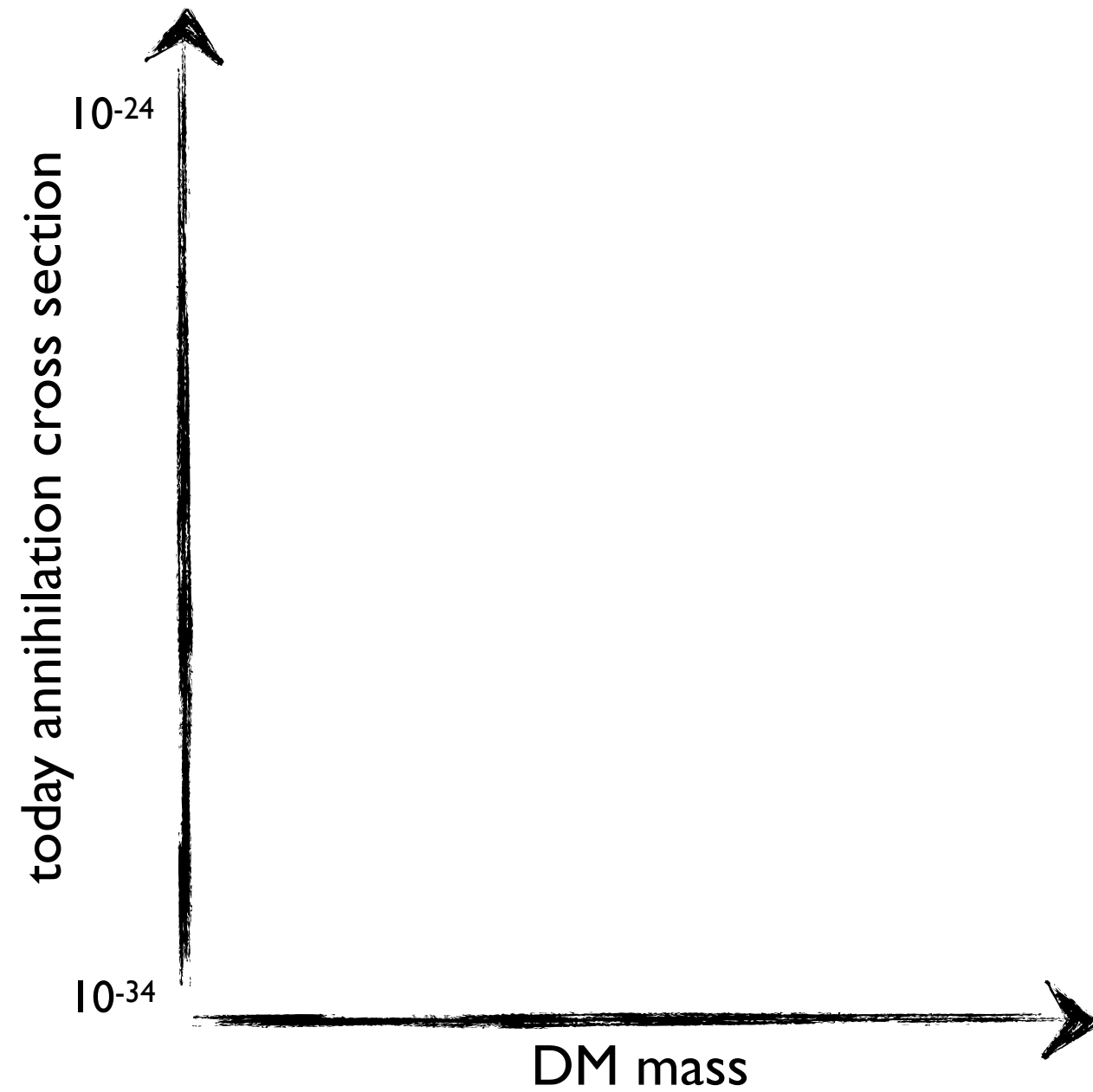
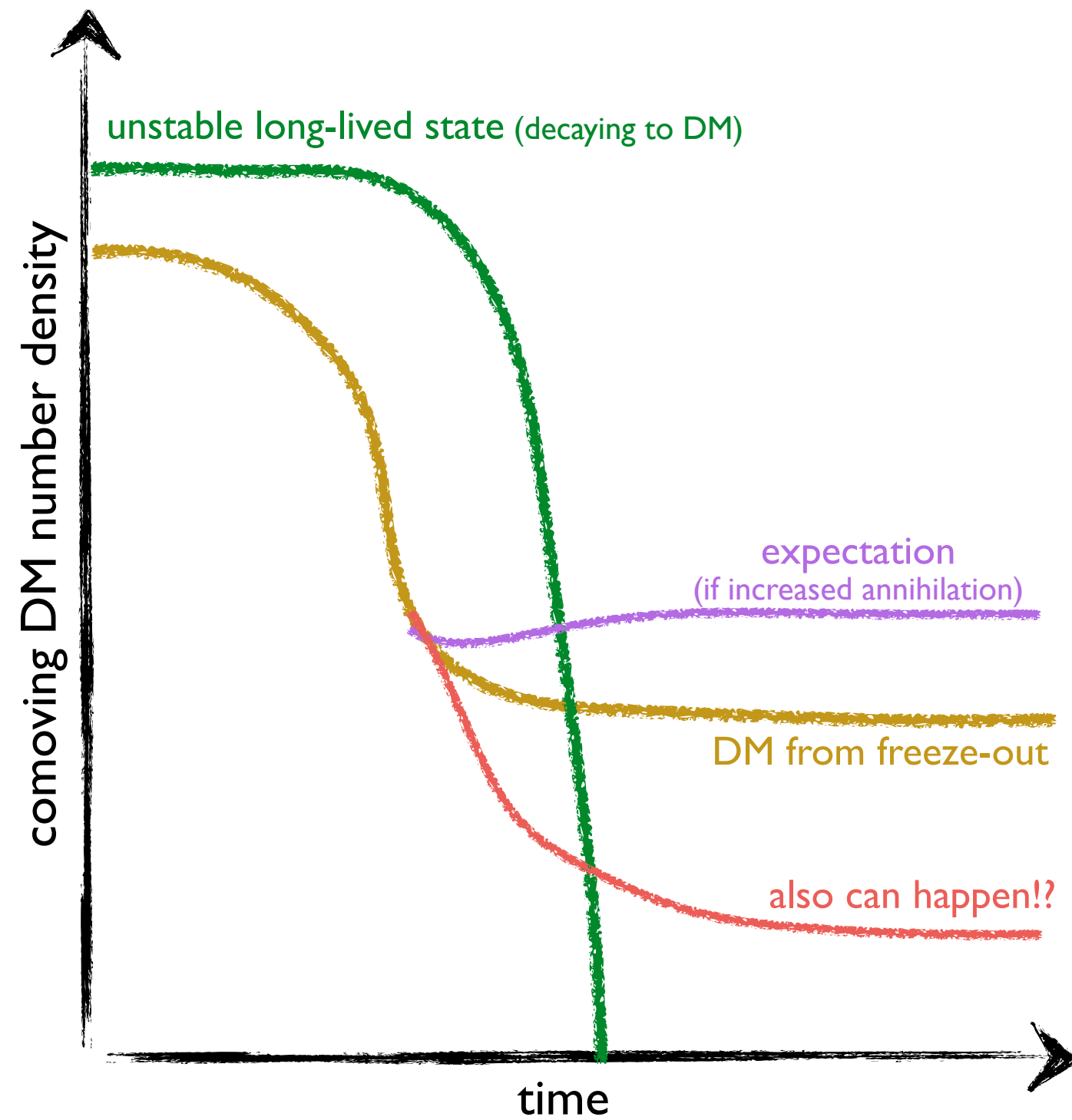
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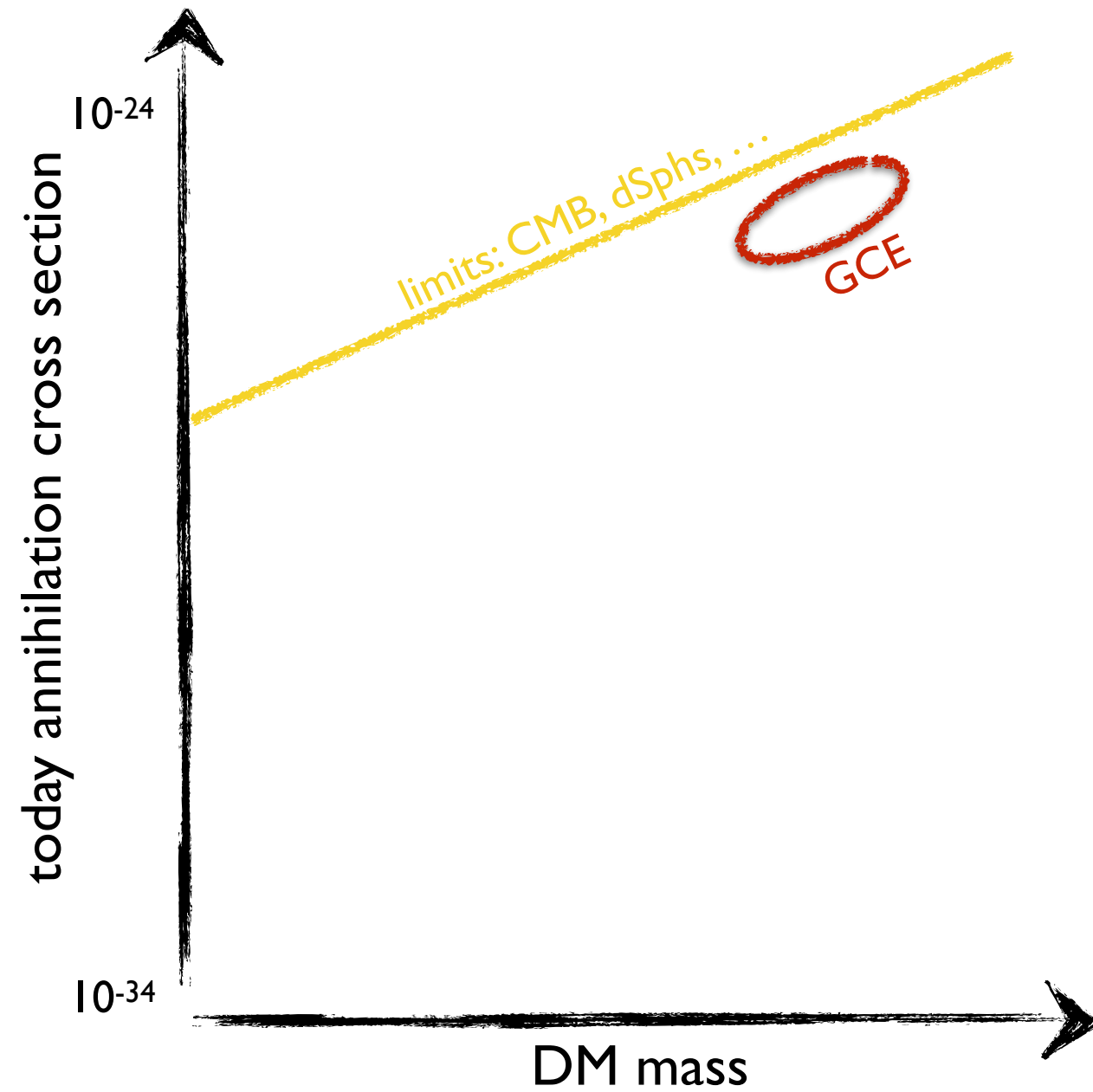
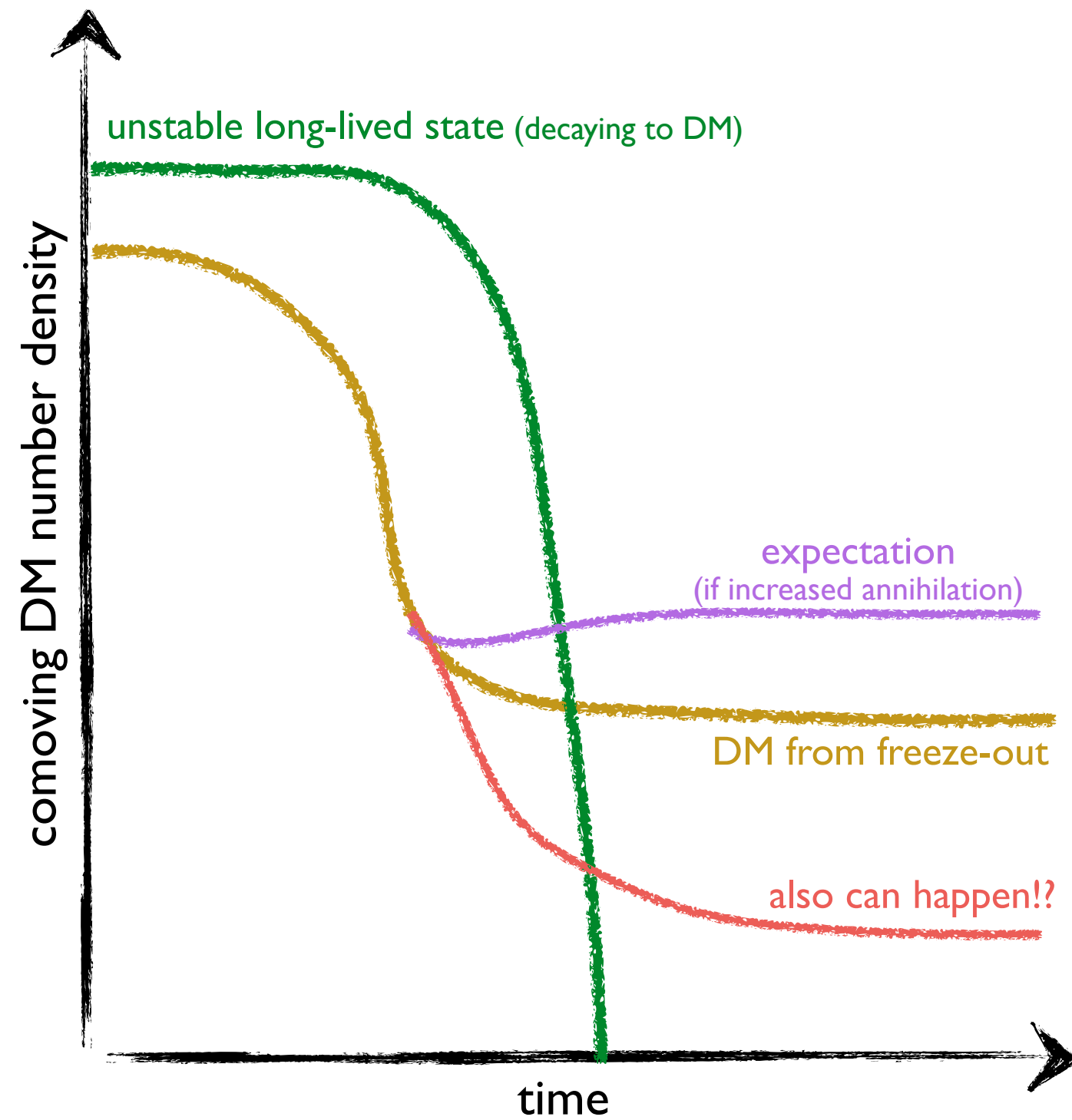
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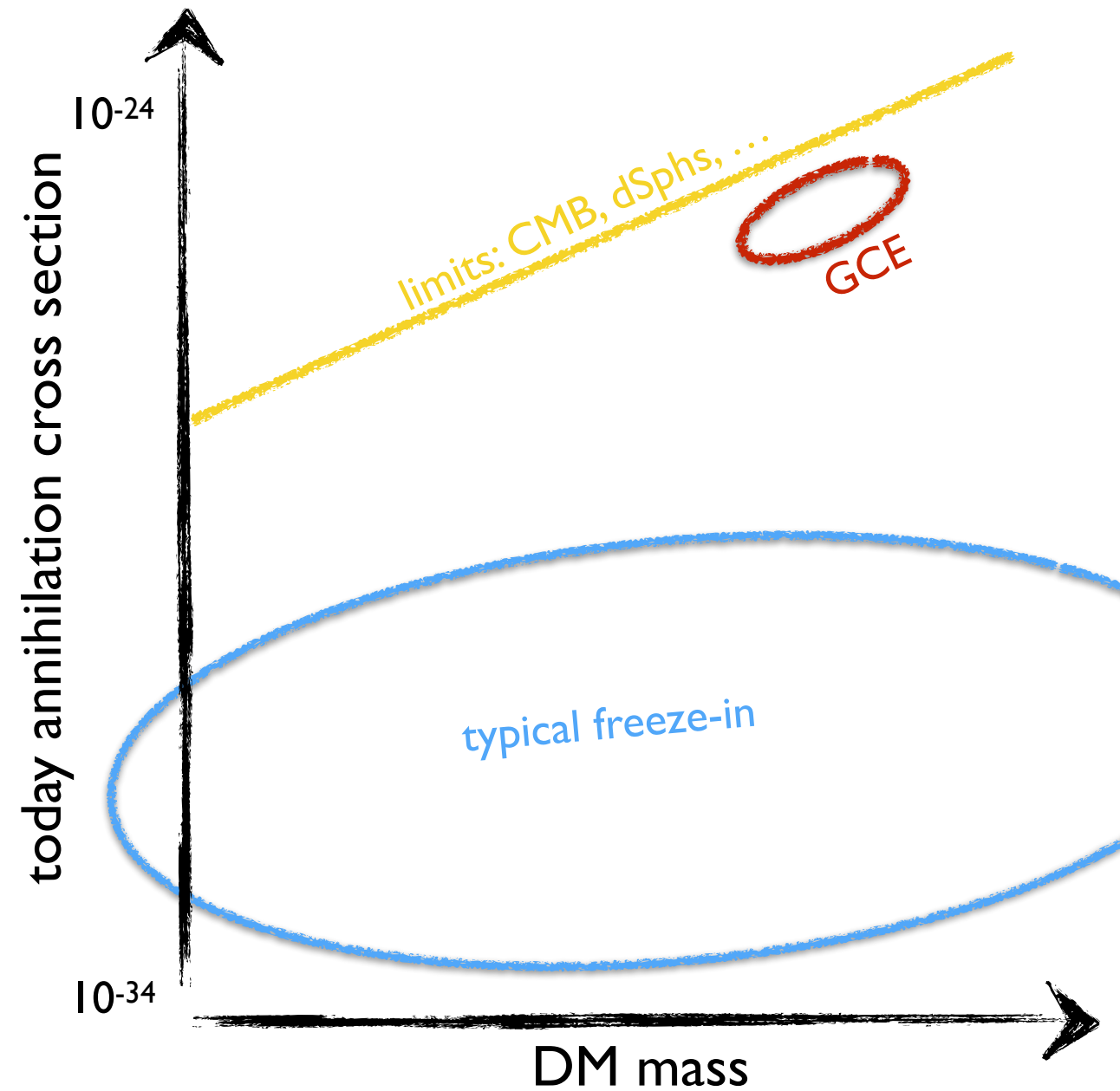
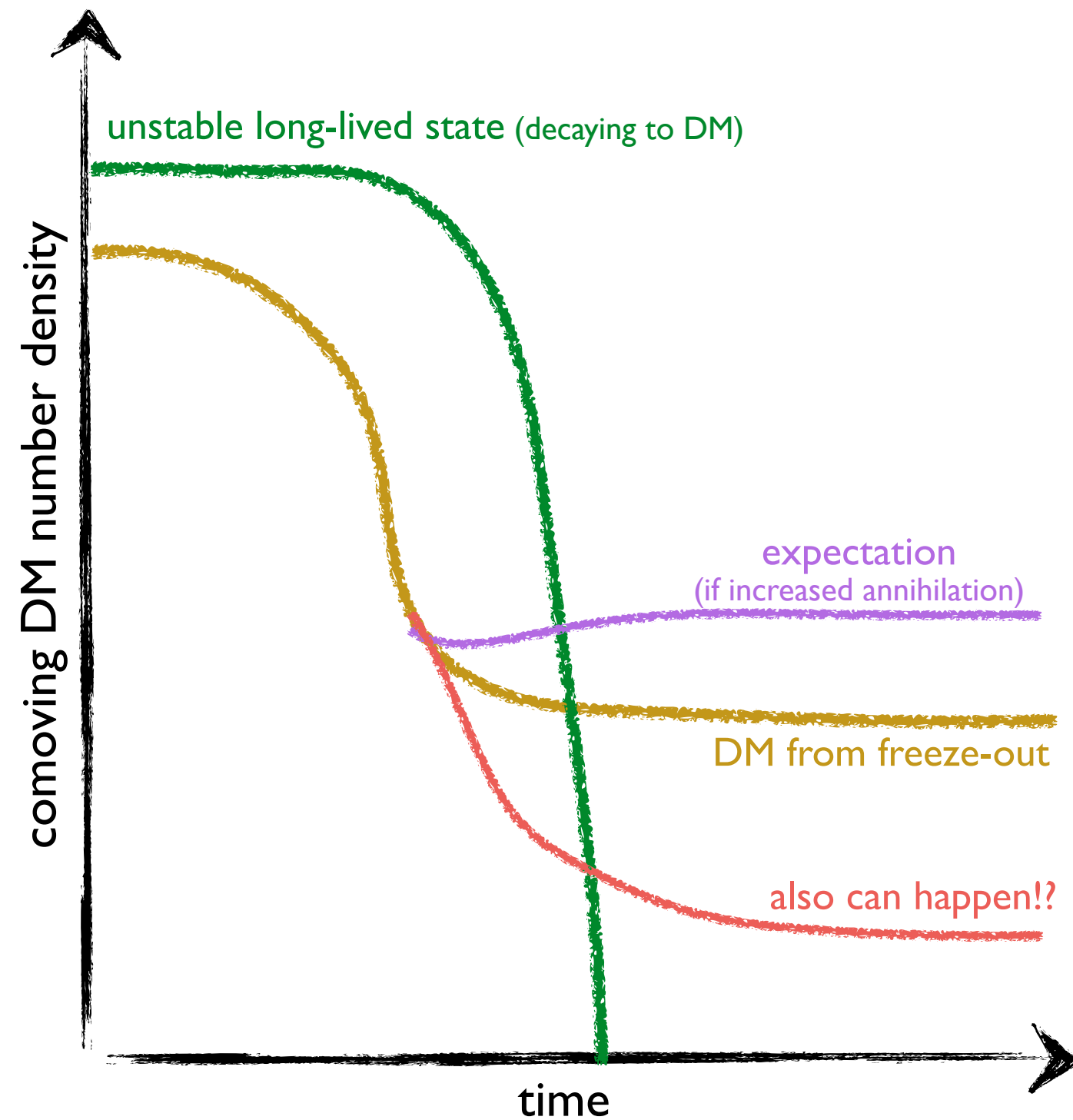
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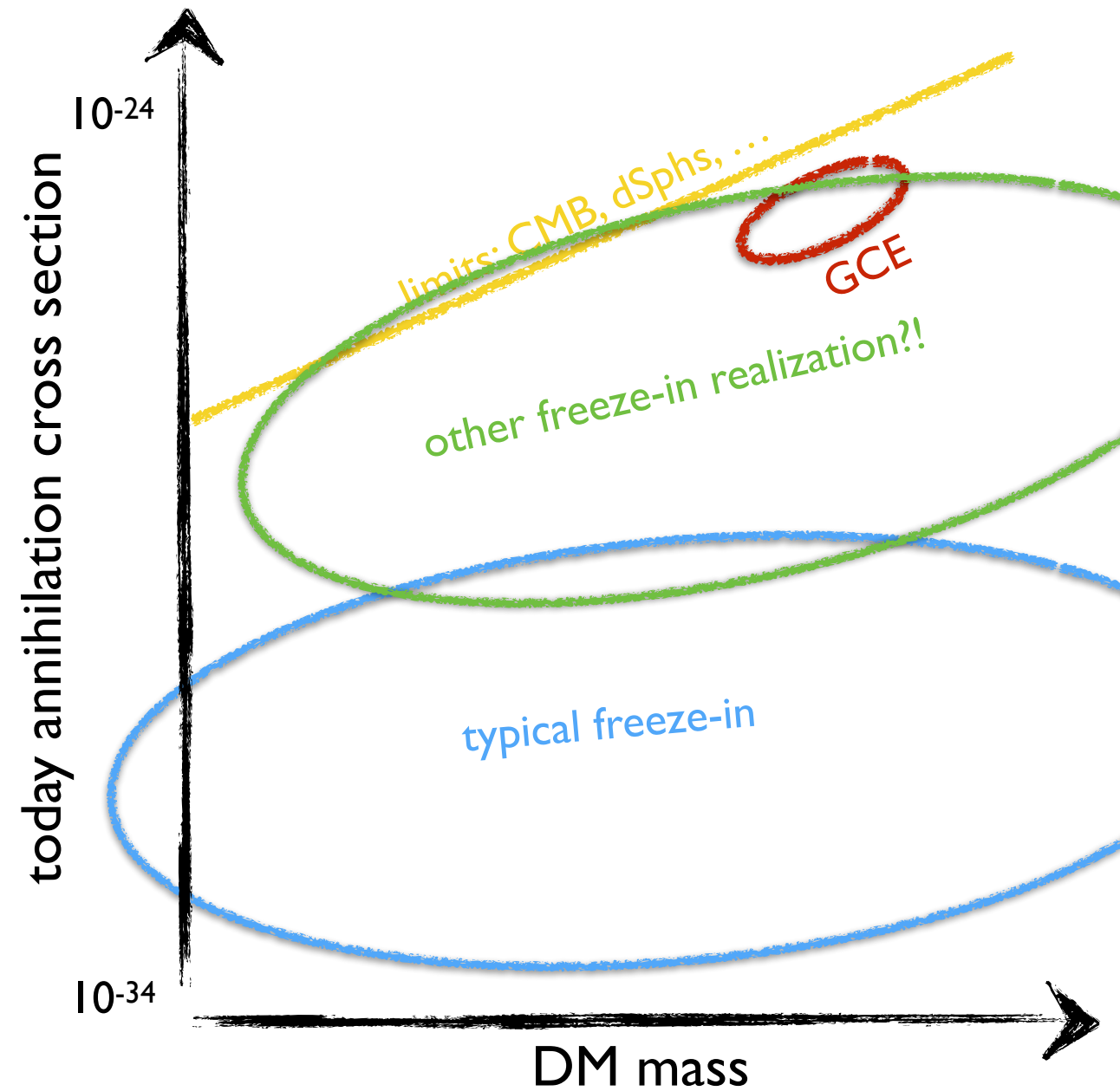
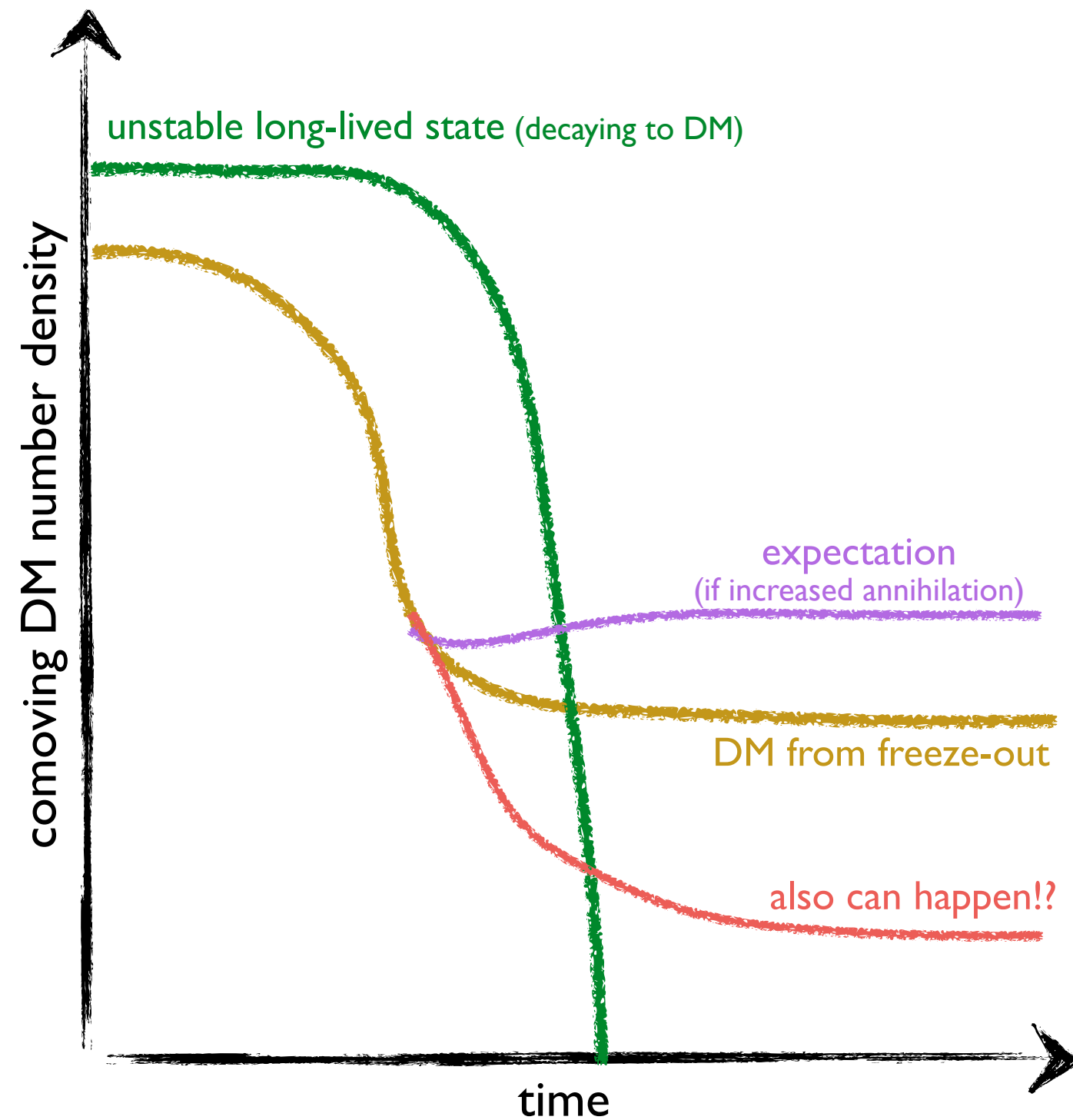
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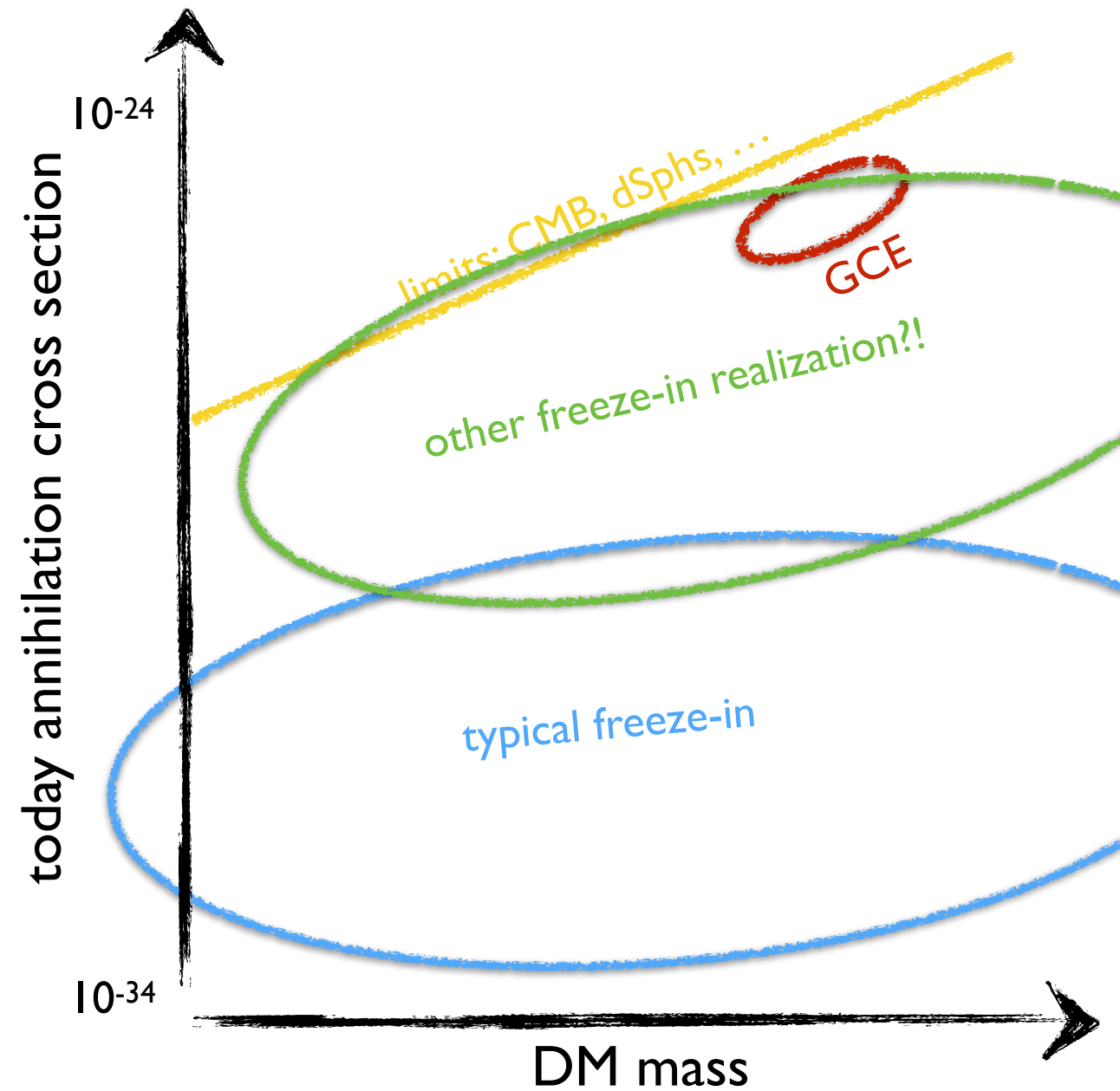
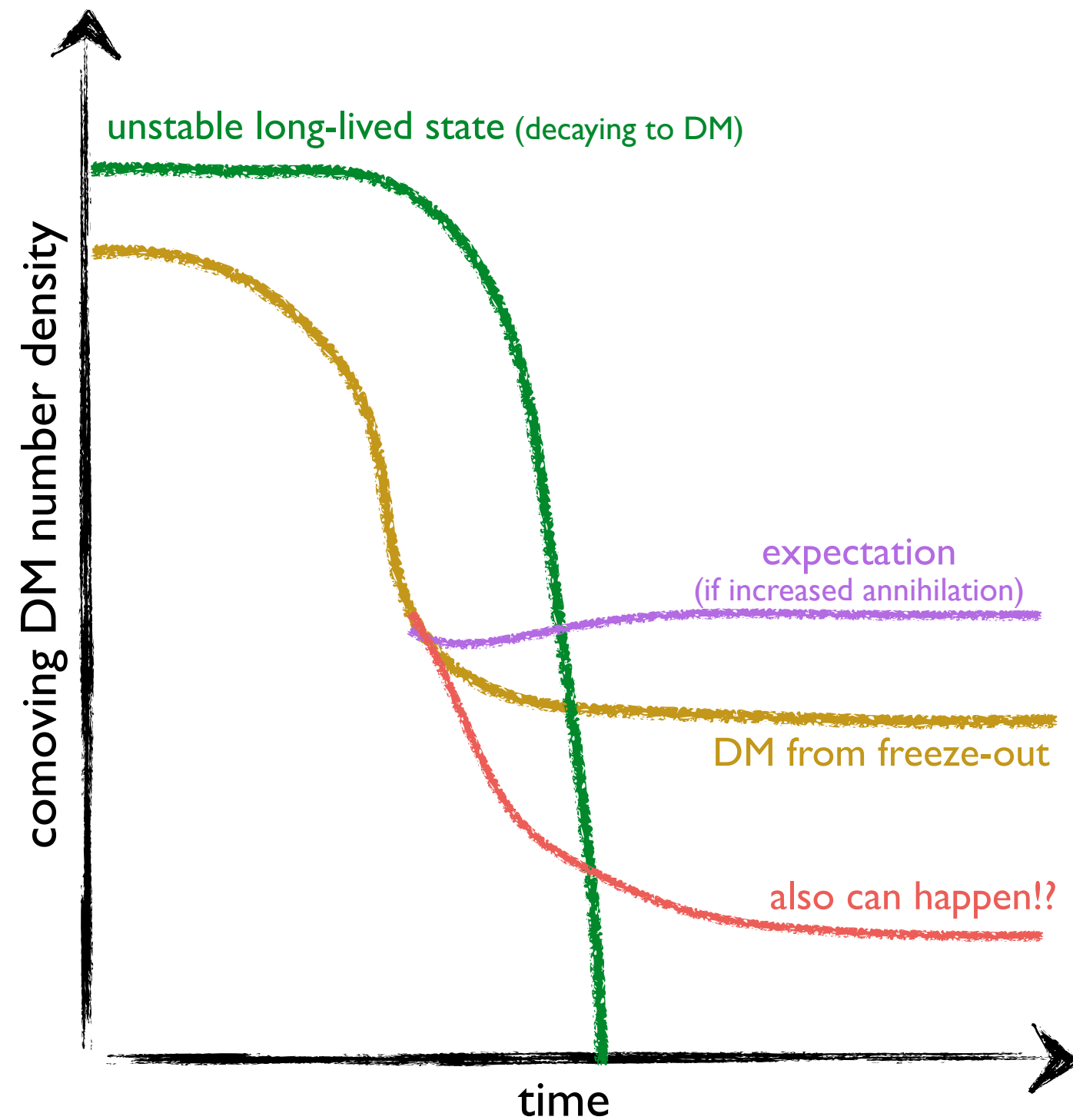
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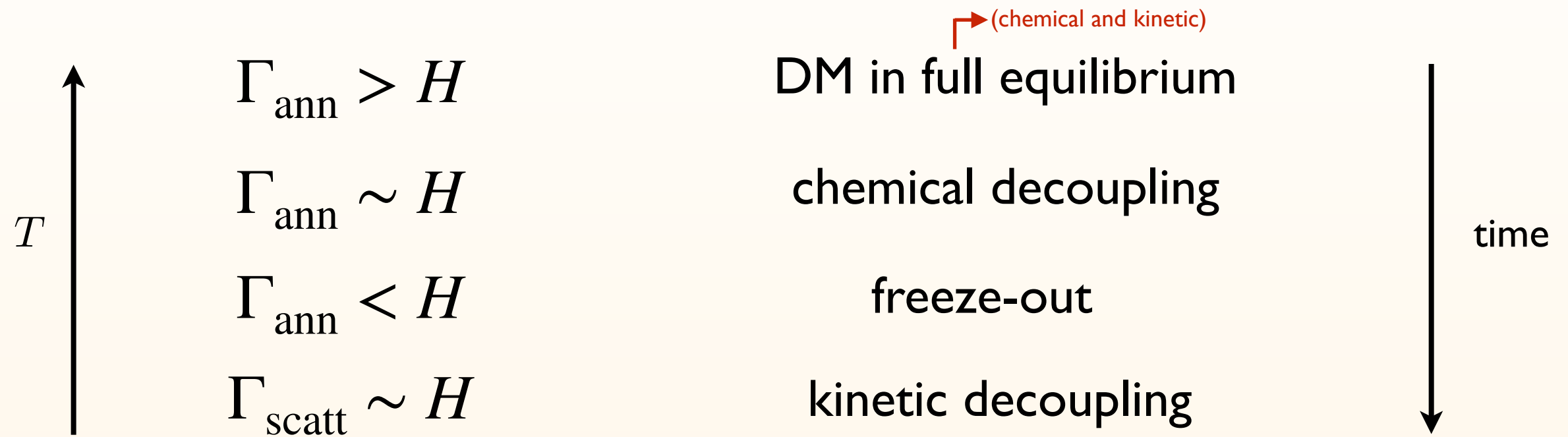
IN CASE YOU'RE NOT INTERESTED IN WHAT FOLLOWS...



TO SEE WHY AND LEARN MORE STAY TUNED :)

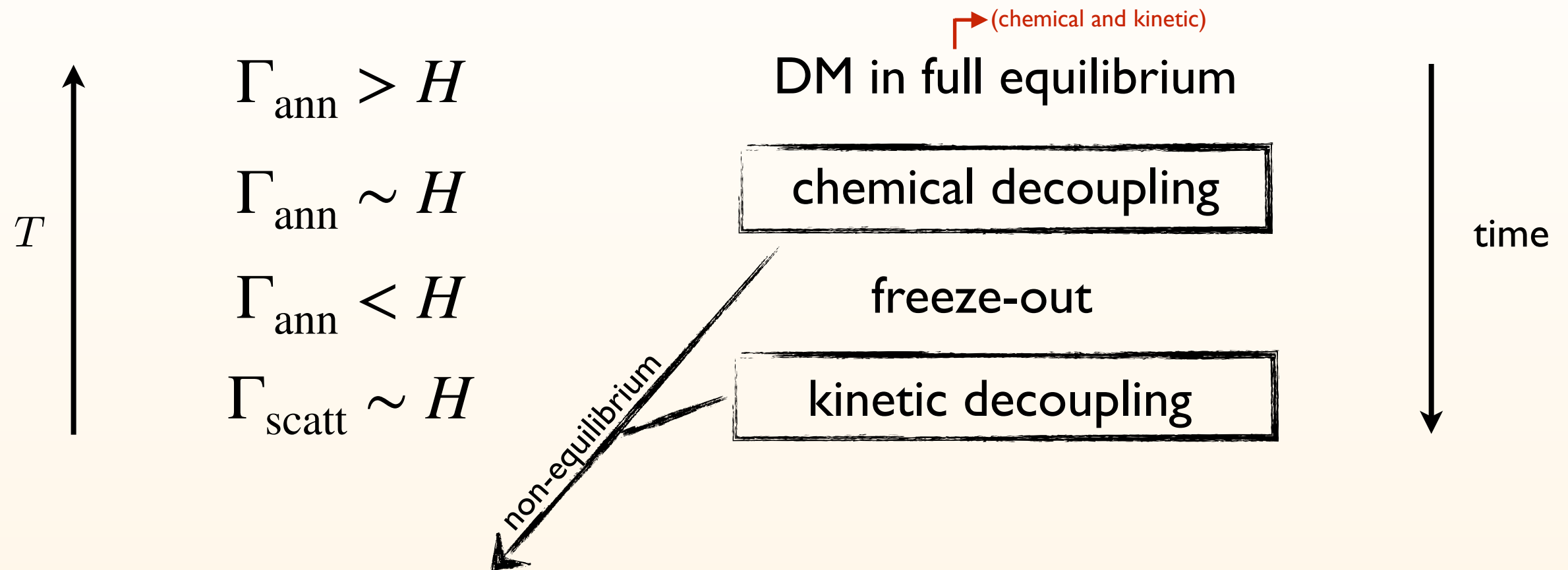
THERMAL RELIC DENSITY

STANDARD SCENARIO



THERMAL RELIC DENSITY

STANDARD SCENARIO



time evolution of $f_\chi(p)$ in kinetic theory:

$$E (\partial_t - H \vec{p} \cdot \nabla_{\vec{p}}) f_\chi = \mathcal{C}[f_\chi]$$

Liouville operator in
FRW background

the collision term

THERMAL RELIC DENSITY

STANDARD APPROACH

Boltzmann equation for $f_\chi(p)$:

$$E (\partial_t - H\vec{p} \cdot \nabla_{\vec{p}}) f_\chi = \mathcal{C}[f_\chi]$$

*assumptions for using Boltzmann eq:
classical limit, molecular chaos,...

...for derivation from thermal QFT
see e.g., 1409.3049

THERMAL RELIC DENSITY

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$$E (\partial_t - H\vec{p} \cdot \nabla_{\vec{p}}) f_\chi = \mathcal{C}[f_\chi]$$

\Downarrow integrate over p
(i.e. take 0th moment)

$$\frac{dn_\chi}{dt} + 3Hn_\chi = - \langle \sigma_{\chi\bar{\chi} \rightarrow ij} \sigma_{\text{rel}} \rangle^{\text{eq}} (n_\chi n_{\bar{\chi}} - n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}})$$

where the **thermally averaged cross section**:

$$\langle \sigma_{\chi\bar{\chi} \rightarrow ij} v_{\text{rel}} \rangle^{\text{eq}} = - \frac{h_\chi^2}{n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}}} \int \frac{d^3\vec{p}_\chi}{(2\pi)^3} \frac{d^3\vec{p}_{\bar{\chi}}}{(2\pi)^3} \sigma_{\chi\bar{\chi} \rightarrow ij} v_{\text{rel}} f_\chi^{\text{eq}} f_{\bar{\chi}}^{\text{eq}}$$

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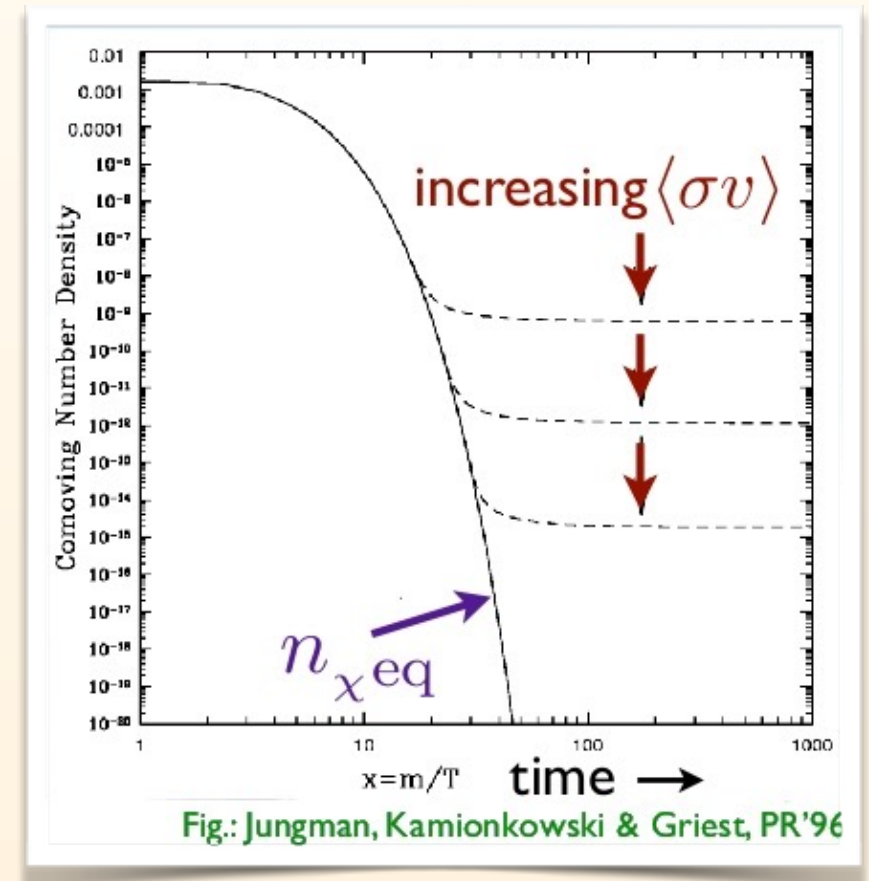
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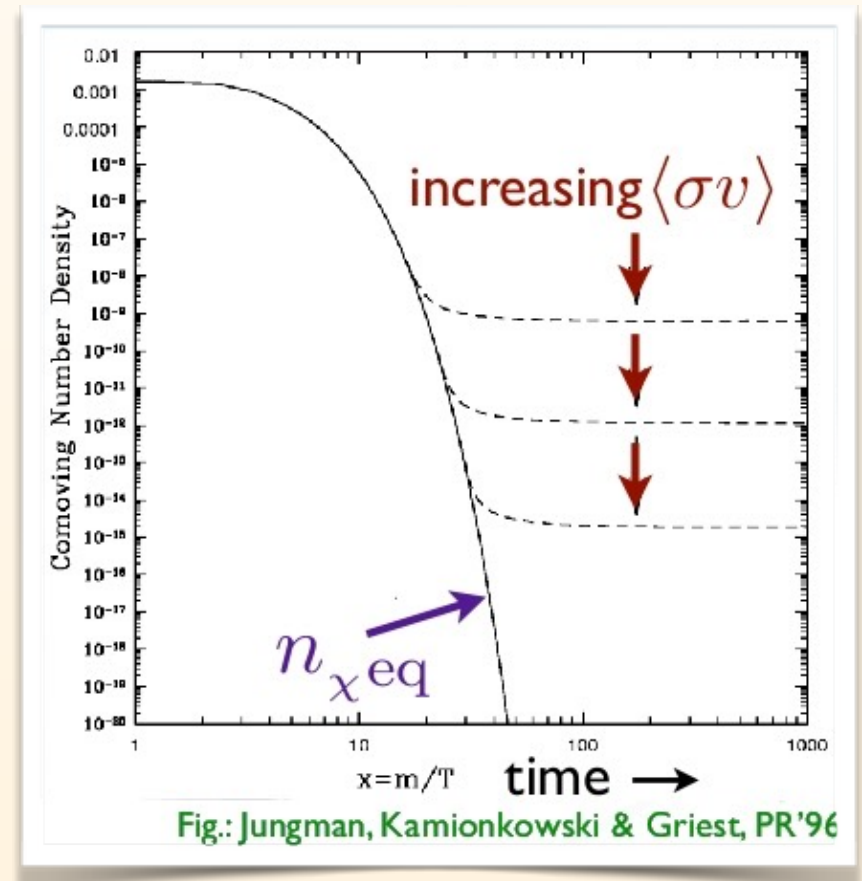
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Critical assumption:
kinetic equilibrium at chemical decoupling

$$f_\chi \sim a(T) f_\chi^{\text{eq}}$$

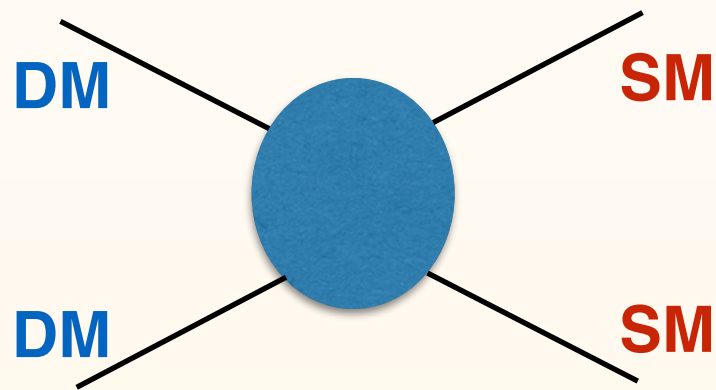
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FREEZE-OUT vs. DECOUPLING

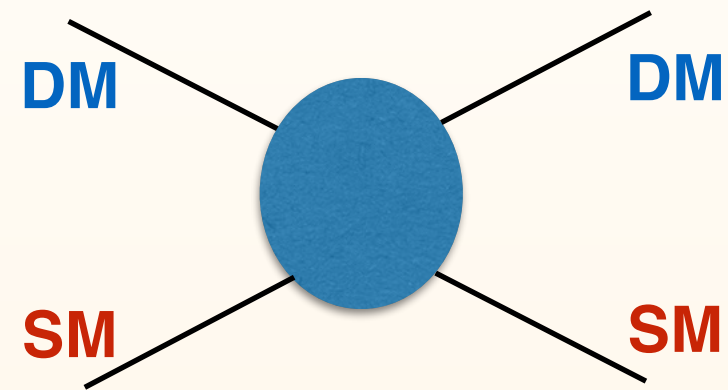
annihilation



$$\sum_{\text{spins}} |\mathcal{M}^{\text{pair}}|^2 = F(p_1, p_2, p'_1, p'_2)$$

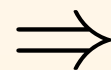
crossing sym.
 \longleftrightarrow

(elastic) scattering



$$\sum_{\text{spins}} |\mathcal{M}^{\text{scatt}}|^2 = F(k, -k', p', -p)$$

Boltzmann suppression of **DM** vs. **SM**



scatterings typically more frequent

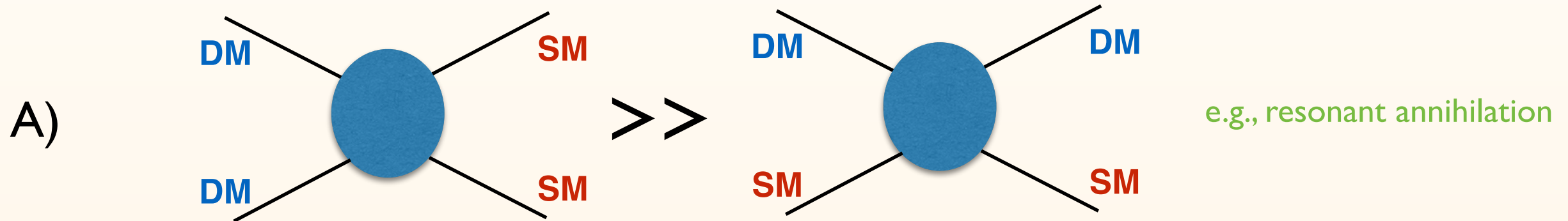
dark matter frozen-out but typically
 still kinetically coupled to the plasma

Schmid, Schwarz, Widern '99; Green, Hofmann, Schwarz '05

EARLY KINETIC DECOUPLING?

A **necessary** and **sufficient** condition: scatterings weaker than annihilation
i.e. rates around freeze-out: $H \sim \Gamma_{\text{ann}} \gtrsim \Gamma_{\text{el}}$

Possibilities:



B) Boltzmann suppression of **SM** as strong as for **DM**
e.g., below threshold annihilation (forbidden-like DM)

C) Scatterings and annihilation have different structure
e.g., semi-annihilation, 3 to 2 models,...

D) Multi-component dark sectors
e.g., additional sources of DM from late decays, ...

HOW TO GO BEYOND KINETIC EQUILIBRIUM?

All information is in the full BE:

both about chemical ("normalization") and kinetic ("shape") equilibrium/decoupling

$$E (\partial_t - H\vec{p} \cdot \nabla_{\vec{p}}) f_{\chi} = \mathcal{C}[f_{\chi}]$$

contains both scatterings and annihilations

Two possible approaches:

fBE

solve numerically
for full $f_{\chi}(p)$

have insight on the distribution
no constraining assumptions

numerically challenging
often an overkill

CBE

consider system of equations
for moments of $f_{\chi}(p)$

0-th moment: n_{χ}
2-nd moment: T_{χ}
...

partially analytic/much easier numerically
manifestly captures all of the relevant physics

finite range of validity
no insight on the distribution

NEW TOOL!

GOING BEYOND THE STANDARD APPROACH

- Home
- Downloads
- Contact



Dark matter Relic Abundance beyond Kinetic Equilibrium

Authors: Tobias Binder, Torsten Bringmann, Michael Gustafsson and Andrzej Hryczuk

DRAKE is a numerical precision tool for predicting the dark matter relic abundance also in situations where the standard assumption of kinetic equilibrium during the freeze-out process may not be satisfied. The code comes with a set of three dedicated Boltzmann equation solvers that implement, respectively, the traditionally adopted equation for the dark matter number density, fluid-like equations that couple the evolution of number density and velocity dispersion, and a full numerical evolution of the phase-space distribution. The code is written in Wolfram Language and includes a Mathematica notebook example program, a template script for terminal usage with the free Wolfram Engine, as well as several concrete example models. DRAKE is a free software licensed under GPL3.

If you use DRAKE for your scientific publications, please cite

- **DRAKE: Dark matter Relic Abundance beyond Kinetic Equilibrium,** Tobias Binder, Torsten Bringmann, Michael Gustafsson and Andrzej Hryczuk, [arXiv:2103.01944]

Currently, a user guide can be found in the Appendix A of this reference. Please cite also quoted other works applying for specific cases.

v1.0 « [Click here to download DRAKE](#)

(March 3, 2021)

<https://drake.hepforge.org>

Applications:

DM relic density for
any (user defined) model*

Interplay between chemical and
kinetic decoupling

Prediction for the DM
phase space distribution

Late kinetic decoupling
and impact on cosmology

see e.g., [1202.5456](#)

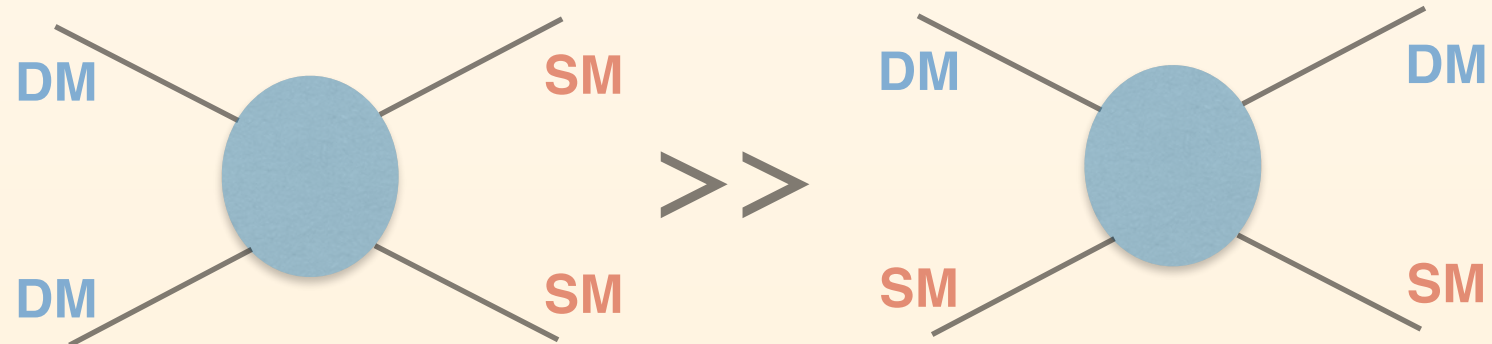
...

(only) prerequisite:
Wolfram Language (or Mathematica)

*at the moment for a single DM species and w/o
co-annihilations... but stay tuned for extensions!

EXAMPLE A: SCALAR SINGLET DM

A)



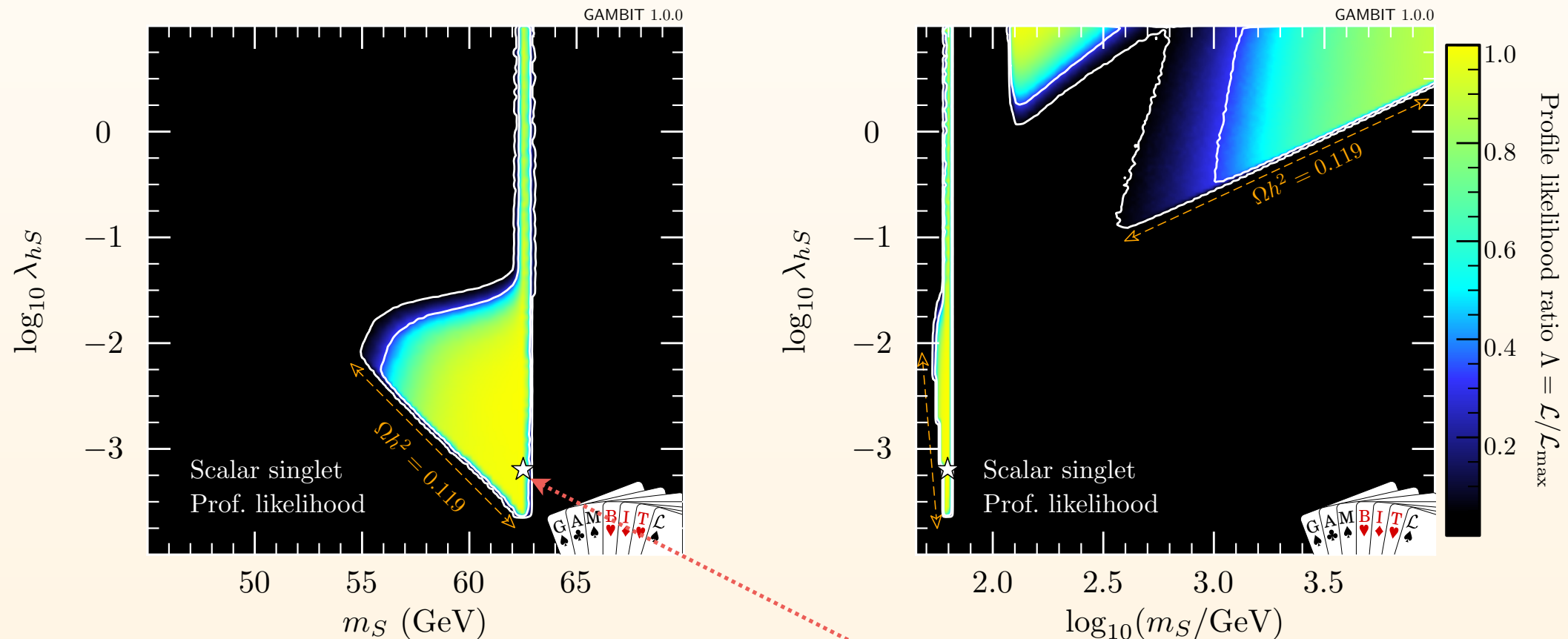
EXAMPLE A

SCALAR SINGLET DM

To the SM Lagrangian add one singlet scalar field S with interactions with the Higgs:

$$\mathcal{L}_S = \frac{1}{2} \partial_\mu S \partial^\mu S - \frac{1}{2} \mu_S^2 S^2 - \frac{1}{2} \lambda_s S^2 |H|^2$$

$$m_s = \sqrt{\mu_S^2 + \frac{1}{2} \lambda_s v_0^2}$$



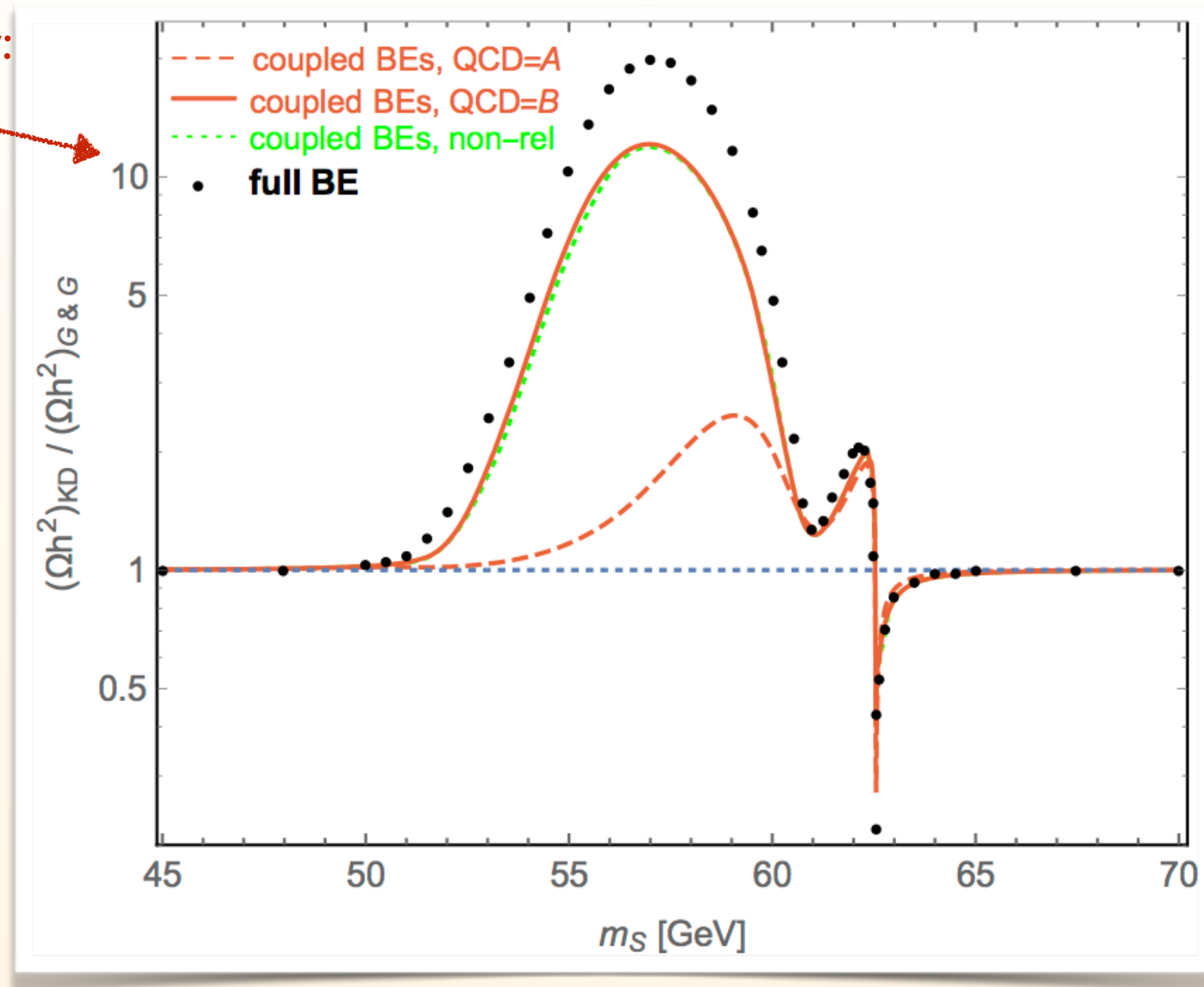
Most of the parameter space excluded, but... even such a simple model is hard to kill

best fit point hides in the **resonance region!**

RESULTS

EFFECT ON THE Ωh^2

effect on relic density:
up to $O(\sim 10)$



[... Freeze-out at few GeV \rightarrow what is the abundance of heavy quarks in QCD plasma?

two scenarios: QCD = A - all quarks are free and present in the plasma down to $T_c = 154$ MeV
 QCD = B - only light quarks contribute to scattering and only down to $4T_c$...] 11

EXAMPLE D: WHEN ADDITIONAL INFLUX OF DM ARRIVES

D) Multi-component dark sectors

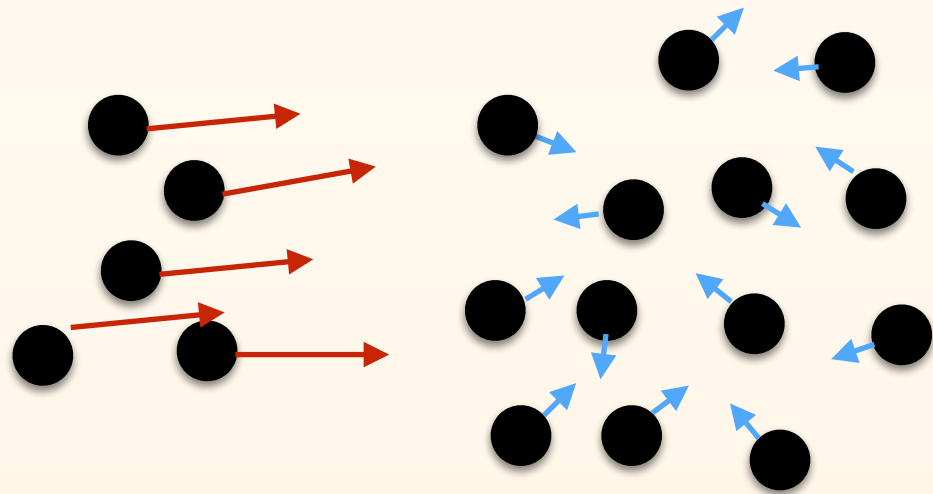
Sudden injection of more DM particles **distorts** $f_\chi(p)$
(e.g. from a decay or annihilation of other states)

- this can **modify the annihilation rate** (if still active)
- how does the **thermalization** due to elastic scatterings happen?

1) DM produced via:

- 1st component from **thermal freeze-out**
- 2nd component from **a decay $\phi \rightarrow \bar{\chi}\chi$**

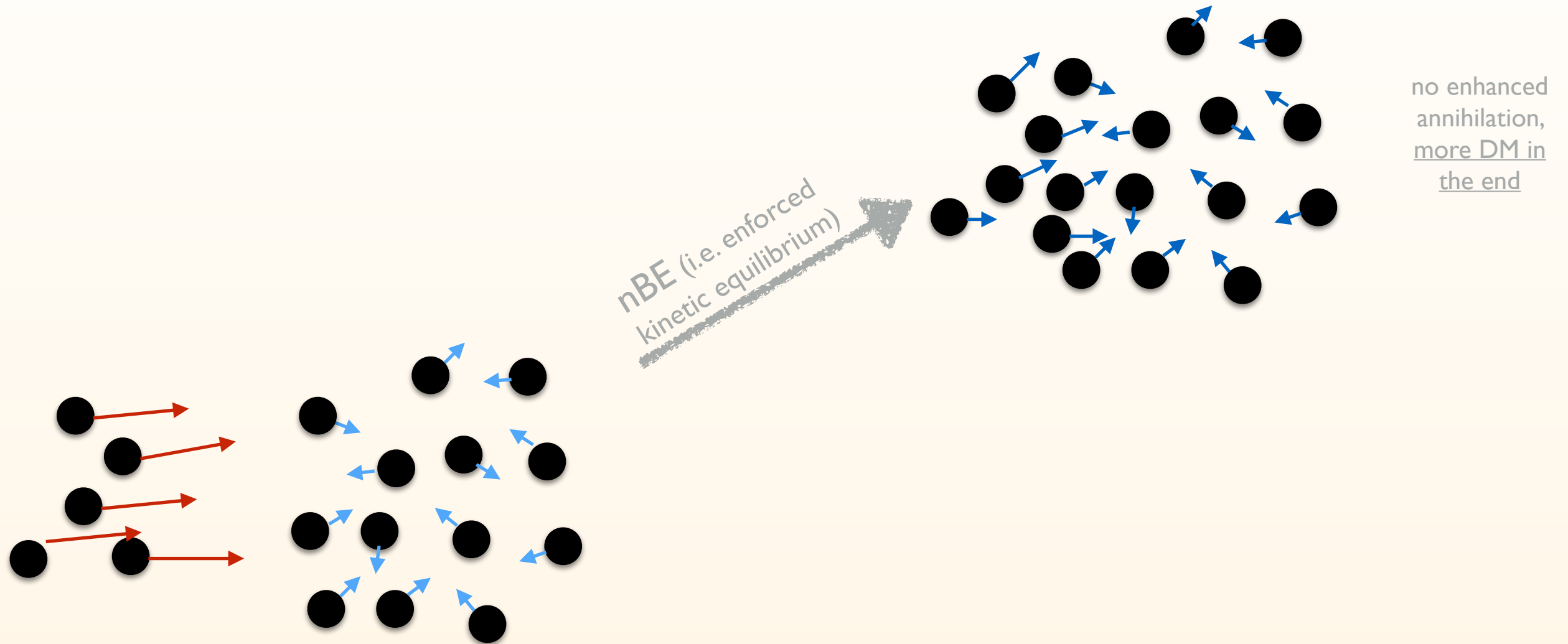
2) DM annihilation has a **threshold**
e.g. $\chi\bar{\chi} \rightarrow f\bar{f}$ with $m_\chi \lesssim m_f$



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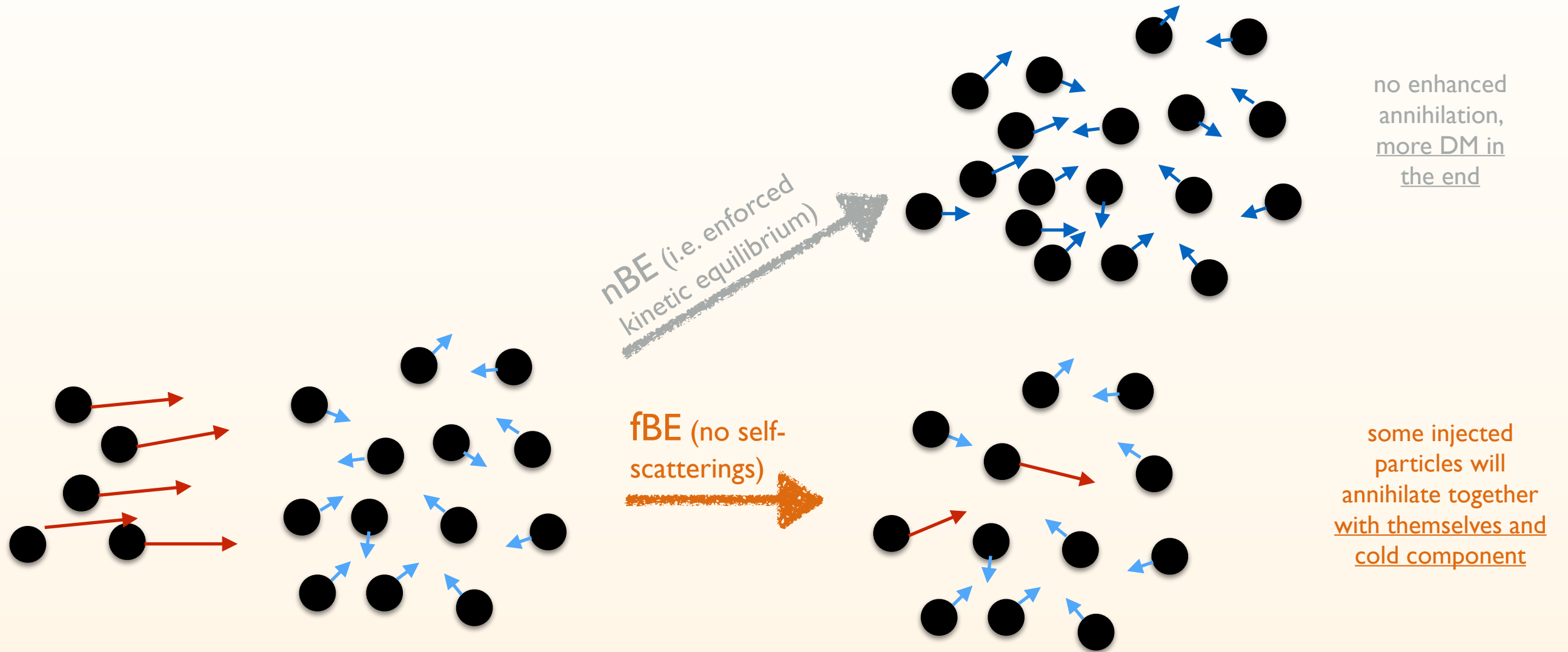
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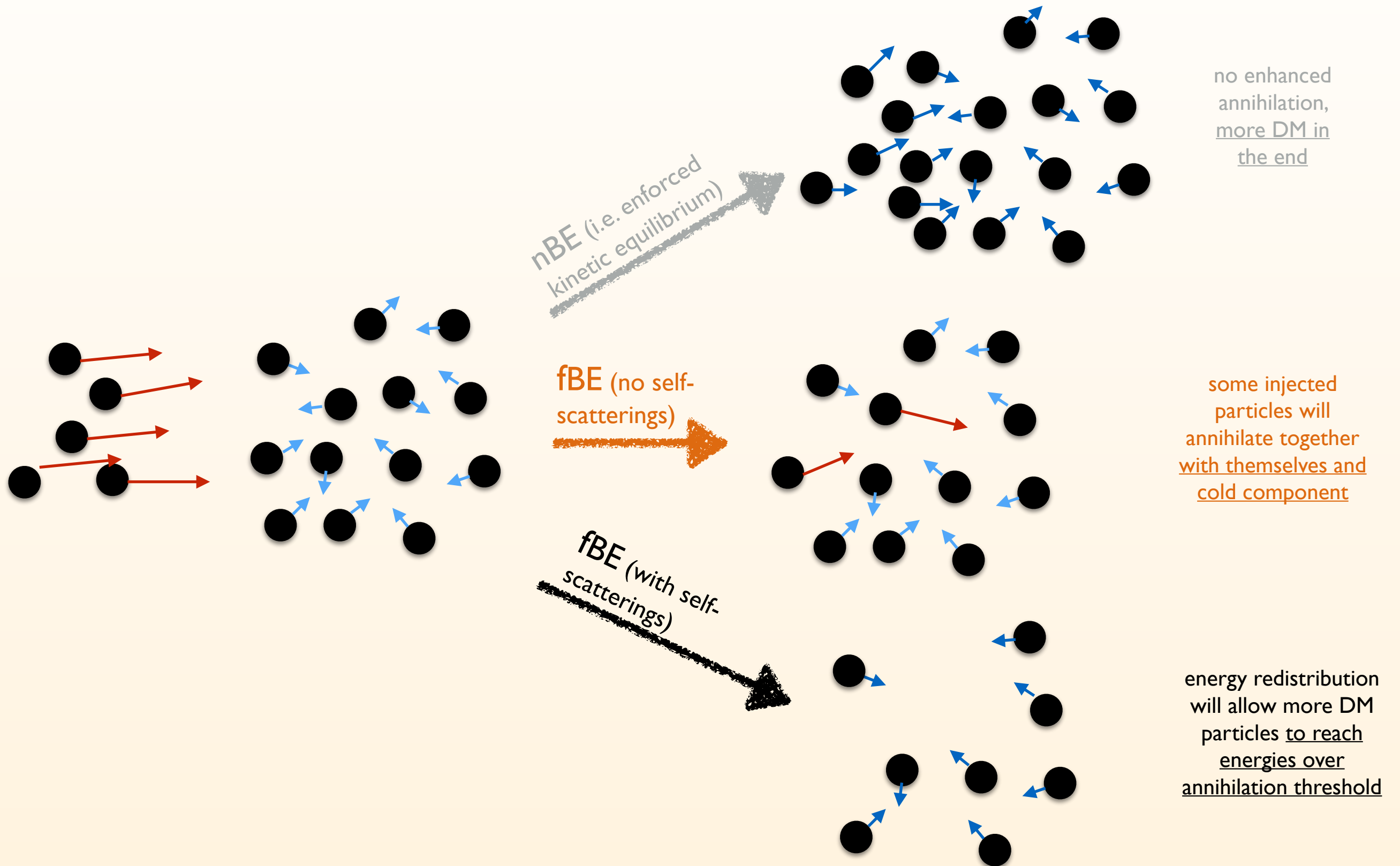
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EXAMPLE EVOLUTION

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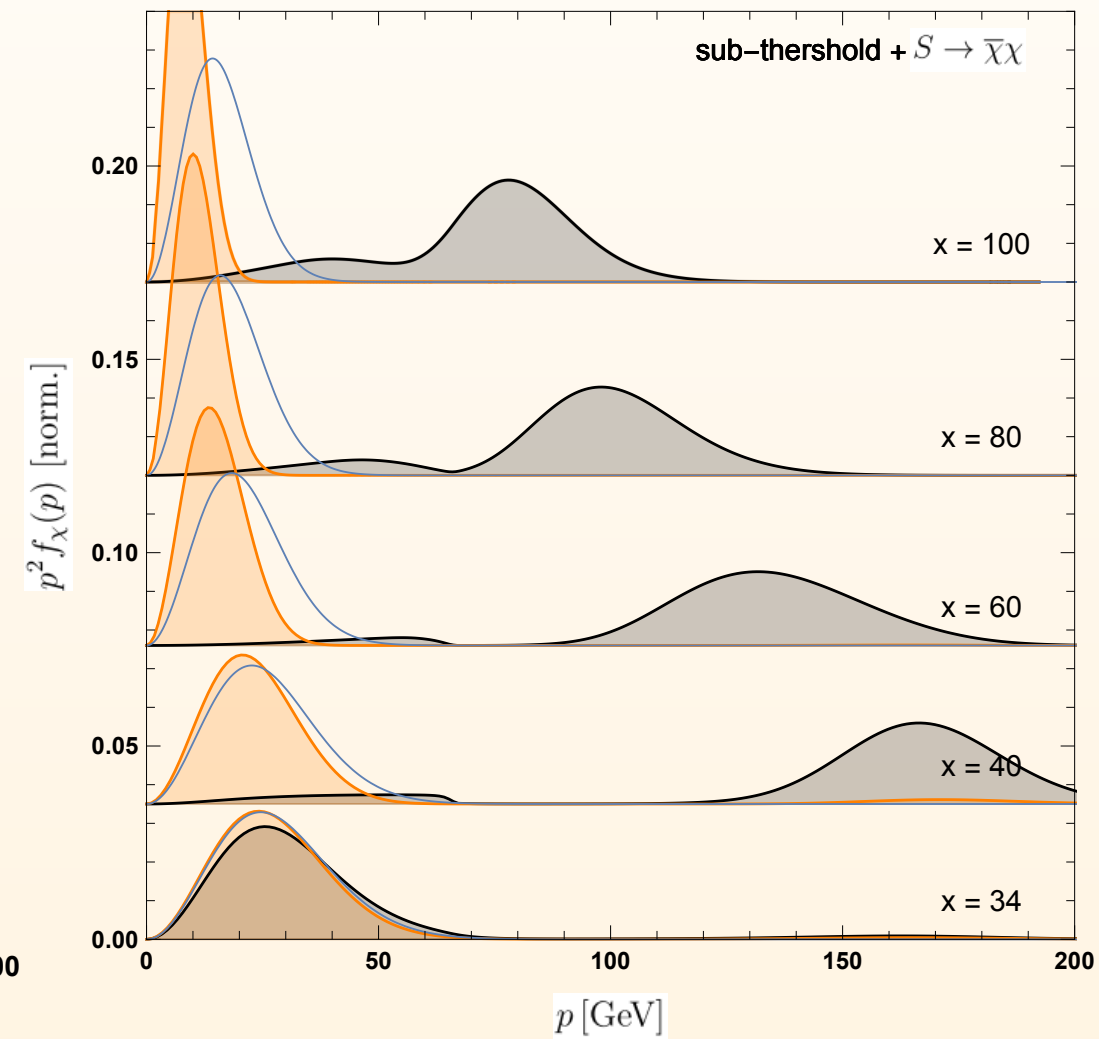
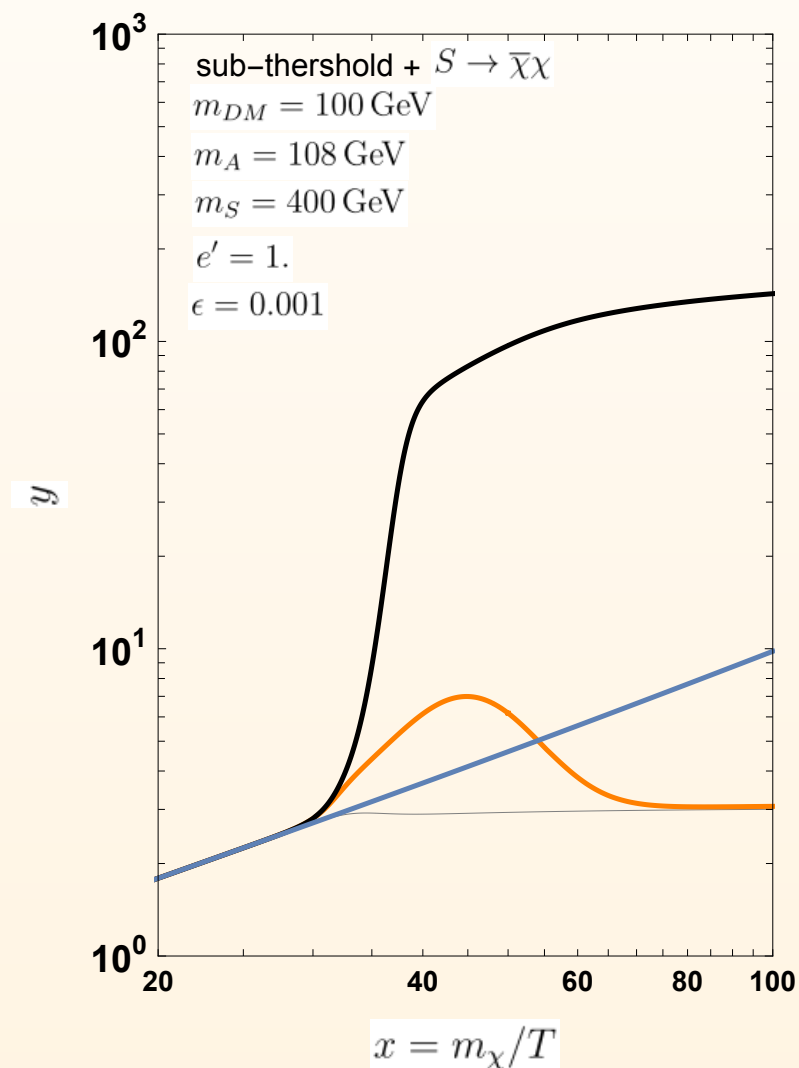
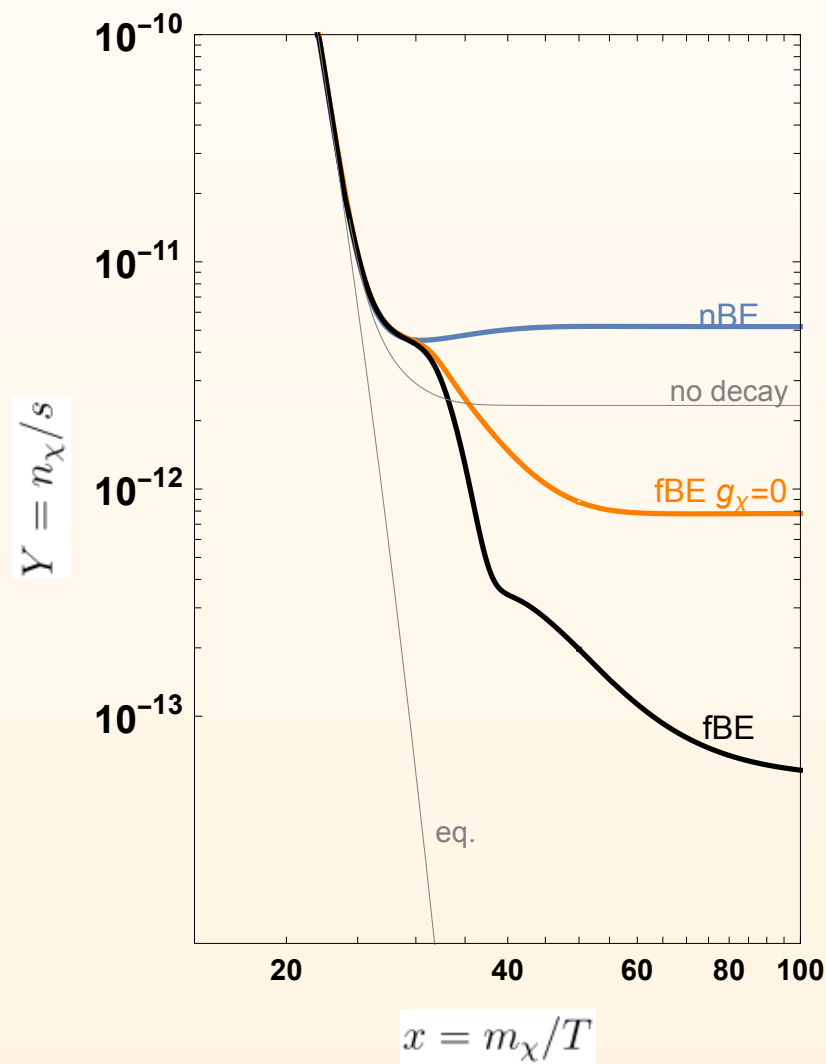
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 e.g. $\chi\bar{\chi} \rightarrow f\bar{f}$ with $m_\chi \lesssim m_f$

$Y \sim$ number density

$y \sim$ temperature

$p^2 f(p) \sim$ momentum distribution



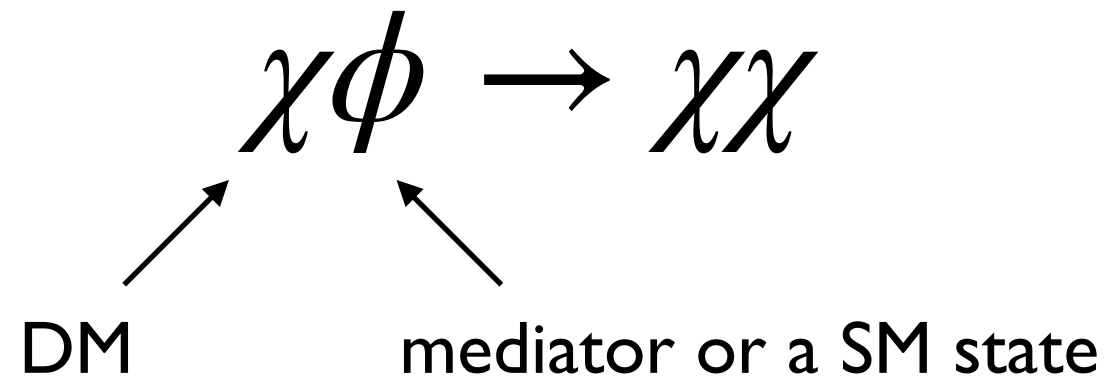
FREEZE-IN:

C) with semi-annihilation process

HOW ABOUT SEMI-PRODUCTION?

AH, Laletin 2104.05684
(see also Bringmann et al. 2103.16572)

Consider process of production that is the **inverse of semi-annihilation**:

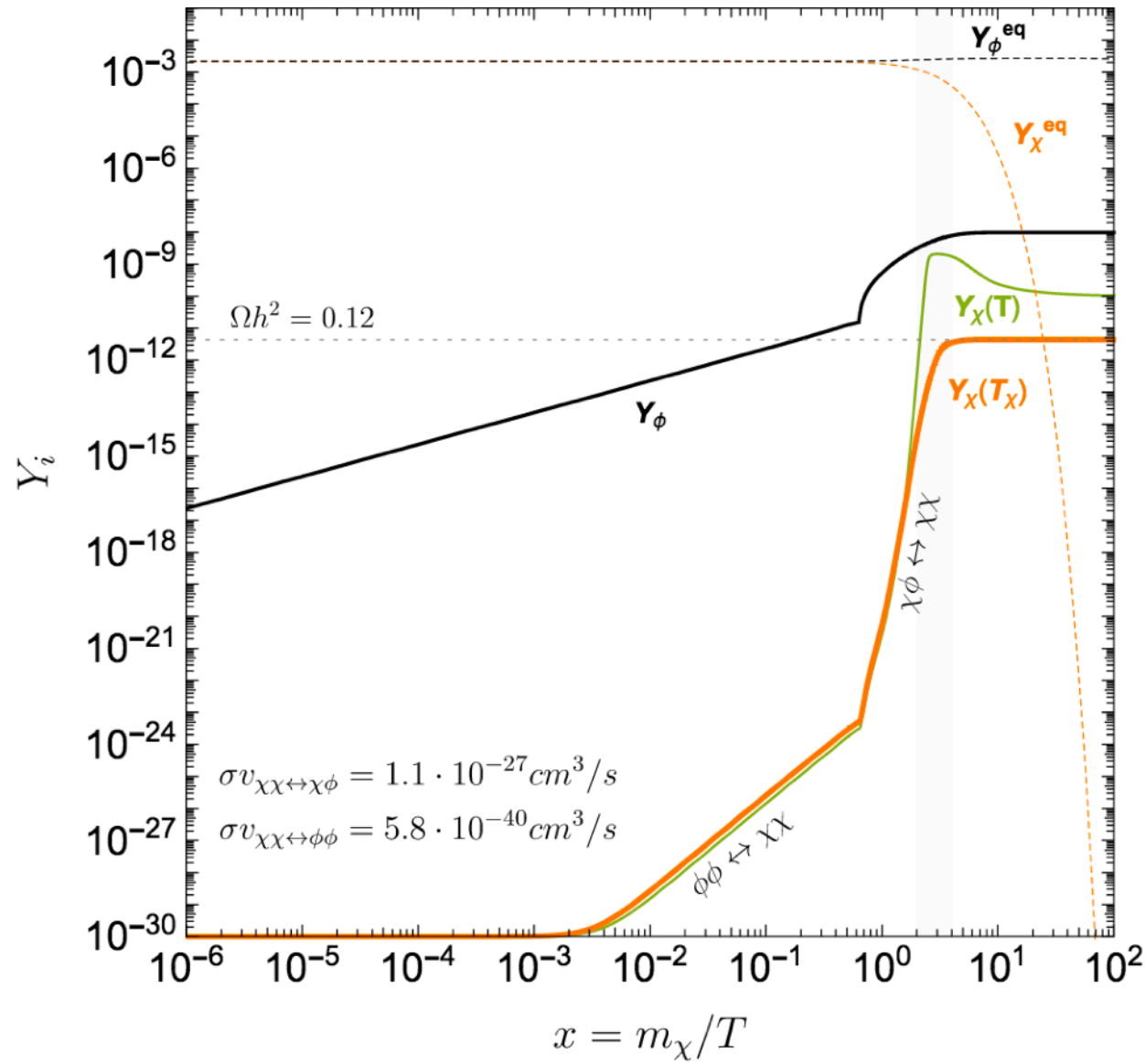


What is different (from the decay/pair-annihilation freeze-in)?

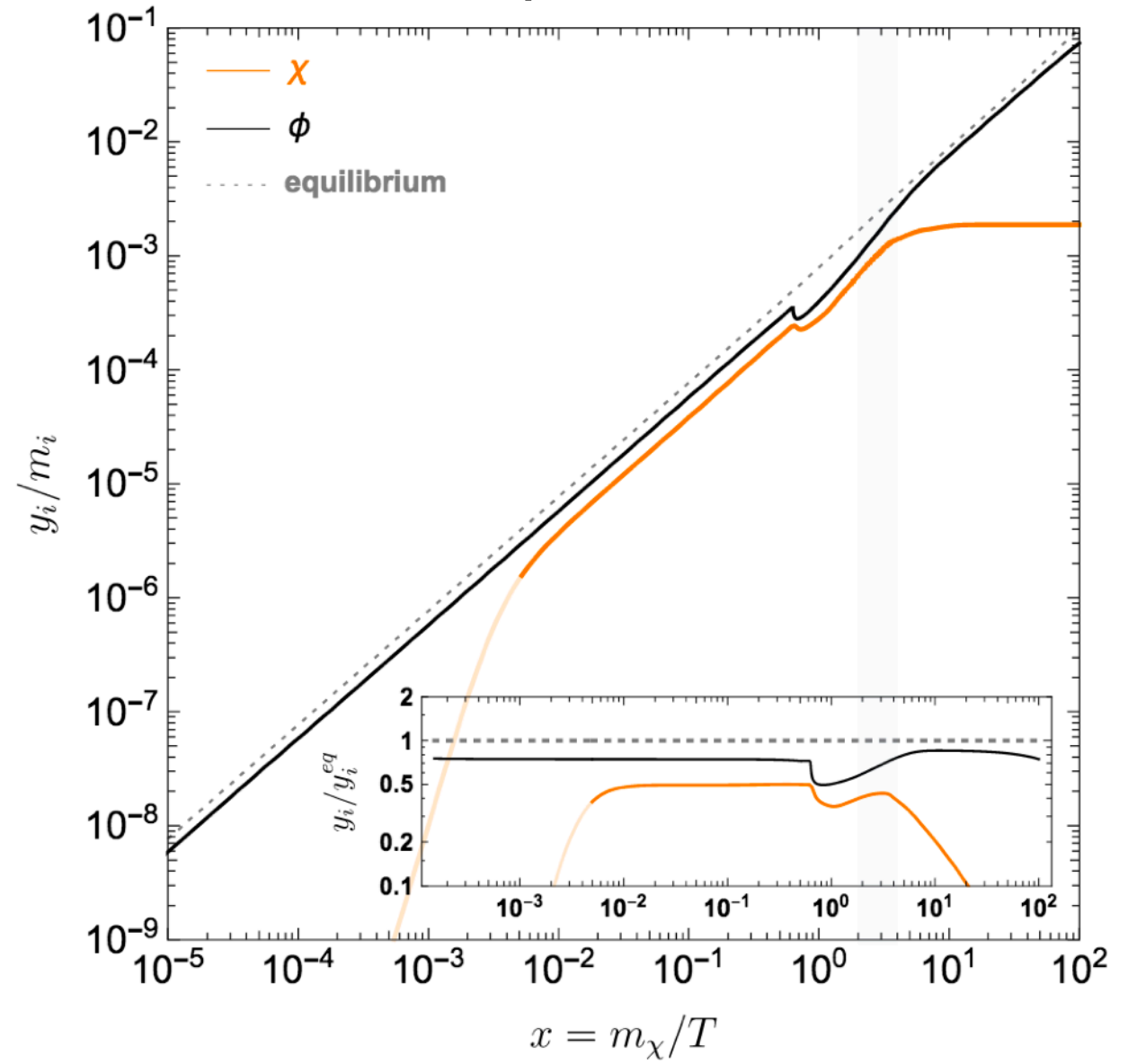
- The production rate is **proportional to the DM density**.
(Smaller initial abundance \rightarrow larger cross section...)
- **Semi-production** modifies the energy of DM particles in a non-trivial way, so the **temperature evolution can affect the relic density**

EVOLUTION

co-moving number density



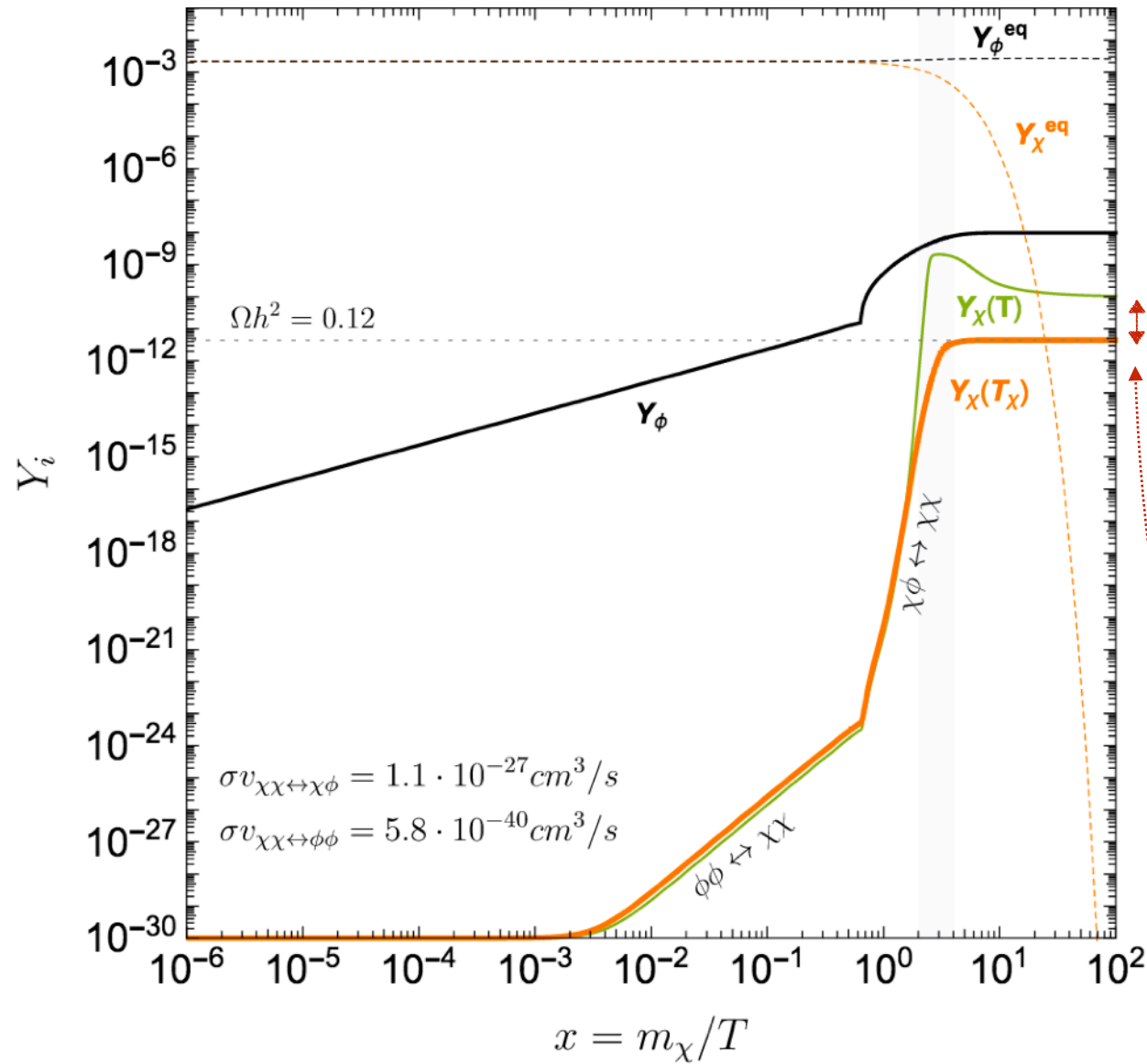
'temperature'



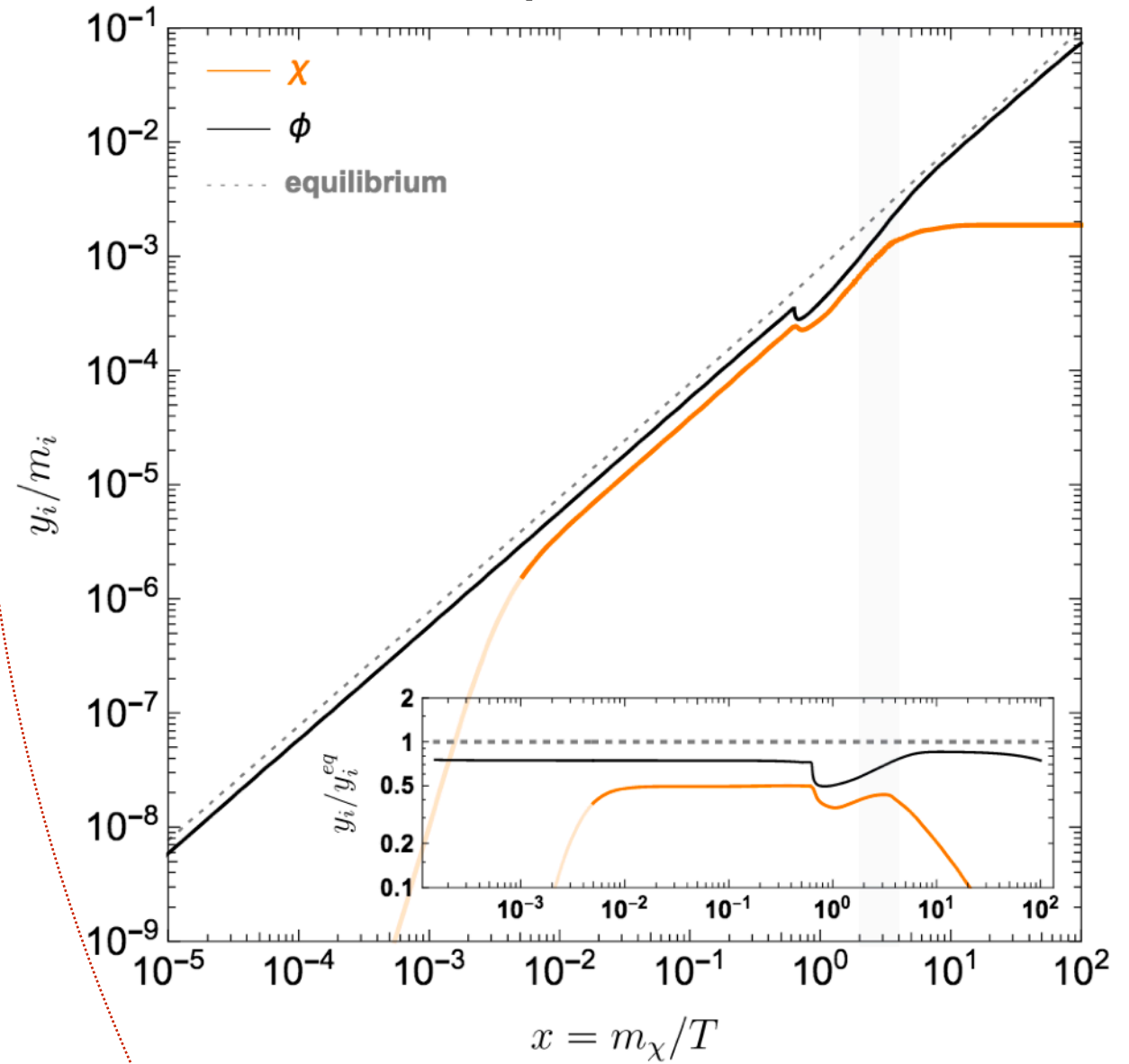
$m_\chi = 100 \text{ GeV}$, $\mu_\phi = 1 \text{ GeV}$, $\lambda_1 = 1.1 \times 10^{-2}$, $\lambda_2 = 10^{-8}$, $\lambda_{h\phi} = 6 \times 10^{-11}$

EVOLUTION

co-moving number density



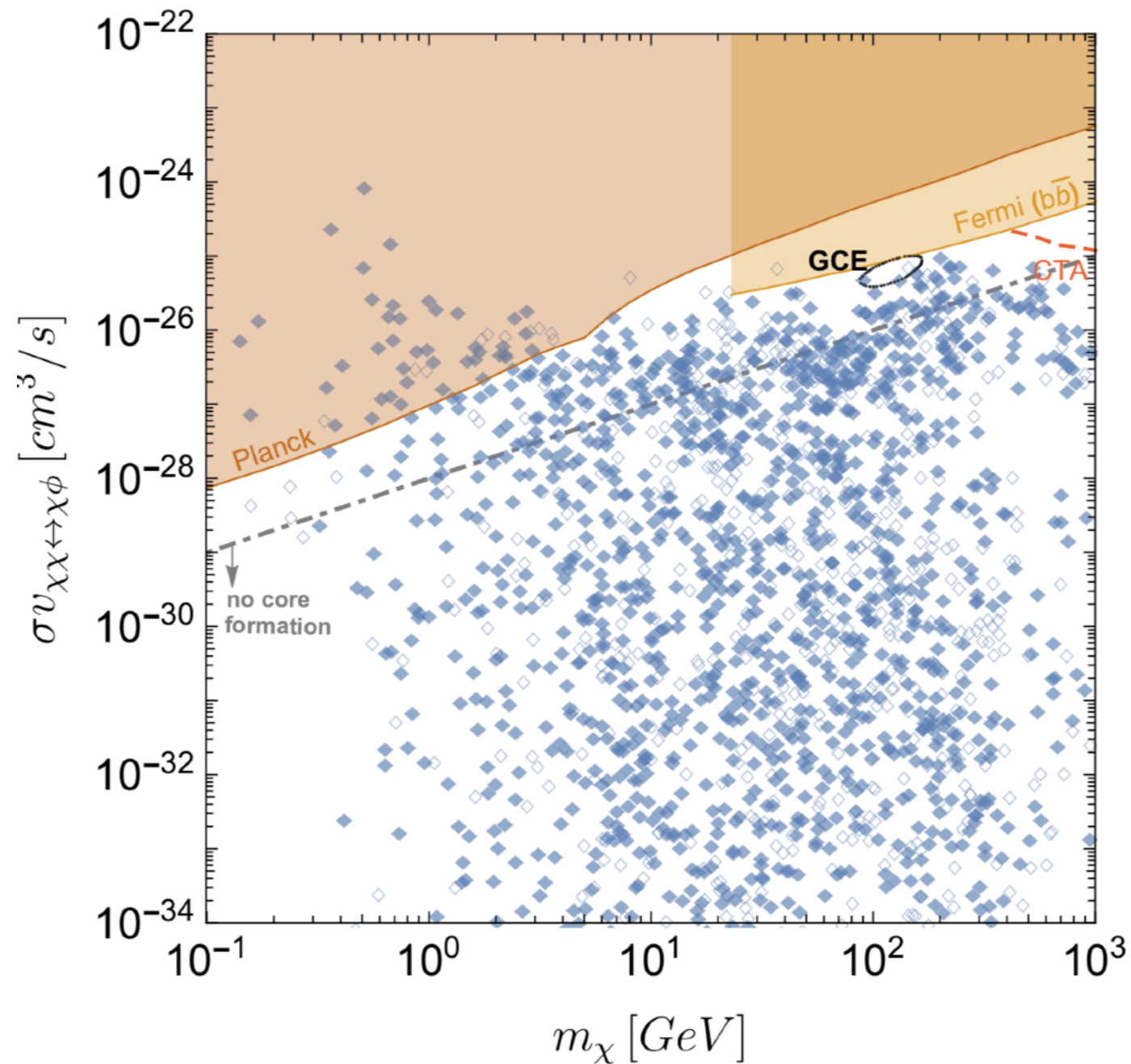
'temperature'



The **full calculation** compared to **one assuming $T_\chi = T$** can differ by more than **order of magnitude!**


$m_\chi = 100 \text{ GeV}$, $\mu_\phi = 1 \text{ GeV}$, $\lambda_1 = 1.1 \times 10^{-2}$, $\lambda_2 = 10^{-8}$, $\lambda_{h\phi} = 6 \times 10^{-11}$

INDIRECT DETECTION



- The results of the scan in the parameter space for the DM production dominated by the **semi-annihilation** processes.
- The **coloured** squares indicate the points, which are **within the reach of the future searches** for the mediator ϕ and the empty ones are beyond these prospects.
- The points above the grey dot-dashed line can potentially **explain the core formation** in dSph [1803.09762]

SUMMARY

- 1. Kinetic equilibrium** is a necessary (often implicit) assumption for standard relic density calculations in all the numerical tools...
...while it is not always warranted!
- 2.** Much more accurate treatment comes from solving the **full phase space Boltzmann equation (fBE)** to obtain result for $f_{\text{DM}}(p)$ where one can study also **self-thermalization from self-scatterings**
- 3.** Introduced **DRAKE**  : a new tool to extend the current capabilities to the regimes **beyond kinetic equilibrium**
- 4. Multi-component sectors**, when studied **at the fBE level**, can reveal quite unexpected behavior